

Optimized Schwartz algorithms for the time-harmonic Maxwell equations discretized by a discontinuous Galerkin method

M. EL BOUAJAJI

NACHOS project-team

INRIA Sophia Antipolis-Méditerranée research center, France

Collaboration with V. Dolean, M. Gander and S. Lanteri

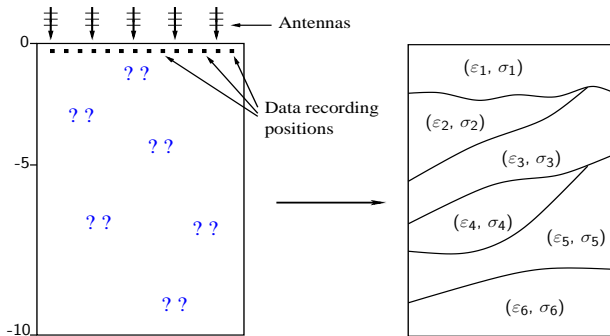


SMAI 2011, Guidel, Bretagne
23 mai-27 mai 2011

- French National Research Agency (ANR) "MAXWELL" project :

Objective : development of a complete microwave imaging system of subsurfaces

- Imaging of a subsurface



- Imaging of a subsurface : solve an inverse problem based in minimizing an objective function
- The minimization = solution of several direct problems at each optimization iteration
- Numerical modeling of electromagnetic wave propagation in heterogeneous media

⇒ Solution of first order time-harmonic Maxwell equations with damping

Discontinuous Galerkin method for the Time-Harmonic Maxwell equations (DGTH) + Domain decomposition method

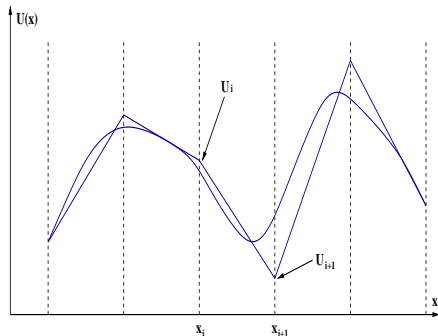
- Time-harmonic Maxwell equations in $\Omega \subset \mathbb{R}^d$ ($d = 2, 3$) :

$$\begin{cases} (i\omega\varepsilon + \sigma)\mathbf{E} - \text{curl}(\mathbf{H}) = -\mathbf{J} \\ i\omega\mu\mathbf{H} + \text{curl}(\mathbf{E}) = \mathbf{0} \end{cases} \Leftrightarrow i\omega\mathbf{Q}\mathbf{W} + \nabla \cdot \mathbf{F}(\mathbf{W}) = \mathbf{S}$$

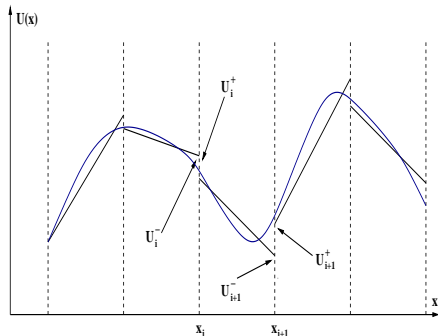
where :

- \mathbf{E} : the electric field, \mathbf{H} : the magnetic field
 - ε the *electric permittivity*, μ the *magnetic permeability*
 - σ is the *electric conductivity*, \mathbf{J} is the applied current density
 - $\mathbf{W} = {}^t(\mathbf{E}, \mathbf{H})$, $\mathbf{S} = {}^t(-\mathbf{J}, \mathbf{0})$
- Boundary conditions on $\partial\Omega = \Gamma = \Gamma^a \cup \Gamma^m$:
 - On Γ^m : $\mathbf{n} \times \mathbf{E} = \mathbf{0}$
 - On Γ^a : $\mathbf{n} \times (\mathbf{E} - \mathbf{E}^{inc}) + z\mathbf{n} \times (\mathbf{n} \times (\mathbf{H} - \mathbf{H}^{inc})) = \mathbf{0}$
 - \mathbf{n} : unitary outwards normal, $z = \sqrt{\mu/\varepsilon}$
 - \mathbf{E}^{inc} , \mathbf{H}^{inc} incident fields

Discontinuous Galerkin method



Continuous P1 interpolation



Discontinuous P1 interpolation

Formulation

- Triangulation : $\mathcal{T}_h = \bigcup_{i=1}^N \tau_i$
- Approximation space :

$$V_h = \{ \mathbf{W} \in L^2(\Omega)^3 \mid \mathbf{W}|_{\tau_i} \in \{\mathcal{P}^{p_i}[\tau_i]^3, \forall \tau_i \in \mathcal{T}_h\}$$

$$\mathcal{P}^{p_i}[\tau_i] = \{\text{polynomial function on } \tau_i \text{ of degree } \leq p_i\}$$

- Variational formulation :

$$\int_{\tau_i} \mathbf{i}\omega \mathbf{Q} \mathbf{W} \varphi \, d\mathbf{x} + \int_{\tau_i} (\nabla \cdot \mathbf{F}(\mathbf{W})) \varphi \, d\mathbf{x} = \int_{\tau_i} \mathbf{S} \varphi \, d\mathbf{x}$$

$$\Leftrightarrow \int_{\tau_i} \mathbf{i}\omega \mathbf{Q} \mathbf{W} \varphi \, d\mathbf{x} - \int_{\tau_i} \nabla \varphi \cdot \mathbf{F}(\mathbf{W}) \, d\mathbf{x} + \int_{\partial \tau_i} (\mathbf{F}(\mathbf{W}) \cdot \mathbf{n}) \varphi \, d\sigma = \int_{\tau_i} \mathbf{S} \varphi \, d\mathbf{x}$$

- Calculation of the boundary term on $\partial \tau_i$: centered or upwind numerical flux

- $(\varphi_{i1}, \varphi_{i2}, \dots, \varphi_{id_i})$ a basis of $\mathcal{P}^{m_i}(\tau_i)$: $\mathbf{W}_i(\mathbf{x}) = \sum_{j=1}^{d_i} \mathbf{W}_{ij} \varphi_{ij}(\mathbf{x})$

- $\mathbf{W}_{ij} \in \mathbb{C}^6$ are the degrees of freedom on τ_i and $d_i = \dim(\mathcal{P}^{m_i}(\tau_i))$

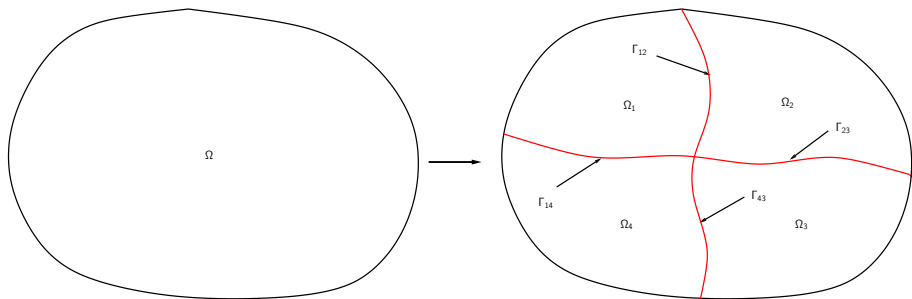
$$\begin{aligned} i\omega Q_i \int_{\tau_i} \mathbf{W}_i \varphi_{ij} d\mathbf{x} &+ \frac{1}{2} \int_{\tau_i} ((\nabla \cdot F(\mathbf{W}_i)) \varphi_{ij} - \nabla \varphi_{ij} \cdot F(\mathbf{W}_i)) d\mathbf{x} \\ &+ \frac{1}{2} \sum_{j \in \mathcal{V}_i} \int_{a_{ij}} (F(\mathbf{W}_j) \cdot \mathbf{n}_{ij}) \varphi_{ij} d\sigma = \int_{\tau_i} \mathbf{S} \varphi_{ij} d\mathbf{x} \end{aligned}$$



Linear system : $\mathcal{A}\mathbf{X} = b$

- \mathcal{A} is of very large dimension, with complex coefficients, sparse and non-hermitian

Domain decomposition method : Motivations

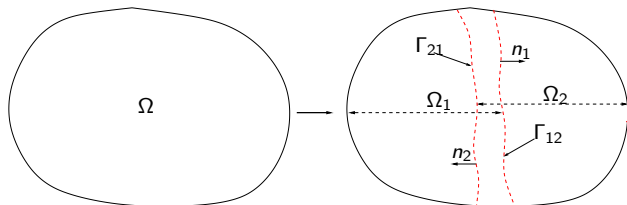


- Each subdomains is assigned to a processor.
- Local problems on the subdomains are solved in parallel : $\mathcal{L}(X^i) = b_i$ on Ω_i
- **Condition on Γ_{ij} ?**
- Schwarz method (iterative process) :

$$\{\mathbf{X}^n = (X^{1,n}, \dots, X^{N,n})\} \mid \begin{cases} X^{i,n} \text{ solution on } \Omega_i \\ \mathcal{B}_i(X^{i,n}) = \mathcal{B}_i(X^{j,n-1}) \text{ on } \Gamma_{ij} \end{cases}$$

Classical Schwarz Method for Maxwell's Equations

■ B. Després, P. Joly and J.E. Roberts, *A domain decomposition method for the harmonic Maxwell equations*, Iterative methods in linear algebra, 1992



Perform for $n = 1, 2, \dots$ the following subdomain iteration

$$\begin{aligned}
 -i\omega\epsilon\mathbf{E}^{1,n} + \text{curl } \mathbf{H}^{1,n} - \sigma\mathbf{E}^{1,n} &= \mathbf{J} && \text{in } \Omega_1 \\
 i\omega\mu\mathbf{H}^{1,n} + \text{curl } \mathbf{E}^{1,n} &= \mathbf{0} && \text{in } \Omega_1 \\
 \mathcal{B}_{n_1}(\mathbf{E}^{1,n}, \mathbf{H}^{1,n}) &= \mathcal{B}_{n_1}(\mathbf{E}^{2,n-1}, \mathbf{H}^{2,n-1}) && \text{on } \Gamma_{12} \\
 -i\omega\epsilon\mathbf{E}^{2,n} + \text{curl } \mathbf{H}^{2,n} - \sigma\mathbf{E}^{2,n} &= \mathbf{J} && \text{in } \Omega_2 \\
 i\omega\mu\mathbf{H}^{2,n} + \text{curl } \mathbf{E}^{2,n} &= \mathbf{0} && \text{in } \Omega_2 \\
 \mathcal{B}_{n_2}(\mathbf{E}^{2,n}, \mathbf{H}^{2,n}) &= \mathcal{B}_{n_2}(\mathbf{E}^{1,n-1}, \mathbf{H}^{1,n-1}) && \text{on } \Gamma_{21}
 \end{aligned}$$

with : $\mathcal{B}_n(\mathbf{E}, \mathbf{H}) := \mathbf{n} \times \frac{\mathbf{E}}{z} + \mathbf{n} \times (\mathbf{H} \times \mathbf{n})$, \mathbf{n} : unitary outwards normal, $z = \sqrt{\mu/\epsilon}$

Convergence analysis

- We consider the domain $\Omega = \mathbb{R}^2$ with the Silver-Muller radiation condition

$$\lim_{r \rightarrow \infty} r (\mathbf{H} \times \mathbf{n} - \mathbf{E}) = 0, \quad r = |\mathbf{x}|, \quad \mathbf{n} = \mathbf{x}/|\mathbf{x}|$$

- Decomposition :

$$\Omega_1 =]-\infty, L] \times \mathbb{R}, \quad \Omega_2 = [0, \infty[\times \mathbb{R}.$$

- Taking a Fourier transform in the y variable of

$$\widehat{\mathbf{W}}(x, k) = (\mathcal{F} \mathbf{W})(x, k) = \int_{\mathbb{R}} \mathbf{W}(x, y) e^{-iky} dy$$

we obtain :

$$\partial_x \begin{pmatrix} \hat{E}_z^{j,n} \\ \hat{H}_y^{j,n} \end{pmatrix} = \begin{pmatrix} 0 & i\omega\mu \\ \frac{k^2 - \omega^2 \varepsilon \mu + i\omega\sigma}{i\omega\mu} & 0 \end{pmatrix} \begin{pmatrix} \hat{E}_z^{j,n} \\ \hat{H}_y^{j,n} \end{pmatrix} =: M \begin{pmatrix} \hat{E}_z^{j,n} \\ \hat{H}_y^{j,n} \end{pmatrix}, \quad j = 1, 2$$

- The eigenvalues of the matrix M , and their corresponding eigenvectors are

$$\lambda_1 = \sqrt{k^2 - \tilde{\omega}^2 + i\tilde{\omega}\sigma Z} := \lambda, \quad \mathbf{v}_1 = \begin{pmatrix} \frac{-i\omega\mu}{\lambda} \\ 1 \end{pmatrix}, \quad \mathbf{v}_2 = \begin{pmatrix} \frac{i\omega\mu}{\lambda} \\ 1 \end{pmatrix}.$$

Convergence analysis

- The solutions of ordinary differential systems are given by :

$$\left(\hat{E}_z^{1,n}, \hat{H}_y^{1,n}\right) = \alpha_1^n \mathbf{v}_1 e^{\lambda(x-L)} + \alpha_2^n \mathbf{v}_2 e^{-\lambda x}, \quad \left(\hat{E}_z^{2,n}, \hat{H}_y^{2,n}\right) = \beta_1^n \mathbf{v}_1 e^{\lambda x} + \beta_2^n \mathbf{v}_2 e^{-\lambda x}.$$

- Using the Silver-Muller radiation condition we have $\alpha_2^n = \beta_1^n = 0$, and inserting the solutions into the interface conditions in we get

$$\alpha_1^n = \frac{\lambda - i\tilde{\omega}}{\lambda + i\tilde{\omega}} \beta_2^{n-1} e^{-\lambda L}, \quad \beta_2^n = \frac{\lambda - i\tilde{\omega}}{\lambda + i\tilde{\omega}} \alpha_1^{n-1} e^{-\lambda L}.$$

Therefore the convergence factor is :

$$\rho(k, \tilde{\omega}, \sigma, L) = \left| \frac{\alpha_1^n}{\alpha_1^{n-2}} \right|^{\frac{1}{2}} = \left| \frac{\lambda - i\tilde{\omega}}{\lambda + i\tilde{\omega}} \frac{\lambda - i\tilde{\omega}}{\lambda + i\tilde{\omega}} e^{-2\lambda L} \right|^{\frac{1}{2}},$$

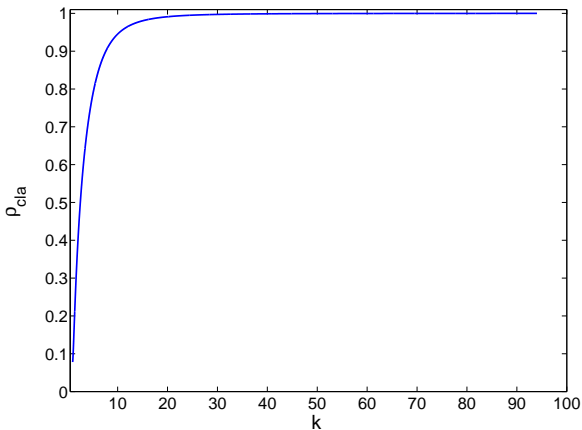


FIG.: Convergence factor ρ_{cla} of the classical Schwarz method as a function of k , for $L = 0$, $\omega = 2\pi$, $\sigma = 2$ and $\mu = \varepsilon = 1$

Better Transmission Conditions Between Subdomains

✎ V. Dolean, L. Gerardo-Giorda and M. J. Gander, *Optimized Schwarz methods for Maxwell equations*, SIAM J. Scient. Comp., 2009

- We propose to modify the algorithm in one aspect only :

$$\begin{aligned}
 -i\omega\varepsilon\mathbf{E}^{1,n} + \operatorname{curl} \mathbf{H}^{1,n} - \sigma\mathbf{E}^{1,n} &= \mathbf{J} && \text{in } \Omega_1 \\
 i\omega\mu\mathbf{H}^{1,n} + \operatorname{curl} \mathbf{E}^{1,n} &= \mathbf{0} && \text{in } \Omega_1 \\
 (\mathcal{B}_{\mathbf{n}_1} + \mathcal{S}_1\mathcal{B}_{\mathbf{n}_2})(\mathbf{E}^{1,n}, \mathbf{H}^{1,n}) &= (\mathcal{B}_{\mathbf{n}_1} + \mathcal{S}_1\mathcal{B}_{\mathbf{n}_2})(\mathbf{E}^{2,n-1}, \mathbf{H}^{2,n-1}) && \text{on } \Gamma_{12} \\
 -i\omega\varepsilon\mathbf{E}^{2,n} + \operatorname{curl} \mathbf{H}^{2,n} - \sigma\mathbf{E}^{2,n} &= \mathbf{J} && \text{in } \Omega_2 \\
 i\omega\mu\mathbf{H}^{2,n} + \operatorname{curl} \mathbf{E}^{2,n} &= \mathbf{0} && \text{in } \Omega_2 \\
 (\mathcal{B}_{\mathbf{n}_2} + \mathcal{S}_2\mathcal{B}_{\mathbf{n}_1})(\mathbf{E}^{2,n}, \mathbf{H}^{2,n}) &= (\mathcal{B}_{\mathbf{n}_2} + \mathcal{S}_2\mathcal{B}_{\mathbf{n}_1})(\mathbf{E}^{1,n-1}, \mathbf{H}^{1,n-1}) && \text{on } \Gamma_{21},
 \end{aligned}$$

where \mathcal{S}_j , $j = 1, 2$ are tangential, possibly pseudo-differential operators.

- *How to Choose the operators \mathcal{S}_j ?*

Optimal convergence

- Taking a Fourier transform in the y variable, we get :

$$\rho(k, \tilde{\omega}, \sigma, L) = \left| \frac{\lambda - i\tilde{\omega} + \mathcal{F}(\mathcal{S}_1)(\lambda + i\tilde{\omega})}{\lambda + i\tilde{\omega} + \mathcal{F}(\mathcal{S}_1)(\lambda - i\tilde{\omega})} \frac{\lambda - i\tilde{\omega} + \mathcal{F}(\mathcal{S}_2)(\lambda + i\tilde{\omega})}{\lambda + i\tilde{\omega} + \mathcal{F}(\mathcal{S}_2)(\lambda - i\tilde{\omega})} e^{-2\lambda L} \right|^{\frac{1}{2}}.$$

- Note that if we choose $\mathcal{F}(\mathcal{S}_1) = \mathcal{F}(\mathcal{S}_2) = -\frac{\lambda - i\tilde{\omega}}{\lambda + i\tilde{\omega}}$, then $\rho = 0$ and the algorithm converges in two iterations.
- Unfortunately, it is difficult to use these operators in practice. The symbols $\mathcal{F}(\mathcal{S}_j)$ contain square roots $\left(\lambda = \sqrt{k^2 - \tilde{\omega}^2 + i\tilde{\omega}\sigma Z} \right)$
- **This choice of $\mathcal{F}(\mathcal{S}_j)$ corresponds to nonlocal operators \mathcal{S}_j .**

Approximation of optimal conditions

- Local operators : if we approximate λ by polynomial in ik , we will get differential operators in physical space.
- The symbols $\mathcal{F}(\mathcal{S}_j)$ can be written in several forms :

$$\begin{aligned} \mathcal{F}(\mathcal{S}_j) &= -\frac{\lambda - i\tilde{\omega}}{\lambda + i\tilde{\omega}} = -\frac{k^2 + i\tilde{\omega}\sigma Z}{k^2 - 2\tilde{\omega}^2 + i\tilde{\omega}\sigma Z + 2i\tilde{\omega}\lambda} \\ &= -\frac{k^2 - 2\tilde{\omega}^2 + i\tilde{\omega}\sigma Z - 2i\tilde{\omega}\lambda}{k^2 + i\tilde{\omega}\sigma Z} \end{aligned}$$

- Approximation of λ by a constant polynomial $s \in \mathbb{C}$.

Lemma

- If S_j , $j = 1, 2$, have the Fourier symbols $\sigma_j = \mathcal{F}(S_j) = -\frac{s - i\tilde{\omega}}{s + i\tilde{\omega}}$, $s \in \mathbb{C}$, then the Schwarz algorithm has the convergence factor

$$\rho_2(k, \tilde{\omega}, \sigma, Z, L, s) = \left| \left(\frac{\sqrt{k^2 - \tilde{\omega}^2 + i\tilde{\omega}\sigma Z} - s}{\sqrt{k^2 - \tilde{\omega}^2 + i\tilde{\omega}\sigma Z} + s} \right) e^{-\sqrt{k^2 - \tilde{\omega}^2 + i\tilde{\omega}\sigma Z} L} \right|.$$

- If S_j , $j = 1, 2$ have the Fourier symbols $\sigma_j = -\frac{k^2 + i\tilde{\omega}\sigma Z}{k^2 - 2\tilde{\omega}^2 + i\tilde{\omega}\sigma Z + 2i\tilde{\omega}s}$, then the Schwarz algorithm has the convergence factor

$$\rho_3(k, \tilde{\omega}, \sigma, Z, L, s) = \left| \left(\frac{\sqrt{k^2 - \tilde{\omega}^2 + i\tilde{\omega}\sigma Z} - i\tilde{\omega}}{\sqrt{k^2 - \tilde{\omega}^2 + i\tilde{\omega}\sigma Z} + i\tilde{\omega}} \right) \right| \rho_2(k, \tilde{\omega}, \sigma, Z, L, s).$$

- If S_j , $j = 1, 2, ,$ have the Fourier symbols $\sigma_j = \mathcal{F}(S_j) = -\frac{s_j - i\tilde{\omega}}{s_j + i\tilde{\omega}}$, $s_2 \in \mathbb{C}$, then the Schwarz algorithm has the convergence factor

$$\rho_4(k, \tilde{\omega}, \sigma, Z, L, s_1, s_2) = \left| \prod_{l=1}^2 \frac{\sqrt{k^2 - \tilde{\omega}^2 + i\tilde{\omega}\sigma Z} - s_l}{\sqrt{k^2 - \tilde{\omega}^2 + i\tilde{\omega}\sigma Z} + s_l} e^{-2\sqrt{k^2 - \tilde{\omega}^2 + i\tilde{\omega}\sigma Z} L} \right|^{\frac{1}{2}}.$$

Hierarchy of Schwarz algorithms

Algorithm 1	$\mathcal{F}(\mathcal{S}_1) = \mathcal{F}(\mathcal{S}_2) = 0$
Algorithm 2	$\mathcal{F}(\mathcal{S}_1) = \mathcal{F}(\mathcal{S}_2) = \frac{s - i\tilde{\omega}}{s + i\tilde{\omega}}$
Algorithm 3	$\mathcal{F}(\mathcal{S}_1) = \mathcal{F}(\mathcal{S}_2) - \frac{k^2 + i\tilde{\omega}\sigma Z}{k^2 - 2\tilde{\omega}^2 + i\tilde{\omega}\sigma Z + 2i\tilde{\omega}s}$
Algorithm 4	$\mathcal{F}(\mathcal{S}_1) = \frac{s_1 - i\tilde{\omega}}{s_1 + i\tilde{\omega}}$ $\mathcal{F}(\mathcal{S}_2) = \frac{s_2 - i\tilde{\omega}}{s_2 + i\tilde{\omega}}$
Algorithm 5	$\mathcal{F}(\mathcal{S}_1) = -\frac{k^2 + i\tilde{\omega}\sigma Z}{k^2 - 2\tilde{\omega}^2 + i\tilde{\omega}\sigma Z + 2i\tilde{\omega}s_1}$ $\mathcal{F}(\mathcal{S}_2) = -\frac{k^2 + i\tilde{\omega}\sigma Z}{k^2 - 2\tilde{\omega}^2 + i\tilde{\omega}\sigma Z + 2i\tilde{\omega}s_2}$

How to Choose the Complex Constants s , s_1 and s_2

- In order to obtain some efficient algorithms, we choose σ_j , $j = 1, 2$, such that ρ , is minimum over a range of frequencies
- We need to solve some min-max problem

$$\min_{s \in \mathbb{C}} \max_{k \in K} \rho(k, \tilde{\omega}, \sigma, Z, L, s), \quad \min_{s_1, s_2 \in \mathbb{C}} \max_{k \in K} \rho(k, \tilde{\omega}, \sigma, Z, L, s_1, s_2), \quad K = [k_{min}, k_{max}]$$

- Solution of min-max problem is based on equioscillation

$$\rho(\bar{k}_i, \tilde{\omega}, \sigma, Z, L, s) = \rho(\bar{k}_l, \tilde{\omega}, \sigma, Z, L, s)$$

How to Choose the Complex Constants s , s_1 and s_2

Theorem

For $\sigma > 0$, treating all frequencies k , we find for the optimal $s = p(1 + i)$ based on equioscillation :

- *Non-overlapping case*

$$p^* = \frac{(\omega\sigma\mu)^{\frac{1}{4}}\sqrt{C}}{2^{\frac{1}{4}}\sqrt{h}} \rightarrow \rho_2^* = 1 - \frac{2^{\frac{3}{4}}(\omega\sigma\mu)^{\frac{1}{4}}\sqrt{h}}{\sqrt{C}}, \quad L = 0$$

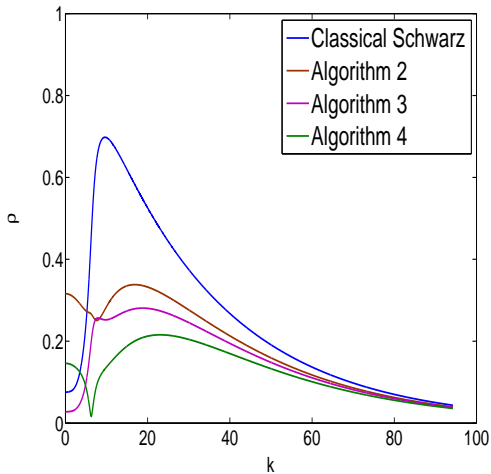
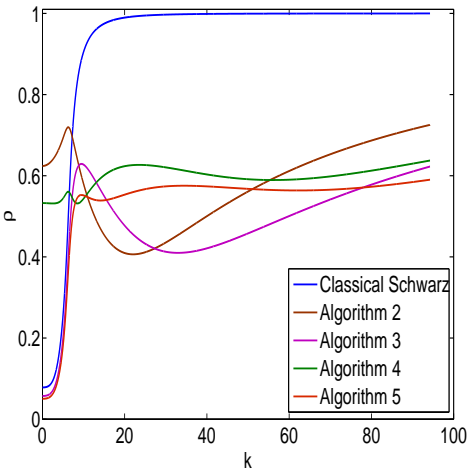
- *Overlapping case*

$$p^* = \frac{(2\omega\sigma\mu)^{1/3}}{2h^{\frac{1}{3}}} \rightarrow \rho_2^* = 1 - 2^{\frac{7}{6}}(\omega\sigma\mu)^{1/6}h^{\frac{1}{3}}, \quad L = h$$

where $k_{max} = \frac{C}{h}$

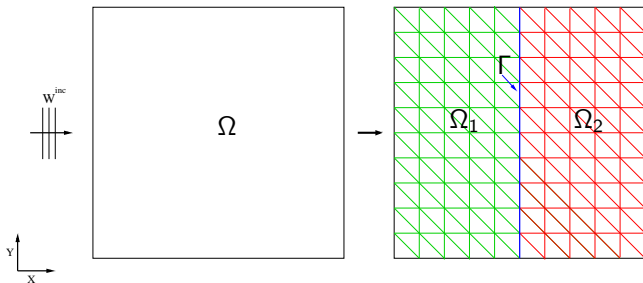
with overlap, $L = h$		
Algorithm	ρ	parameters
1	$1 - \frac{4}{3} (9\omega^4 \sigma^2 \mu^3 \varepsilon)^{\frac{1}{8}} h^{\frac{3}{4}}$	none
2	$1 - 2^{\frac{7}{6}} (\omega \sigma \mu)^{1/6} h^{\frac{1}{3}}$	$\rho = \frac{(2\omega \sigma \mu)^{1/3}}{2h^{\frac{1}{3}}}$
3	$1 - \frac{2^{\frac{17}{10}} (\omega^4 \varepsilon \mu^3 \sigma^2)^{\frac{1}{20}} h^{\frac{3}{10}}}{3^{\frac{3}{10}}}$	$\rho = \frac{2^{\frac{2}{5}} (\omega^4 \varepsilon \mu^3 \sigma^2)^{\frac{1}{10}}}{3^{\frac{3}{5}} h^{\frac{2}{5}}}$
4	$1 - 4\sqrt{2} (\omega \sigma \mu)^{\frac{1}{10}} h^{\frac{1}{5}}$	$\rho_1 = \frac{(\omega \sigma \mu)^{\frac{1}{5}}}{2h^{\frac{1}{5}}}, \rho_2 = \frac{(\omega \sigma \mu)^{\frac{2}{5}}}{2h^{\frac{1}{5}}}$
5	$1 - \frac{2^{\frac{23}{8}} (\omega^4 \varepsilon \mu^3 \sigma^2)^{\frac{1}{32}} h^{\frac{3}{16}}}{3^{\frac{3}{16}}}$	$\rho_1 = \frac{(\omega^4 \varepsilon \mu^3 \sigma^2)^{\frac{1}{16}}}{2^{\frac{1}{4}} 3^{\frac{3}{8}} h^{\frac{5}{8}}}, \rho_2 = \frac{\sqrt{2} (\omega^4 \varepsilon \mu^3 \sigma^2)^{\frac{1}{8}}}{3^{\frac{3}{4}} h^{\frac{1}{4}}}$
without overlap, $L = 0$		
1	$1 - \frac{\omega^2 \sigma \sqrt{\mu^3 \varepsilon}}{C^3} h^3$	none
2	$1 - \frac{2^{\frac{3}{4}} (\omega \sigma \mu)^{\frac{1}{4}} \sqrt{h}}{\sqrt{C}}$	$\rho = \frac{(\omega \sigma \mu)^{\frac{1}{4}} \sqrt{C}}{2^{\frac{1}{4}} \sqrt{h}}$
3	$1 - \frac{2^{\frac{11}{7}} (\omega^4 \varepsilon \mu^3 \sigma^2)^{\frac{1}{14}} h^{\frac{3}{7}}}{3^{\frac{3}{7}} C^{\frac{3}{7}}}$	$\rho = \frac{2^{\frac{4}{7}} (\omega^4 \varepsilon \mu^3 \sigma^2)^{\frac{1}{14}} C^{\frac{4}{7}}}{3^{\frac{3}{7}} h^{\frac{4}{7}}}$
4	$1 - \frac{(2\omega \sigma \mu)^{\frac{1}{8}} h^{\frac{1}{4}}}{C^{\frac{1}{4}}}$	$\rho_1 = \frac{(2\omega \sigma \mu)^{\frac{1}{8}} C^{\frac{3}{4}}}{h^{\frac{3}{4}}}, \rho_2 = \frac{(2\omega \sigma \mu)^{3/8} C^{1/4}}{2h^{\frac{1}{4}}}$
5	$1 - \frac{2^{\frac{14}{13}} (\omega^4 \varepsilon \mu^3 \sigma^2)^{\frac{1}{26}} h^{\frac{3}{13}}}{3^{\frac{3}{13}} C^{\frac{3}{13}}}$	$\rho_1 = \frac{2^{\frac{8}{13}} (\omega^4 \varepsilon \mu^3 \sigma^2)^{\frac{1}{26}} C^{\frac{10}{13}}}{3^{\frac{3}{13}} h^{\frac{10}{13}}}, \rho_2 = \frac{2^{\frac{11}{13}} (\omega^4 \varepsilon \mu^3 \sigma^2)^{\frac{1}{26}} C^{\frac{4}{13}}}{3^{\frac{9}{13}} h^{\frac{4}{13}}}$

Convergence factor for the optimized Schwarz algorithms



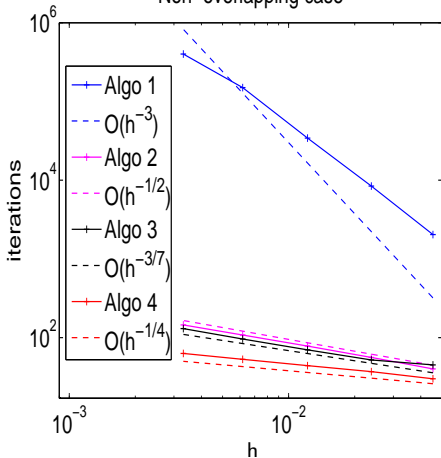
Performance for a Two Subdomain Decomposition

$$\left\{ \begin{array}{l} \Omega = (0, 1)^2, \varepsilon = \mu = 1, \sigma = 5, \omega = 2\pi \\ \mathbf{W}^{inc} = (H_x^{inc}, H_y^{inc}, E_z^{inc}) = \left(\frac{k_y}{\mu\omega}, \frac{-k_x}{\mu\omega}, 1 \right) e^{-i\mathbf{k}\cdot\mathbf{x}} \\ \mathbf{k} = (k_x, k_y) = (\omega\sqrt{\varepsilon - i\frac{\sigma}{\omega}}, 0), \mathbf{x} = (x, y) \\ \Omega_1 = (0, \frac{1}{2}) \times (0, 1), \Omega_2 = (\frac{1}{2}, 1) \times (0, 1), \Gamma = \Omega_1 \cap \Omega_2 \end{array} \right.$$

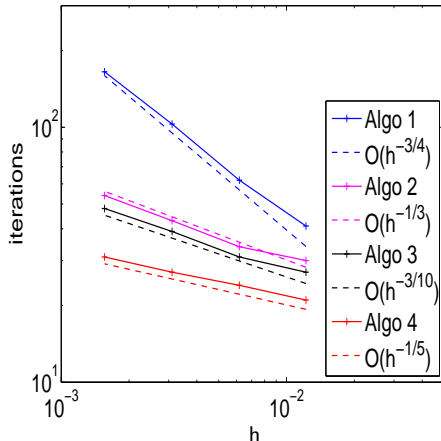


Performance for a Two Subdomain Decomposition

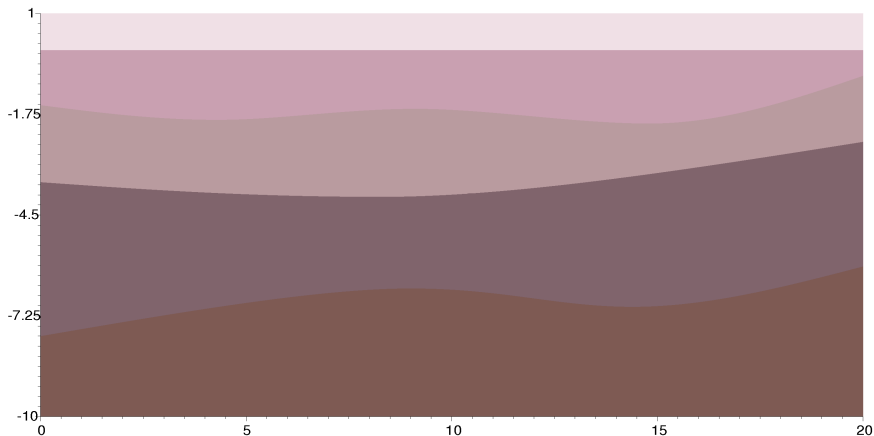
Non-overlapping case



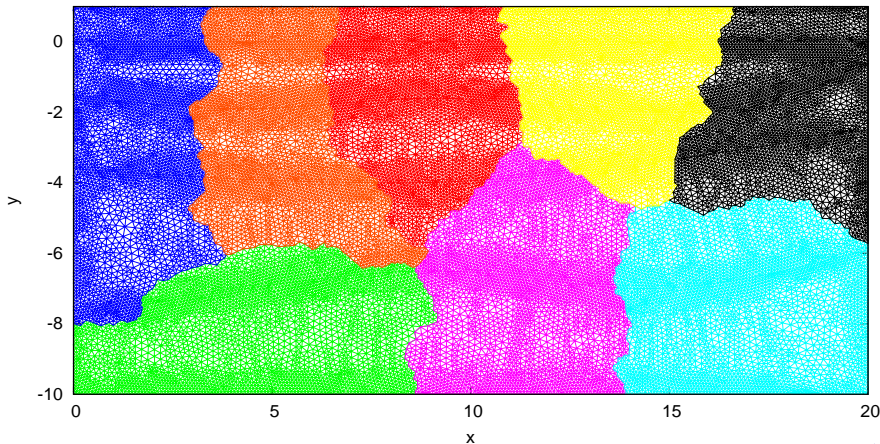
Overlapping case



Electromagnetic wave propagation in heterogeneous subsurface



Electromagnetic wave propagation in heterogeneous subsurface



Electromagnetic wave propagation in heterogeneous subsurface

# processor	2	4	8	16
Algorithm 1	59	72	80	85
Algorithm 2	39	47	53	61
Algorithm 4	35	41	49	57

- Imaging of a subsurface : solve an inverse problem based in minimizing an objective function
- The minimization = solution of several direct problems at each optimization iteration

