# Mathématiorsas <br> universitė PARIS-SACLAY 

On Optimal Transport, variational Mean Field Games and beyond

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## Overview

1. Optimal Transport

The three formulations of quadratic Optimal Transport
On geodesics
2. The Schrödinger problem

The three formulations of Schrödinger
3. Variational Mean Field Games

Eulerian and Lagrangian formulation for MFG with quadratic Hamiltonian
4. Towards a numerical method:

Optimal Transport

## Optimal Transportation Theory

Let $\mu, \nu \in \mathcal{P}(\Omega), \Omega$ compact subset of $R^{n}$, the Optimal Transport (OT) problem is defined as follows

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\begin{equation*}
(\mathcal{M K}) \quad E_{c}(\mu, \nu)=\inf \left\{\mathcal{E}_{c}(\gamma) \mid \gamma \in \Pi(\mu, \nu)\right\} \tag{1}
\end{equation*}
$$

where $\Pi(\mu, \nu):=\left\{\gamma \in \mathcal{P}\left(\Omega^{2}\right) \mid \quad \pi_{1, \sharp} \gamma=\mu, \pi_{2, \sharp} \gamma=\nu\right\}$ and

$$
\mathcal{E}_{c}(\gamma):=\int c(x, y) d \gamma(x, y)
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Solution à la Monge: the transport plan $\gamma$ is deterministic (or à la Monge) if $\gamma=(I d, T)_{\sharp} \mu$ where $T_{\sharp} \mu=\nu$.


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The $(\mathcal{M} \mathcal{K})$ problem admits a dual formulation:

$$
\begin{equation*}
\sup \{\mathcal{J}(\phi, \psi) \mid(\phi, \psi) \in \mathcal{K}\} \tag{2}
\end{equation*}
$$

where

$$
\mathcal{J}(\phi, \psi):=\int_{\Omega} \phi d \mu(x)+\int_{\Omega} \psi d \nu(y)
$$

and $\mathcal{K}$ is the set of bounded and continuous functions $\phi, \psi$ such that $\phi(x)+\psi(y) \leq c(x, y)$.

The three formulations of quadratic Optimal Transport

The static: $\inf \left\{\left.\int_{X \times Y} \frac{1}{2}|x-y|^{2} d \gamma \right\rvert\, \gamma \in \Pi(\mu, \nu)\right\}$

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The dynamic (Eulerian), aka the Benamou-Brenier formulation

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\begin{array}{r}
\inf \int_{0}^{1} \int_{\Omega} \frac{1}{2}\left|v_{t}\right|^{2} \rho_{t} d x d t \quad \text { s.t. } \partial_{t} \rho_{t}+\operatorname{div}\left(\rho_{t} v_{t}\right)=0 \\
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\sup \left\{\int_{\Omega} \varphi(1, x) d \nu-\left.\int \varphi(0, x) d \mu\left|\partial_{t} \varphi+\frac{1}{2}\right| \nabla \varphi\right|^{2} \leq 0\right\}
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The dynamic (Lagrangian) $\left(C=H^{1}([0,1] ; \Omega)\right.$ and $\left.e_{t}:[0,1] \rightarrow \Omega\right)$

$$
\inf \left\{\left.\int_{C} \int_{0}^{1} \frac{1}{2}|\dot{\omega}|^{2} d t d Q(\omega) \right\rvert\, Q \in \mathcal{P}(C),\left(e_{0}\right)_{\sharp} Q=\mu,\left(e_{1}\right)_{\sharp} Q=\nu\right\}
$$

## On geodesics

- Quadratic optimal transport is indeed a distance between probability measures, aka the $\mathcal{W}_{2}^{2}$ Wasserstein distance, and $\left(\mathcal{P}(\Omega), \mathcal{W}_{2}\right)$ is a metric space;
- Gives a way to compare and interpolate between probability measures.
- Consider the optimal solutions for the three formulations $\gamma^{\star}, Q^{\star}, \rho_{t}^{\star}$ then

$$
\pi_{t}(x, y)_{\sharp \gamma^{\star}}=\left(e_{t}\right)_{\sharp} Q^{\star}=\rho_{t}^{\star},
$$

where $\pi_{t}(x, y)=(1-t) x+t y$ and $\rho_{t}$ is the geodesic between $\mu$ and $\nu$, the so called McCann's interpolant.


- Hidden convexity: look at convexity along the Wasserstein geodesics.


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$\Rightarrow$ The Schrödinger problem
Definition (Relative entropy)
Let $\rho$ and $\pi$ probability measures on $\Omega$ then the relative entropy is defined as

$$
\mathcal{H}(\rho \mid \pi)= \begin{cases}\int_{\Omega \times \Omega}\left(\log \left(\frac{d \rho(x, y)}{d \pi(x, y)}\right)-1\right) d \rho(x, y), & \text { if } \rho \ll \pi \\ +\infty, & \text { otherwise }\end{cases}
$$

The three formulations of Schrödinger
The static: $\inf \left\{\left.\int_{\Omega \times \Omega} \frac{1}{2}|x-y|^{2} d \gamma+\varepsilon \mathcal{H}(\gamma \mid \mu \otimes \nu) \right\rvert\, \gamma \in \Pi(\mu, \nu)\right\}$

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where $R^{\varepsilon}$ is the Wiener measure $R^{\varepsilon}:=\int \delta_{x+B^{\varepsilon}} d x$ of variance $\varepsilon$.
Remark: static formulation can be defined for a general cost function $c(x, y)$.

## The "bridge" between quadratic Monge-Kantorovich and Schrödinger

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Convergence to optimal transport solution for quadratic cost as $\varepsilon \rightarrow 0$


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Sinkhorn algorithm (we will see it in 15 minutes...if I do not talk too much).

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## Variational Mean Field Games

## Lagrangian formulation for 1st order MFG

Consider a first order MFG system then we have the following "equivalence" (see (Lasry and Lions 2007))

A MFG system

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\left\{\begin{array}{l}
-\partial_{t} \varphi+\frac{1}{2}|\nabla \varphi|^{2}=g(x, \rho), \varphi(1, x)=\Psi(x) \\
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The (Eulerian) Variational Formulation

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\begin{gathered}
\inf \int_{0}^{1} \int_{\Omega}\left(\frac{1}{2}\left|v_{t}\right|^{2} \rho_{t}+G\left(x, \rho_{t}\right)\right) d x d t+F\left(\rho_{1}\right) \quad \text { s.t. } \partial_{t} \rho_{t}+\operatorname{div}\left(\rho_{t} v_{t}\right)=0 \\
\rho(0, \cdot)=\rho_{0}
\end{gathered}
$$

where $G$ is the anti-derivative of $g$ w.r.t its second variable and $F\left(\rho_{1}\right)=\int_{\Omega} \Psi d \rho_{1}$ is a final cost.

## Lagrangian formulation for 1st order MFG

Consider a first order MFG system then we have the following "equivalence" (see (Lasry and Lions 2007))

A MFG system

$$
\left\{\begin{array}{l}
-\partial_{t} \varphi+\frac{1}{2}|\nabla \varphi|^{2}=g(x, \rho), \varphi(1, x)=\Psi(x) \\
\partial_{t} \rho-\operatorname{div}(\rho \nabla \varphi)=0, \rho(0, \cdot)=\rho_{0} .
\end{array}\right.
$$

The (Eulerian) Variational Formulation

$$
\begin{gathered}
\inf \int_{0}^{1} \int_{\Omega}\left(\frac{1}{2}\left|v_{t}\right|^{2} \rho_{t}+G\left(x, \rho_{t}\right)\right) d x d t+F\left(\rho_{1}\right) \quad \text { s.t. } \partial_{t} \rho_{t}+\operatorname{div}\left(\rho_{t} v_{t}\right)=0 \\
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The (Lagrangian) Variational Formulation (J.-D. Benamou, G. Carlier, and Santambrogio 2017)

$$
\inf _{Q \in \mathcal{P}(C)}\left\{\left.\int_{C} \int_{0}^{1} \frac{1}{2}|\dot{\omega}|^{2} d t d Q+\int_{0}^{1} \mathcal{G}\left(e_{t, \sharp} Q\right) d t+F\left(e_{1, \sharp} Q\right) \right\rvert\,\left(e_{0}\right)_{\sharp} Q=\rho_{0}\right\},
$$

where $\mathcal{G}(\rho)=\int G(x, \rho) d x$ if $\rho \ll \mathcal{L}$ and $+\infty$ otherwise.

## A Lagrangian formulation via Entropy minimization

A MFG system

$$
\left\{\begin{array}{l}
-\partial_{t} \varphi-\frac{\varepsilon}{2} \Delta \phi+\frac{1}{2}|\nabla \varphi|^{2}=g(x, \rho), \varphi(1, x)=\Psi(x) \\
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where $G$ is the anti-derivative of $g$ w.r.t its second variable and $F\left(\rho_{1}\right)=\int_{\Omega} \psi d \rho_{1}$ is a final cost.

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$$

Towards a numerical method:

## The discretised (in time) problems

We can solve the Lagrangian problems by firstly discretising them in time as follows

$$
\begin{aligned}
& \text { Regularized 1st order MFG } \\
& \begin{aligned}
& \inf \int_{\Omega^{N+1}} K_{N} d Q_{N}\left(x_{0}, \cdots, x_{N}\right)+\varepsilon \mathcal{H}\left(Q_{N} \mid \mathcal{L}\right)+\sum_{i=1}^{N-1} \int_{\Omega} G\left(x, \pi_{i, \sharp} Q_{N}\right) d x_{i}+F\left(\pi_{N, \sharp} Q_{N}\right) \\
& \text { s.t. } Q_{N} \in \mathcal{P}\left(\Omega^{N+1}\right), \pi_{0, \sharp} Q_{N}=\rho_{0},
\end{aligned}
\end{aligned}
$$

where $K_{N}=\frac{1}{2 N} \sum_{i=0}^{N-1}\left|x_{i+1}-x_{i}\right|^{2}$ and $\pi_{i}: \Omega^{N+1} \rightarrow \Omega$ is the canonical projection.

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2nd order MFG

$$
\inf \left\{\mathcal{H}\left(Q_{N} \mid R_{N}^{\varepsilon}\right)+\sum_{i=1}^{T-1} \int_{\Omega} G\left(x, \pi_{i, \sharp} Q_{N}\right) d x_{i}+F\left(\pi_{N, \sharp} Q_{N}\right) \mid \pi_{0, \sharp} Q_{N}=\rho_{0}\right\},
$$

where $R_{N}^{\varepsilon} \stackrel{\text { def }}{=} \prod_{n=0}^{N} \xi_{n, n+1}$ and $\xi_{i j}=\exp ^{-\frac{\left|x_{i}-x_{j}\right|^{2}}{2 N \varepsilon}}$.

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Remarks: (i) for small $\varepsilon$ the regularized 1st MFG approximate the unreg pb (ii) both problems can be re-written in the same way.

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Remarks: (i) for small $\varepsilon$ the regularized 1st MFG approximate the unreg pb (ii) both problems can be re-written in the same way.
IDEA: an alternate coordinate ascent algorithm (equivalent to a generalised Sinkhorn)

## Dual formulation

## The dual problem

The Lagrangian problem can be re-written as the following dual optimization problem:
$\sup _{\left(\phi_{0}, \cdots, \phi_{N}\right)}-\tilde{F}^{\star}\left(-\phi_{0}\right)-\frac{1}{N} \sum_{k=1}^{N-1} G^{\star}\left(-\phi_{k}\right)-F^{\star}\left(-\phi_{N}\right)-\int\left(\exp \left(\oplus_{k=0}^{N} \phi_{k}\right)-1\right) R_{N}^{\epsilon}$,
where $\tilde{F}^{\star}, G^{\star}$ and $F^{\star}$ are the Legendre transforms of $i_{\rho 0}, G$ and $G$.

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## Proposition ((J.-D. Benamou, G. Carlier, S. Di Marino, and L. Nenna 2018))

Strong duality holds, namely sup $=$ inf.
Moreover, denoting by $\phi_{k}^{\star}$ and $Q_{N}^{\star}$ the optimal solutions to the dual and primal problem respectively, it follows that the unique solution to the primal has the form

$$
\begin{equation*}
Q_{N}^{\star}\left(x_{0}, \cdots, x_{N}\right):=\left(\otimes_{k=0}^{N} e^{\phi_{k}^{\star}\left(x_{k}\right)}\right) R_{N}^{\varepsilon}\left(x_{0}, \cdots, x_{N}\right) . \tag{3}
\end{equation*}
$$

## A coordinate ascent algorithm (or generalised Sinkhorn)

Generalizing a result of (Peyré 2015; Chizat, Peyré, B. Schmitzer, and Vialard 2016), we get the iterative method computing a sequence of potentials (denoted with the superscripts ${ }^{(n)}$ ) :

$$
\begin{aligned}
& \phi_{0}^{(n)}:=\operatorname{argmax}_{\phi}-\tilde{F}^{\star}(-\phi)-\int \exp (\phi) l_{k}^{\phi} \mathrm{d} x_{1} \cdots \mathrm{~d} x_{N}, \\
& \phi_{k}^{(n)}:=\operatorname{argmax}_{\phi}-\frac{1}{N} G^{\star}(-\phi)-\int \exp (\phi) l_{k}^{\phi} \mathrm{d} x_{0} \cdots \mathrm{~d} x_{k-1} \mathrm{~d} x_{k+1} \cdots \mathrm{~d} x_{N} \text { for } k \neq 0, N, \\
& \phi_{N}^{(n)}:=\operatorname{argmax}_{\phi}-F^{\star}(-\phi)-\int \exp (\phi) l_{k}^{u} \mathrm{~d} x_{0} \cdots \mathrm{~d} x_{N-1},
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I_{k}^{\phi}:=\exp \left(\oplus_{i=0}^{k-1} \phi_{i}^{(n)}\right) \exp \left(\oplus_{i=k+1}^{N} \phi_{i}^{(n-1)}\right) R^{N} .
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## Remarks:

- For many interesting energies $F$ and $G$, the relaxed maximizations can be computed point-wise in space and analytically;


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## Remarks:

- For many interesting energies $F$ and $G$, the relaxed maximizations can be computed point-wise in space and analytically;
- with the same method one can compute dynamic optimal transport by imposing $G=0$ and $F=i_{\rho_{1}}$


## Sinkhornizing the world!!

- Wasserstein Barycenter (Jean-David Benamou, Guillaume Carlier, Cuturi, Luca Nenna, and Peyré 2015);
- Matching for teams (Luca Nenna 2016);
- Optimal transport with capacity constraint (Jean-David Benamou, Guillaume Carlier, Cuturi, Luca Nenna, and Peyré 2015);
- Partial Optimal Transport (Jean-David Benamou, Guillaume Carlier, Cuturi, Luca Nenna, and Peyré 2015; Chizat, Peyré, B. Schmitzer, and Vialard 2016);
- Multi-Marginal Optimal Transport (Luca Nenna 2016; J.-D. Benamou, G. Carlier, and L. Nenna 2016; Jean-David Benamou, Guillaume Carlier, and Luca Nenna 2018; Jean-David Benamou, Guillaume Carlier, Cuturi, Luca Nenna, and Peyré 2015);
- Wasserstein Gradient Flows (JKO) (Peyré 2015);
- Unbalanced Optimal Transport (Chizat, Peyré, B. Schmitzer, and Vialard 2016);
- Cournot-Nash equilibria (Blanchet, Guillaume Carlier, and Luca Nenna 2017)
- Mean Field Games (J.-D. Benamou, G. Carlier, S. Di Marino, and L. Nenna 2018);
- Grand Canonical Optimal transport (Simone Di Marino, Lewin, and Luca Nenna 2022);
- and more...


## Dynamic OT

## Data:

- $T=32$ time steps;
- grid: uniform discretization of $[0,1]^{2}$ with $N \times N$ points $N=250$;
- Given final density;




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## Planning MFG with obstacles on the torus, behaviour as $\varepsilon \rightarrow 0$

## Data:

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Thank you!

