

On Optimal Transport, variational Mean Field Games and beyond

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1. Optimal Transport

The three formulations of quadratic Optimal Transport On geodesics

2. The Schrödinger problem

The three formulations of Schrödinger

3. Variational Mean Field Games

Eulerian and Lagrangian formulation for MFG with quadratic Hamiltonian

4. Towards a numerical method:

Optimal Transport

Let $\mu, \nu \in \mathcal{P}(\Omega)$, Ω compact subset of \mathbb{R}^n , the Optimal Transport (OT) problem is defined as follows

$$(\mathcal{MK}) \quad \mathcal{E}_{c}(\mu,\nu) = \inf \left\{ \mathcal{E}_{c}(\gamma) \mid \gamma \in \Pi(\mu,\nu) \right\}$$
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where $\Pi(\mu,\nu) := \left\{ \gamma \in \mathcal{P}(\Omega^{2}) \mid \pi_{1,\sharp}\gamma = \mu, \ \pi_{2,\sharp}\gamma = \nu \right\}$ and
 $\mathcal{E}_{c}(\gamma) := \int c(x,y) d\gamma(x,y).$

Solution à la Monge: the transport plan γ is deterministic (or à la Monge) if $\gamma = (Id, T)_{\sharp}\mu$ where $T_{\sharp}\mu = \nu$.



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The (\mathcal{MK}) problem admits a dual formulation: $\sup \{ \mathcal{J}(\phi, \psi) \mid (\phi, \psi) \in \mathcal{K} \}.$ (2)

where

$$\mathcal{J}(\phi,\psi) := \int_{\Omega} \frac{\phi}{\partial t} d\mu(x) + \int_{\Omega} \frac{\psi}{\partial t} d\nu(y)$$

and \mathcal{K} is the set of bounded and continuous functions ϕ, ψ such that $\phi(x) + \psi(y) \leq c(x, y)$.

The three formulations of quadratic Optimal Transport

The static:
$$\inf \left\{ \int_{X \times Y} \frac{1}{2} |x - y|^2 d\gamma \mid \gamma \in \Pi(\mu, \nu) \right\}$$

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The dynamic (Eulerian), aka the Benamou-Brenier formulation

$$\inf \int_0^1 \int_\Omega \frac{1}{2} |v_t|^2 \rho_t dx dt \quad s.t. \ \partial_t \rho_t + \operatorname{div}(\rho_t v_t) = 0$$
$$\rho(0, \cdot) = \mu, \ \rho(1, \cdot) = \nu$$

And its "dual"

$$\sup\Big\{\int_{\Omega}\varphi(1,x)d\nu-\int\varphi(0,x)d\mu\ |\partial_t\varphi+\frac{1}{2}|\nabla\varphi|^2\leq 0\Big\}$$

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The dynamic (Lagrangian) $(C = H^1([0,1];\Omega) \text{ and } e_t : [0,1] \to \Omega)$ $\inf \left\{ \int_C \int_0^1 \frac{1}{2} |\dot{\omega}|^2 dt dQ(\omega) \mid Q \in \mathcal{P}(C), \ (e_0)_{\sharp}Q = \mu, \ (e_1)_{\sharp}Q = \nu \right\}$

• Gives a way to **compare** and **interpolate** between probability measures.

• Consider the optimal solutions for the three formulations $\gamma^{\star}, Q^{\star}, \rho_t^{\star}$ then

$$\pi_t(x,y)_{\sharp}\gamma^{\star} = (e_t)_{\sharp}Q^{\star} = \rho_t^{\star},$$

where $\pi_t(x, y) = (1 - t)x + ty$ and ρ_t is the geodesic between μ and ν , the so called McCann's interpolant.

0 50 100 150 200 250 50 -100 -150 -200 -250 -

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The Schrödinger problem

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 \Rightarrow The Schrödinger problem

Definition (Relative entropy)

Let ρ and π probability measures on Ω then the relative entropy is defined as

$$\mathcal{H}(\rho|\pi) = \begin{cases} \int_{\Omega \times \Omega} \Big(\log \Big(\frac{d\rho(x,y)}{d\pi(x,y)} \Big) - 1 \Big) d\rho(x,y), & \text{if } \rho \ll \pi \\ +\infty, & \text{otherwise} \end{cases}$$

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where R^{ε} is the Wiener measure $R^{\varepsilon} := \int \delta_{x+B^{\varepsilon}} dx$ of variance ε .

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Remark: static formulation can be defined for a general cost function c(x, y).

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Variational Mean Field Games

Lagrangian formulation for 1st order MFG

Consider a first order MFG system then we have the following "equivalence" (see (Lasry and Lions 2007))

A MFG system

$$\begin{split} &-\partial_t \varphi + \frac{1}{2} |\nabla \varphi|^2 = g(x, \rho), \ \varphi(1, x) = \Psi(x) \\ &\partial_t \rho - \operatorname{div}(\rho \nabla \varphi) = 0, \ \rho(0, \cdot) = \rho_0. \end{split}$$

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The (Eulerian) Variational Formulation

$$\inf \int_0^1 \int_\Omega \left(\frac{1}{2} |v_t|^2 \rho_t + G(x, \rho_t) \right) dx dt + F(\rho_1) \quad s.t. \ \partial_t \rho_t + \operatorname{div}(\rho_t v_t) = 0$$
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where G is the anti-derivative of g w.r.t its second variable and $F(\rho_1) = \int_{\Omega} \Psi d\rho_1$ is a final cost.

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The (Lagrangian) Variational Formulation (J.-D. Benamou, G. Carlier, and Santambrogio 2017)

$$\inf_{Q\in\mathcal{P}(C)}\left\{\int_C\int_0^1\frac{1}{2}|\dot{\omega}|^2dtdQ+\int_0^1\mathcal{G}(e_{t,\sharp}Q)dt+F(e_{1,\sharp}Q)\mid(e_0)_\sharp Q=\rho_0\right\},$$

where $\mathcal{G}(\rho) = \int \mathcal{G}(x, \rho) dx$ if $\rho \ll \mathcal{L}$ and $+\infty$ otherwise.

A Lagrangian formulation via Entropy minimization

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Towards a numerical method:

We can solve the Lagrangian problems by firstly discretising them in time as follows

Regularized 1st order MFG

$$\inf \int_{\Omega^{N+1}} \mathcal{K}_N dQ_N(x_0, \cdots, x_N) + \varepsilon \mathcal{H}(Q_N | \mathcal{L}) + \sum_{i=1}^{N-1} \int_{\Omega} G(x, \pi_{i,\sharp} Q_N) dx_i + F(\pi_{N,\sharp} Q_N)$$

s.t. $Q_N \in \mathcal{P}(\Omega^{N+1}), \ \pi_{0,\sharp} Q_N = \rho_0,$

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2nd order MFG

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$$\inf \left\{ \mathcal{H}(Q_N | \mathcal{R}_N^{\varepsilon}) + \sum_{i=1}^{T-1} \int_{\Omega} G(x, \pi_{i,\sharp} Q_N) dx_i + F(\pi_{N,\sharp} Q_N) \mid \pi_{0,\sharp} Q_N = \rho_0 \right\},$$

where $\mathcal{R}_N^{\varepsilon} \stackrel{def}{=} \prod_{n=0}^N \xi_{n,n+1}$ and $\xi_{ij} = \exp^{-\frac{|x_i - x_j|^2}{2N\varepsilon}}.$

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Remarks: (i) for small ε the regularized 1st MFG approximate the unreg pb (ii) both problems can be re-written in the same way.

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2nd order MFG

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$$\inf \left\{ \mathcal{H}(Q_N | R_N^{\varepsilon}) + \sum_{i=1}^{T-1} \int_{\Omega} G(x, \pi_{i,\sharp} Q_N) dx_i + F(\pi_{N,\sharp} Q_N) \mid \pi_{0,\sharp} Q_N = \rho_0 \right\},$$

where $R_N^{\varepsilon} \stackrel{def}{=} \prod_{i=1}^{N} \sum_{j=1}^{N} \xi_{n,\sharp+1}$ and $\xi_{ii} = \exp^{-\frac{|x_i - x_j|^2}{2N\varepsilon}}.$

Remarks: (i) for small ε the regularized 1st MFG approximate the unreg pb (ii) both problems can be re-written in the same way.

IDEA: an alternate coordinate ascent algorithm (equivalent to a generalised Sinkhorn)

The dual problem

The Lagrangian problem can be re-written as the following dual optimization problem:

$$\sup_{\substack{(\phi_{\mathbf{0}},\cdots,\phi_{N})}} -\tilde{F}^{\star}(-\phi_{\mathbf{0}}) - \frac{1}{N} \sum_{k=1}^{N-1} G^{\star}(-\phi_{k}) - F^{\star}(-\phi_{N}) - \int \left(\exp(\bigoplus_{k=0}^{N} \phi_{k}) - 1\right) R_{N}^{\epsilon},$$

where \tilde{F}^* , G^* and F^* are the Legendre transforms of i_{ρ_0} , G and G.

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Proposition ((J.-D. Benamou, G. Carlier, S. Di Marino, and L. Nenna 2018)) Strong duality holds, namely sup = inf. Moreover, denoting by ϕ_k^* and Q_N^* the optimal solutions to the dual and

primal problem respectively, it follows that the unique solution to the primal has the form

$$Q_N^*(x_0,\cdots,x_N) := \left(\bigotimes_{k=0}^N e^{\phi_k^*(x_k)} \right) R_N^{\varepsilon}(x_0,\cdots,x_N). \tag{3}$$

A coordinate ascent algorithm (or generalised Sinkhorn)

Generalizing a result of (Peyré 2015; Chizat, Peyré, B. Schmitzer, and Vialard 2016), we get the iterative method computing a sequence of potentials (denoted with the superscripts $.^{(n)}$) :

$$\begin{split} \phi_{0}^{(n)} &:= \operatorname{argmax}_{\phi} - \tilde{F}^{\star}(-\phi) - \int \exp(\phi) I_{k}^{\phi} dx_{1} \cdots dx_{N}, \\ \phi_{k}^{(n)} &:= \operatorname{argmax}_{\phi} - \frac{1}{N} G^{\star}(-\phi) - \int \exp(\phi) I_{k}^{\phi} dx_{0} \cdots dx_{k-1} dx_{k+1} \cdots dx_{N} \text{ for } k \neq 0, N, \\ \phi_{N}^{(n)} &:= \operatorname{argmax}_{\phi} - F^{\star}(-\phi) - \int \exp(\phi) I_{k}^{u} dx_{0} \cdots dx_{N-1}, \\ \end{split}$$
where

$$I_k^{\phi} := \exp(\oplus_{i=0}^{k-1} \phi_i^{(n)}) \exp(\oplus_{i=k+1}^N \phi_i^{(n-1)}) R^N.$$

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Remarks:

• For many interesting energies *F* and *G*, the relaxed maximizations can be computed point-wise in space and analytically;

• with the same method one can compute dynamic optimal transport by imposing G=0 and $F=i_{
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Sinkhornizing the world!!

- Wasserstein Barycenter (Jean-David Benamou, Guillaume Carlier, Cuturi, Luca Nenna, and Peyré 2015);
- Matching for teams (Luca Nenna 2016);
- Optimal transport with capacity constraint (Jean-David Benamou, Guillaume Carlier, Cuturi, Luca Nenna, and Peyré 2015);
- Partial Optimal Transport (Jean-David Benamou, Guillaume Carlier, Cuturi, Luca Nenna, and Peyré 2015; Chizat, Peyré, B. Schmitzer, and Vialard 2016);
- Multi-Marginal Optimal Transport (Luca Nenna 2016; J.-D. Benamou,
 G. Carlier, and L. Nenna 2016; Jean-David Benamou, Guillaume Carlier, and
 Luca Nenna 2018; Jean-David Benamou, Guillaume Carlier, Cuturi,
 Luca Nenna, and Peyré 2015);
- Wasserstein Gradient Flows (JKO) (Peyré 2015);
- Unbalanced Optimal Transport (Chizat, Peyré, B. Schmitzer, and Vialard 2016);
- Cournot-Nash equilibria (Blanchet, Guillaume Carlier, and Luca Nenna 2017)
- Mean Field Games (J.-D. Benamou, G. Carlier, S. Di Marino, and L. Nenna 2018);
- Grand Canonical Optimal transport (Simone Di Marino, Lewin, and Luca Nenna 2022);
- and more...

- *T* = 32 time steps;
- grid: uniform discretization of $[0,1]^2$ with $N \times N$ points N = 250;
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Thank you!