

Numerical methods and uncertainty quantification for kinetic equations

Lecture 3: Appendix on Asymptotic Preserving schemes On the spread of a virus

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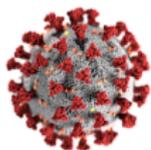
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Outline

- 1 Introduction and motivation
- 2 A multiscale kinetic transport model
- 3 A numerical method capturing the diffusive limit
- 4 Application to the emergence of COVID-19 in Italy

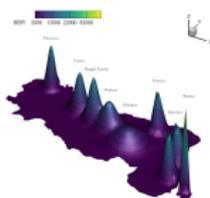
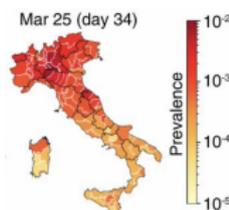
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The ongoing **COVID-19 pandemic** has led to a strong interest from researchers around the world in building and studying new epidemiological models capable of describing the progress of the epidemic¹

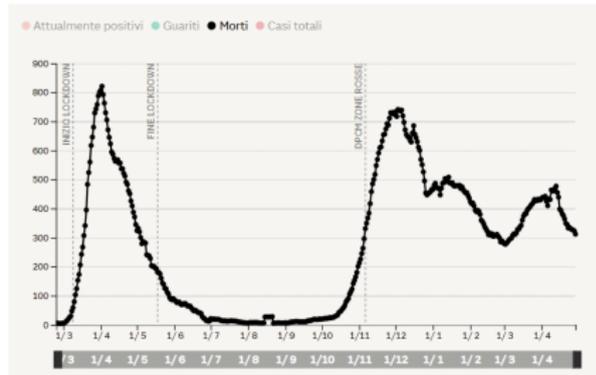
Most classical compartmental models represent the spread of the epidemic only concerning the temporal evolution of the disease among the population, but not taking into account **spatial effects**



Spatial effects can be modeled using **networks** of interacting components (cities, regions, ...) or more in general by considering a fully **two-dimensional** space dynamics

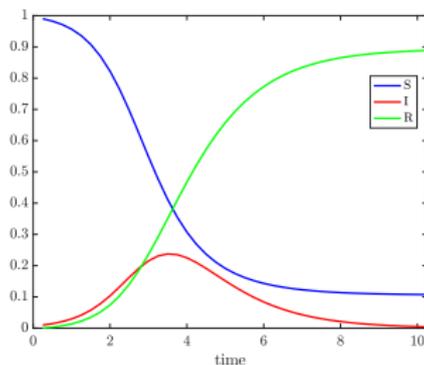
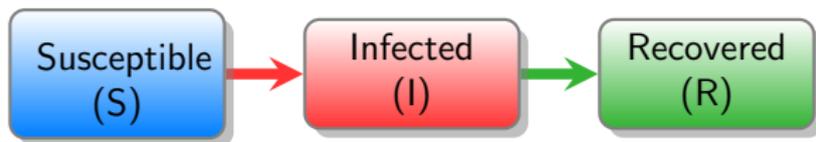
Bellomo et al. '20; Buonomo, Della Marca '20; Colombo et al. '20; Gatto et al. 2020; Giordano et al. '20; Peirlinck et al. '20; Tang et al. 2020; Viguerie et al. '20; Vollmer et al. '20; and many many more....

Additionally, any realistic data-driven model must take into account the large **uncertainty** in the values reported by official sources, such as the amount of infectious individuals



Detected cases (left) and deaths (right) in Italy from the beginning of the pandemic

The SIR model²



$$\begin{aligned}\frac{dS}{dt} &= -\beta \frac{SI}{N} \\ \frac{dI}{dt} &= \beta \frac{SI}{N} - \gamma I \\ \frac{dR}{dt} &= \gamma I\end{aligned}$$

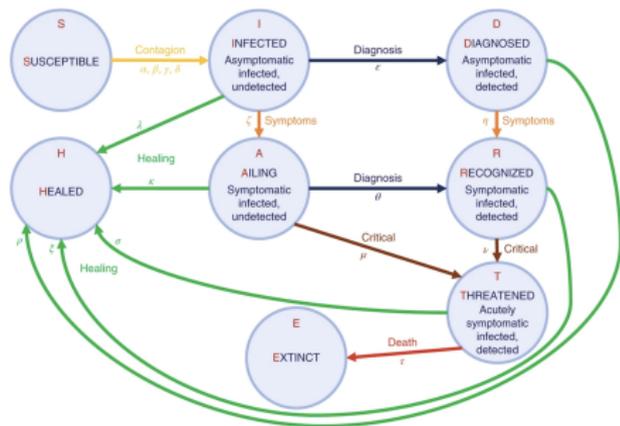
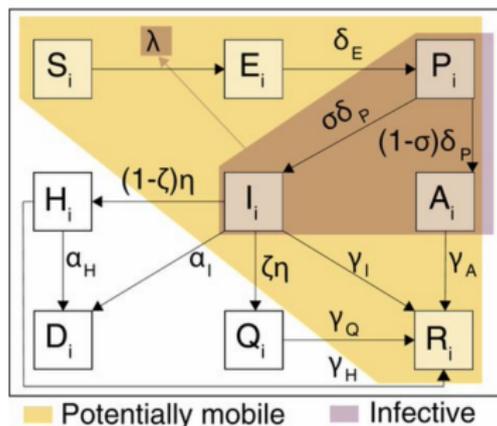
- $N = S + I + R \Rightarrow$ total population (normalized $N = 1$)
- $\beta, \gamma \Rightarrow$ transmission and recovery rates
- $R_0 = \beta/\gamma \Rightarrow$ basic reproduction number

Deterministic model with **no spatial information** on the epidemic spread

²Kermack, McKendrick '27; Hethcote '00

The compartmentalization game

More realistic models involve additional compartmentalizations depending on the specific characteristic of the infectious disease³



Matching models with available data may be a real challenge ([data-driven models](#)). Introducing some degree of uncertainty into the data is an essential feature of analyzing realistic scenarios.

³Hethcote '00; Gatto et al. '20; Giordano et al. '20

The parabolic SIR model⁴

$$\frac{\partial S}{\partial t} = -\beta SI + \frac{\partial}{\partial x} \left(D_S \frac{\partial S}{\partial x} \right)$$

$$\frac{\partial I}{\partial t} = \beta SI - \gamma I + \frac{\partial}{\partial x} \left(D_I \frac{\partial I}{\partial x} \right)$$

$$\frac{\partial R}{\partial t} = \gamma I + \frac{\partial}{\partial x} \left(D_R \frac{\partial R}{\partial x} \right)$$

- $S = S(x, t)$, $I = I(x, t)$, $R = R(x, t)$, $x \in \Omega \subset \mathbb{R}$
- $D_S, D_I, D_R \Rightarrow$ self-diffusion coefficients
- $D_S > D_I$ (population dynamics), $D_I > D_S = 0$ (infection dynamics)
- the diffusion coefficients might also be space-dependent (or nonlinear)

The parabolic character of the model may lead the disease to **propagate instantaneously** over large distances

⁴Webb '86; Murray '01; Berestycki, Roquejoffre, Rossi '21

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A compartmental kinetic transport model for commuters

We consider a population of commuters at position $x \in \Omega$ moving with velocity directions $v \in \mathbb{S}^1$. The kinetic densities of the commuters satisfy the transport equations

$$\begin{aligned}
 \frac{\partial f_S}{\partial t} + \nabla_x \cdot (v_S f_S) &= -F_I(f_S, I_T) - F_A(f_S, A_T) + \frac{1}{\tau_S} (S - f_S) \\
 \frac{\partial f_E}{\partial t} + \nabla_x \cdot (v_E f_E) &= F_I(f_S, I_T) + F_A(f_S, A_T) - a f_E + \frac{1}{\tau_E} (E - f_E) \\
 \frac{\partial f_I}{\partial t} + \nabla_x \cdot (v_I f_I) &= a \sigma f_E - \gamma_I f_I + \frac{1}{\tau_I} (I - f_I) \\
 \frac{\partial f_A}{\partial t} + \nabla_x \cdot (v_A f_A) &= a(1 - \sigma) f_E - \gamma_A f_A + \frac{1}{\tau_A} (A - f_A) \\
 \frac{\partial f_R}{\partial t} + \nabla_x \cdot (v_R f_R) &= \gamma_I f_I + \gamma_A f_A + \frac{1}{\tau_R} (R - f_R)
 \end{aligned} \tag{1}$$

The number of Susceptible, Exposed, Infected and Recovered is

$$\begin{aligned}
 S(x, t) &= \frac{1}{2\pi} \int_{\mathbb{S}^1} f_S(x, v, t) dv, \quad E(x, t) = \frac{1}{2\pi} \int_{\mathbb{S}^1} f_E(x, v, t) dv, \quad R(x, t) = \frac{1}{2\pi} \int_{\mathbb{S}^1} f_R(x, v, t) dv \\
 I(x, t) &= \frac{1}{2\pi} \int_{\mathbb{S}^1} f_I(x, v, t) dv, \quad A(x, t) = \frac{1}{2\pi} \int_{\mathbb{S}^1} f_A(x, v, t) dv,
 \end{aligned}$$

A diffusion compartmental model for non commuters

The unknowns $S_u(x, t)$, $E_u(x, t)$, $I_u(x, t)$, $A_u(x, t)$, $R_u(x, t)$ are the **density fractions of the non-commuters** who, by assumption, move only on an urban scale and satisfy

$$\begin{aligned}
 \frac{\partial S_u}{\partial t} &= -F_I(S_u, I_T) - F_A(S_u, A_T) + \nabla_x \cdot (D_S^u \nabla_x S_u) \\
 \frac{\partial E_u}{\partial t} &= F_I(S_u, I_T) + F_A(S_u, A_T) - aE_u + \nabla_x \cdot (D_E^u \nabla_x E_u) \\
 \frac{\partial I_u}{\partial t} &= a\sigma E_u - \gamma_I I_u + \nabla_x \cdot (D_I^u \nabla_x I_u) \\
 \frac{\partial A_u}{\partial t} &= a(1 - \sigma)E_u - \gamma_A A_u + \nabla_x \cdot (D_A^u \nabla_x A_u) \\
 \frac{\partial R_u}{\partial t} &= \gamma_I I_u + \gamma_A A_u + \nabla_x \cdot (D_R^u \nabla_x R_u).
 \end{aligned} \tag{2}$$

The velocities $v_i = \lambda_i(x)v$ in the kinetic model, the diffusion coefficients $D_i^u = D_i^u(x)$, $i \in \{S, E, I, A, R\}$ and the relaxation times $\tau_i = \tau_i(x)$, $i \in \{S, E, I, A, R\}$ take **into account the heterogeneity of geographical areas**.

The coupled model

The **total densities** are defined by

$$S_T(x, t) = S(x, t) + S_u(x, t), \quad E_T(x, t) = E(x, t) + E_u(x, t), \quad R_T(x, t) = R(x, t) + R_u(x, t),$$

$$I_T(x, t) = I(x, t) + I_u(x, t), \quad A_T(x, t) = A(x, t) + A_u(x, t).$$

The **transmission of the infection** is governed by the incidence functions $F_I(\cdot, I_T)$ and $F_A(\cdot, A_T)$. We assume local interactions to characterize the nonlinear **incidence functions**

$$F_I(g, I_T) = \beta_I \frac{g I_T^p}{1 + \kappa_I I_T}, \quad F_A(g, A_T) = \beta_A \frac{g A_T^p}{1 + \kappa_A A_T},$$

Alternative incidence functions are

$$F_I(g, I_T) = \beta_I \frac{g I_T^p}{1 + \kappa_I \int_{\bar{\Omega}} I_T dx}, \quad F_A(g, A_T) = \beta_A \frac{g A_T^p}{1 + \kappa_A \int_{\bar{\Omega}} A_T dx},$$

$\beta_I = \beta_I(x, t)$ and $\beta_A = \beta_A(x, t)$ characterize the **contact rates** of highly symptomatic and mildly symptomatic/asymptomatic infectious individuals.

$\kappa_I = \kappa_I(x, t)$ and $\kappa_A = \kappa_A(x, t)$ are the **incidence damping coefficients** based on the self-protective behavior.

Commuters behavior in urban areas

A **diffusion behavior** can be recovered when $\tau_{S,I,R} \rightarrow 0$ while the diffusion coefficients

$$D_S = \frac{1}{2} \lambda_{S^2}^2 \tau_S, \quad D_E = \frac{1}{2} \lambda_{E^2}^2 \tau_E, \quad D_I = \frac{1}{2} \lambda_I^2 \tau_I, \quad D_A = \frac{1}{2} \lambda_{A^2}^2 \tau_A, \quad D_R = \frac{1}{2} \lambda_R^2 \tau_R.$$

Let us introduce the **flux functions**

$$J_S = \frac{\lambda_S}{2\pi} \int_{\mathbb{S}^1} v f_S(x, v, t) dv, \quad J_E = \frac{\lambda_E}{2\pi} \int_{\mathbb{S}^1} v f_E(x, v, t) dv, \quad J_I = \frac{\lambda_I}{2\pi} \int_{\mathbb{S}^1} v f_I(x, v, t) dv$$

$$J_A = \frac{\lambda_A}{2\pi} \int_{\mathbb{S}^1} v f_A(x, v, t) dv, \quad J_R = \frac{\lambda_R}{2\pi} \int_{\mathbb{S}^1} v f_R(x, v, t) dv.$$

Integrating system (1) in v we see that the macroscopic densities of commuters obey to

$$\begin{aligned} \frac{\partial S}{\partial t} + \nabla_x \cdot J_S &= -F_I(S, I_T) - F_A(S, A_T) \\ \frac{\partial E}{\partial t} + \nabla_x \cdot J_E &= F_I(S, I_T) + F_A(S, A_T) - aE \\ \frac{\partial I}{\partial t} + \nabla_x \cdot J_I &= a\sigma E - \gamma_I I \\ \frac{\partial A}{\partial t} + \nabla_x \cdot J_A &= a(1 - \sigma)E - \gamma_A A \\ \frac{\partial R}{\partial t} + \nabla_x \cdot J_R &= \gamma_I I + \gamma_A A \end{aligned}$$

Diffusion limit

The flux functions satisfy

$$\begin{aligned} \frac{\partial J_S}{\partial t} + \nabla_x \cdot \left(\frac{\lambda_S^2}{2\pi} \int_{\mathbb{S}^1} ((v \otimes v) f_S) dv \right) &= -F_I(J_S, I_T) - F_A(J_S, A_T) - \frac{1}{\tau_S} J_S \\ \frac{\partial J_E}{\partial t} + \nabla_x \cdot \left(\frac{\lambda_E^2}{2\pi} \int_{\mathbb{S}^1} ((v \otimes v) f_E) dv \right) &= \frac{\lambda_E}{\lambda_S} (F_I(J_S, I_T) + F_A(J_S, A_T)) - aJ_E - \frac{1}{\tau_E} J_E \\ \frac{\partial J_I}{\partial t} + \nabla_x \cdot \left(\frac{\lambda_I^2}{2\pi} \int_{\mathbb{S}^1} ((v \otimes v) f_I) dv \right) &= \frac{\lambda_I}{\lambda_E} a\sigma J_E - \gamma_I J_I - \frac{1}{\tau_I} J_I \\ \frac{\partial J_A}{\partial t} + \nabla_x \cdot \left(\frac{\lambda_A^2}{2\pi} \int_{\mathbb{S}^1} ((v \otimes v) f_A) dv \right) &= \frac{\lambda_A}{\lambda_E} a(1 - \sigma) J_E - \gamma_A J_A - \frac{1}{\tau_A} J_A \\ \frac{\partial J_R}{\partial t} + \nabla_x \cdot \left(\frac{\lambda_R^2}{2\pi} \int_{\mathbb{S}^1} ((v \otimes v) f_R) dv \right) &= \frac{\lambda_R}{\lambda_I} \gamma_I J_I + \frac{\lambda_R}{\lambda_A} \gamma_A J_A - \frac{1}{\tau_R} J_R. \end{aligned}$$

Letting $\tau_{S,I,R} \rightarrow 0$, we get from the r.h.s. in (1)

$$f_S = S, \quad f_E = E, \quad f_I = I, \quad f_A = A, \quad f_R = R,$$

and

$$J_S = -D_S \nabla_x S, \quad J_E = -D_E \nabla_x E, \quad J_I = -D_I \nabla_x I, \quad J_A = -D_A \nabla_x A, \quad J_R = -D_R \nabla_x R,$$

Diffusion limit II

Thus, the following diffusion system for the population of commuters is obtained

$$\begin{aligned}
 \frac{\partial S}{\partial t} &= -F_I(S, I_T) - F_A(S, A_T) + \nabla_x \cdot (D_S \nabla_x S) \\
 \frac{\partial E}{\partial t} &= F_I(S, I_T) + F_A(S, A_T) - aE + \nabla_x \cdot (D_E \nabla_x E) \\
 \frac{\partial I}{\partial t} &= a\sigma E - \gamma_I I + \nabla_x \cdot (D_I \nabla_x I) \\
 \frac{\partial A}{\partial t} &= a(1 - \sigma)E - \gamma_A A + \nabla_x \cdot (D_A \nabla_x A) \\
 \frac{\partial R}{\partial t} &= \gamma_I I + \gamma_A A + \nabla_x \cdot (D_R \nabla_x R)
 \end{aligned} \tag{3}$$

The capability of the model to account for **different regimes**, hyperbolic or parabolic, accordingly to the space dependent relaxation times τ_i , $i \in \{S, E, I, A, R\}$, makes it suitable for describing the dynamics of **human beings**.

Including data uncertainty

- To extend the model to the case of **uncertainties**, let us suppose that the population depend on an additional random vector $\mathbf{z} = (z_1, \dots, z_d)^T \in \Omega_{\mathbf{z}} \subseteq \mathbb{R}^d$.
- Possible **sources** of uncertainty are due to lack of information on the actual number of infected or specific epidemic characteristics.
- This gives the following high-dimensional unknowns

$$f_S(x, v, t, \mathbf{z}), f_E(x, v, t, \mathbf{z}), f_I(x, v, t, \mathbf{z}), f_A(x, v, t, \mathbf{z}), f_R(x, v, t, \mathbf{z}).$$

$$S_u(x, t, \mathbf{z}), E_u(x, t, \mathbf{z}), I_u(x, t, \mathbf{z}), A_u(x, t, \mathbf{z}), R_u(x, t, \mathbf{z}).$$

- Notice that, the **structure of the model** does not change, i.e. there is no direct variation of the unknowns with respect to \mathbf{z} .
- One common choice is to consider the parameters acting inside the incidence function to have a dependence of the form

$$\beta_I = \beta_I(x, t, \mathbf{z}), \beta_A = \beta_A(x, t, \mathbf{z}).$$

$$k_I = k_I(x, t, \mathbf{z}), k_A = k_A(x, t, \mathbf{z}),$$

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The numerical method

- For the **commuters**, the numerical scheme for the deterministic case is based on a discrete ordinate method in velocity in which the **even and odd parity** formulation is employed ⁵.
- Then a **finite volume** method working on two-dimensional **unstructured meshes** ⁶ approximate the discrete ordinate formulation.
- The full discretization of the equations is obtained through the use of suitable **IMEX Runge-Kutta** schemes ⁷.
- The above choices permit to obtain a scheme which consistently **captures the diffusion limit** from the kinetic system when the scaling parameters $\tau_{S,I,R}$ tends toward zero.
- The discretization of the stochastic part is performed by standard **generalized Polynomial Chaos (gPC)** expansion technique ⁸.

⁵Jin-Pareschi '00, Klar '98, D.-Pareschi-Rispoli '14

⁶Boscheri-D. '20

⁷Boscarino-Pareschi-Russo '13, D.-Pareschi-Rispoli '14

⁸Xiu '10

Even and odd parities formulation

We rewrite the commutators system denoting $v = (\eta, \xi) \in \mathbb{S}^1$. This gives

$$r_i^{(1)}(\xi, \eta) = \frac{1}{2}(f_i(\xi, -\eta) + f_i(-\xi, \eta)), \quad r_i^{(2)}(\xi, \eta) = \frac{1}{2}(f_i(\xi, \eta) + f_i(-\xi, -\eta))$$

while for the scalar fluxes one has

$$j_i^{(1)}(\xi, \eta) = \frac{\lambda_i}{2}(f_i(\xi, -\eta) + f_i(-\xi, \eta)), \quad j_i^{(2)}(\xi, \eta) = \frac{\lambda_i}{2}(f_i(\xi, \eta) + f_i(-\xi, -\eta))$$

with $i = S, E, I, A, R$. An equivalent formulation with respect to (1) can then be obtained thanks to this change of variables

$$\frac{\partial r_S^{(1)}}{\partial t} + \xi \frac{\partial j_S^{(1)}}{\partial x} - \eta \frac{\partial j_S^{(1)}}{\partial y} = -F_I(r_S^{(1)}, I_T) - F_A(r_S^{(1)}, A_T) + \frac{1}{\tau_S} (S - r_S^{(1)})$$

$$\frac{\partial r_S^{(2)}}{\partial t} + \xi \frac{\partial j_S^{(2)}}{\partial x} + \eta \frac{\partial j_S^{(2)}}{\partial y} = -F_I(r_S^{(2)}, I_T) - F_A(r_S^{(2)}, A_T) + \frac{1}{\tau_S} (S - r_S^{(2)})$$

$$\frac{\partial r_E^{(1)}}{\partial t} + \xi \frac{\partial j_E^{(1)}}{\partial x} - \eta \frac{\partial j_E^{(1)}}{\partial y} = F_I(r_S^{(1)}, I_T) + F_A(r_S^{(1)}, A_T) - ar_E^{(1)} + \frac{1}{\tau_E} (E - r_E^{(1)})$$

$$\frac{\partial r_E^{(2)}}{\partial t} + \xi \frac{\partial j_E^{(2)}}{\partial x} + \eta \frac{\partial j_E^{(2)}}{\partial y} = F_I(r_S^{(2)}, I_T) + F_A(r_S^{(2)}, A_T) - ar_E^{(2)} + \frac{1}{\tau_E} (E - r_E^{(2)})$$

Even and odd parities formulation II

$$\frac{\partial r_I^{(1)}}{\partial t} + \xi \frac{\partial j_I^{(1)}}{\partial x} - \eta \frac{\partial j_I^{(1)}}{\partial y} = a\sigma r_E^{(1)} - \gamma_I r_I^{(1)} + \frac{1}{\tau_I} (I - r_I^{(1)})$$

$$\frac{\partial r_I^{(2)}}{\partial t} + \xi \frac{\partial j_S^{(2)}}{\partial x} + \eta \frac{\partial j_I^{(2)}}{\partial y} = a\sigma r_E^{(2)} - \gamma_I r_I^{(2)} + \frac{1}{\tau_I} (I - r_I^{(2)})$$

$$\frac{\partial r_A^{(1)}}{\partial t} + \xi \frac{\partial j_A^{(1)}}{\partial x} - \eta \frac{\partial j_A^{(1)}}{\partial y} = a(1 - \sigma)r_E^{(1)} - \gamma_A r_A^{(1)} + \frac{1}{\tau_A} (A - r_A^{(1)})$$

$$\frac{\partial r_A^{(2)}}{\partial t} + \xi \frac{\partial j_A^{(2)}}{\partial x} + \eta \frac{\partial j_A^{(2)}}{\partial y} = a(1 - \sigma)r_E^{(2)} - \gamma_A r_A^{(2)} + \frac{1}{\tau_A} (A - r_A^{(2)})$$

$$\frac{\partial r_R^{(1)}}{\partial t} + \xi \frac{\partial j_R^{(1)}}{\partial x} - \eta \frac{\partial j_R^{(1)}}{\partial y} = \gamma_I r_I^{(1)} + \gamma_A r_A^{(1)} + \frac{1}{\tau_R} (R - r_R^{(1)})$$

$$\frac{\partial r_R^{(2)}}{\partial t} + \xi \frac{\partial j_R^{(2)}}{\partial x} + \eta \frac{\partial j_R^{(2)}}{\partial y} = \gamma_I r_I^{(2)} + \gamma_A r_A^{(2)} + \frac{1}{\tau_R} (R - r_R^{(2)})$$

In addition we have

Even and odd parities formulation III

$$\frac{\partial j_S^{(1)}}{\partial t} + \lambda_S^2 \xi \frac{\partial r_S^{(1)}}{\partial x} - \lambda_S^2 \eta \frac{\partial r_S^{(1)}}{\partial y} = -F_I(j_S^{(1)}, I_T) - F_A(j_S^{(1)}, A_T) - \frac{1}{\tau_S} j_S^{(1)}$$

$$\frac{\partial j_S^{(2)}}{\partial t} + \lambda_S^2 \xi \frac{\partial r_S^{(2)}}{\partial x} + \lambda_S^2 \eta \frac{\partial r_S^{(2)}}{\partial y} = -F_I(j_S^{(2)}, I_T) - F_A(j_S^{(2)}, A_T) - \frac{1}{\tau_S} j_S^{(2)}$$

$$\frac{\partial j_E^{(1)}}{\partial t} + \lambda_E^2 \xi \frac{\partial r_E^{(1)}}{\partial x} - \lambda_E^2 \eta \frac{\partial r_E^{(1)}}{\partial y} = \frac{\lambda_E}{\lambda_S} \left(F_I(j_S^{(1)}, I_T) + F_A(j_S^{(1)}, A_T) \right) - a j_E^{(1)} - \frac{1}{\tau_E} j_E^{(1)}$$

$$\frac{\partial j_E^{(2)}}{\partial t} + \lambda_E^2 \xi \frac{\partial r_E^{(2)}}{\partial x} + \lambda_E^2 \eta \frac{\partial r_E^{(2)}}{\partial y} = \frac{\lambda_E}{\lambda_S} \left(F_I(j_S^{(2)}, I_T) + F_A(j_S^{(2)}, A_T) \right) - a j_E^{(2)} - \frac{1}{\tau_E} j_E^{(2)}$$

$$\frac{\partial j_I^{(1)}}{\partial t} + \lambda_I^2 \xi \frac{\partial r_I^{(1)}}{\partial x} - \lambda_I^2 \eta \frac{\partial r_I^{(1)}}{\partial y} = \frac{\lambda_I}{\lambda_E} a \sigma j_E^{(1)} - \gamma_I j_I^{(1)} - \frac{1}{\tau_I} j_I^{(1)}$$

$$\frac{\partial j_I^{(2)}}{\partial t} + \lambda_I^2 \xi \frac{\partial r_I^{(2)}}{\partial x} + \lambda_I^2 \eta \frac{\partial r_I^{(2)}}{\partial y} = \frac{\lambda_I}{\lambda_E} a \sigma j_E^{(2)} - \gamma_I j_I^{(2)} - \frac{1}{\tau_I} j_I^{(2)}$$

$$\frac{\partial j_A^{(1)}}{\partial t} + \lambda_A^2 \xi \frac{\partial r_A^{(1)}}{\partial x} - \lambda_A^2 \eta \frac{\partial r_A^{(1)}}{\partial y} = \frac{\lambda_A}{\lambda_E} a(1 - \sigma) j_E^{(1)} - \gamma_A j_A^{(1)} - \frac{1}{\tau_A} j_A^{(1)}$$

$$\frac{\partial j_A^{(2)}}{\partial t} + \lambda_A^2 \xi \frac{\partial r_A^{(2)}}{\partial x} + \lambda_A^2 \eta \frac{\partial r_A^{(2)}}{\partial y} = \frac{\lambda_A}{\lambda_E} a(1 - \sigma) j_E^{(2)} - \gamma_A j_A^{(2)} - \frac{1}{\tau_A} j_A^{(2)}$$

$$\frac{\partial j_R^{(1)}}{\partial t} + \lambda_R^2 \xi \frac{\partial r_R^{(1)}}{\partial x} - \lambda_R^2 \eta \frac{\partial r_R^{(1)}}{\partial y} = \frac{\lambda_R}{\lambda_I} \gamma_I j_I^{(1)} + \frac{\lambda_R}{\lambda_A} \gamma_A j_A^{(1)} - \frac{1}{\tau_R} j_R^{(1)}$$

$$\frac{\partial j_R^{(2)}}{\partial t} + \lambda_R^2 \xi \frac{\partial r_R^{(2)}}{\partial x} + \lambda_R^2 \eta \frac{\partial r_R^{(2)}}{\partial y} = \frac{\lambda_R}{\lambda_I} \gamma_I j_I^{(2)} + \frac{\lambda_R}{\lambda_A} \gamma_A j_A^{(2)} - \frac{1}{\tau_R} j_R^{(2)}$$

Space discretization on unstructured grids

- We consider a spatial two-dimensional computational domain Ω discretized by a set of non overlapping polygons $P_i, i = 1, \dots, N_p$.
- Each element P_i exhibits an arbitrary number N_{S_i} of edges e_{ij} . The boundary of the cell is given by $\partial P_i = \bigcup_{j=1}^{N_{S_i}} e_{ij}$.

- The governing equations discretized by means of a finite volume scheme

$$\frac{\partial \mathbf{Q}}{\partial t} + \nabla_x \cdot \mathbf{F}(\mathbf{Q}) = \mathbf{S}(\mathbf{Q}), \quad (x, y) \in \Omega \subset \mathbb{R}^2, \quad t \in \mathbb{R}_0^+,$$

where \mathbf{Q} is the vector of conserved variables

$$\mathbf{Q} = \left(r_i^{(1)}, r_i^{(2)}, j_i^{(1)}, j_i^{(2)} \right)^\top, \quad i = S, E, I, A, R$$

while $\mathbf{F}(\mathbf{Q})$ is the linear flux tensor and $\mathbf{S}(\mathbf{Q})$ represents the stiff source term.

- A first order in time finite volume method is obtained by

$$\mathbf{Q}_i^{n+1} = \mathbf{Q}_i^n - \frac{\Delta t}{|P_i|} \sum_{P_j \in \mathcal{N}_{S_i}} \int_{e_{ij}} \mathbf{F}_{ij}^n \cdot \mathbf{n}_{ij} dl + \int_{P_i} \mathbf{S}_i^n dx.$$

Space discretization on unstructured grids II

- Higher order in space is then achieved by substituting the cell averages by piecewise high order polynomials: $\mathbf{w}_i(\mathbf{x})$.
- These are obtained from the given cell averages relying on a second order Central WENO (CWENO) reconstruction procedure.
- The numerical flux function $\mathbf{F}_{ij} \cdot \mathbf{n}_{ij}$ is a simple and robust local Lax-Friedrichs flux

$$\mathbf{F}_{ij} \cdot \mathbf{n}_{ij} = \frac{1}{2} (\mathbf{F}(\mathbf{w}_{i,j}^+) + \mathbf{F}(\mathbf{w}_{i,j}^-)) \cdot \mathbf{n}_{ij} - \frac{1}{2} s_{\max} (\mathbf{w}_{i,j}^+ - \mathbf{w}_{i,j}^-),$$

where $\mathbf{w}_{i,j}^+, \mathbf{w}_{i,j}^-$ are the high order boundary extrapolated data evaluated through the CWENO reconstruction procedure.

- In the diffusion limit, i.e. as $(\tau_S, \tau_I, \tau_R) \rightarrow 0$, the source term $\mathbf{S}(\mathbf{Q})$ becomes stiff, therefore in order to avoid prohibitive time steps we need to discretize the commuters system partly implicitly.
- A second order IMEX method which preserves the asymptotic limit given by the diffusion equations is proposed.

Time integration

We consider again system (1) formulated using the parities. We assume $\tau_{S,I,R} = \tau$ and rewrite it in partitioned form as

$$\begin{aligned} \frac{\partial \mathbf{u}}{\partial t} + \frac{\partial \mathbf{f}(\mathbf{v})}{\partial x} + \frac{\partial \mathbf{g}(\mathbf{v})}{\partial y} &= \mathbf{E}(\mathbf{u}) + \frac{1}{\tau} (\mathbf{U} - \mathbf{u}) \\ \frac{\partial \mathbf{v}}{\partial t} + \Lambda^2 \frac{\partial \mathbf{f}(\mathbf{u})}{\partial x} + \Lambda^2 \frac{\partial \mathbf{g}(\mathbf{u})}{\partial y} &= \mathbf{E}(\mathbf{v}) - \frac{1}{\tau} \mathbf{v}, \end{aligned}$$

where

$$\mathbf{u} = \left(r_S^{(1)}, r_S^{(2)}, r_E^{(1)}, r_E^{(2)}, r_I^{(1)}, r_I^{(2)}, r_A^{(1)}, r_A^{(2)}, r_R^{(1)}, r_R^{(2)} \right)^T,$$

$$\mathbf{v} = \left(j_S^{(1)}, j_S^{(2)}, j_E^{(1)}, j_E^{(2)}, j_I^{(1)}, j_I^{(2)}, j_A^{(1)}, j_A^{(2)}, j_R^{(1)}, j_R^{(2)} \right)^T,$$

$$\mathbf{f}(\mathbf{v}) = \xi \mathbf{v}, \quad \mathbf{g}(\mathbf{v}) = \eta \mathbf{J} \mathbf{v}, \quad \mathbf{J} = \text{diag}\{-1, 1, -1, 1, -1, 1, -1, 1, -1, 1\},$$

$$\mathbf{E}(\mathbf{u}) = \left(-F_I(r_S^{(i)}, l_T) - F_A(r_S^{(i)}, A_T), F_I(r_S^{(i)}, l_T) + F_A(r_S^{(i)}, A_T) - a r_E^{(i)}, a \sigma r_E^{(i)} - \gamma_I r_I^{(i)}, \right. \\ \left. a(1 - \sigma) r_E^{(i)} - \gamma_A r_A^{(i)}, \gamma_I r_I^{(i)} \gamma_A r_A^{(i)} \right)^T, \quad i = 1, 2$$

$$\mathbf{U} = (S, S, E, E, I, I, A, A, R, R)^T, \quad \Lambda = \text{diag}\{\lambda_S, \lambda_S, \lambda_E, \lambda_E, \lambda_I, \lambda_I, \lambda_A, \lambda_A, \lambda_R, \lambda_R\},$$

and $\mathbf{f}(\mathbf{u})$, $\mathbf{g}(\mathbf{u})$, $\mathbf{E}(\mathbf{v})$ are defined similarly.

Time integration II

The Implicit-Explicit Runge-Kutta (IMEX-RK) approach applied to the partitioned system reads as

$$\mathbf{u}^{(k)} = \mathbf{u}^n - \Delta t \sum_{j=1}^k a_{kj} \left(\frac{\partial \mathbf{f}(\mathbf{v}^{(j)})}{\partial x} + \frac{\partial \mathbf{g}(\mathbf{v}^{(j)})}{\partial y} - \frac{1}{\tau} (\mathbf{U}^{(j)} - \mathbf{u}^{(j)}) \right) + \Delta t \sum_{j=1}^{k-1} \tilde{a}_{kj} \mathbf{E}(\mathbf{u}^{(j)})$$

$$\mathbf{v}^{(k)} = \mathbf{v}^n - \Delta t \sum_{j=1}^{k-1} \tilde{a}_{kj} \left(\Lambda^2 \frac{\partial \mathbf{f}(\mathbf{u}^{(j)})}{\partial x} + \Lambda^2 \frac{\partial \mathbf{g}(\mathbf{u}^{(j)})}{\partial y} - \mathbf{E}(\mathbf{v}^{(j)}) \right) + \Delta t \sum_{j=1}^k a_{kj} \frac{1}{\tau} \mathbf{v}^{(j)},$$

where $\mathbf{u}^{(k)}, \mathbf{v}^{(k)}$ are the so-called internal stages. The numerical solution reads

$$\mathbf{u}^{n+1} = \mathbf{u}^n - \Delta t \sum_{k=1}^s b_k \left(\frac{\partial \mathbf{f}(\mathbf{v}^{(k)})}{\partial x} + \frac{\partial \mathbf{g}(\mathbf{v}^{(k)})}{\partial y} - \frac{1}{\tau} (\mathbf{U}^{(k)} - \mathbf{u}^{(k)}) \right) + \Delta t \sum_{k=1}^s \tilde{b}_k \mathbf{E}(\mathbf{u}^{(k)})$$

$$\mathbf{v}^{n+1} = \mathbf{v}^n - \Delta t \sum_{k=1}^s \tilde{b}_k \left(\Lambda^2 \frac{\partial \mathbf{f}(\mathbf{u}^{(k)})}{\partial x} + \Lambda^2 \frac{\partial \mathbf{g}(\mathbf{u}^{(k)})}{\partial y} - \mathbf{E}(\mathbf{v}^{(k)}) \right) + \Delta t \sum_{k=1}^s b_k \frac{1}{\tau} \mathbf{v}^{(k)}.$$

Furthermore, we choose the Runge-Kutta scheme in such a way that the following relations hold true

$$a_{kj} = b_j, \quad j = 1, \dots, s, \quad \tilde{a}_{kj} = \tilde{b}_j, \quad j = 1, \dots, s-1.$$

Numerical diffusion limit

Assuming for simplicity $D_{S,I,R}$ independent from space, we can write

$$\tau \mathbf{v}^{(k)} = \tau \mathbf{v}^n - \Delta t \sum_{j=1}^{k-1} \tilde{a}_{kj} \left(\tau \Lambda^2 \frac{\partial \mathbf{f}(\mathbf{u}^{(j)})}{\partial x} + \tau \Lambda^2 \frac{\partial \mathbf{g}(\mathbf{u}^{(j)})}{\partial y} - \tau \mathbf{E}(\mathbf{v}^{(j)}) \right) + \Delta t \sum_{j=1}^k a_{kj} \mathbf{v}^{(j)},$$

therefore the limit $\tau \rightarrow 0$ yields

$$\sum_{j=1}^k a_{kj} \mathbf{v}^{(j)} = \sum_{j=1}^{k-1} \tilde{a}_{kj} \left(2\mathbf{D} \frac{\partial \mathbf{f}(\mathbf{U}^{(j)})}{\partial x} + 2\mathbf{D} \frac{\partial \mathbf{g}(\mathbf{U}^{(j)})}{\partial y} \right),$$

where we used $\mathbf{u}^{(j)} = \mathbf{U}^{(j)}$ as $\tau \rightarrow 0$. Using now the same identity $\mathbf{u}^{(j)} = \mathbf{U}^{(j)}$ we get

$$\mathbf{U}^{(k)} = \mathbf{U}^n - \Delta t \sum_{j=1}^k a_{kj} \left(\frac{\partial \mathbf{f}(\mathbf{v}^{(j)})}{\partial x} + \frac{\partial \mathbf{g}(\mathbf{v}^{(j)})}{\partial y} \right) + \Delta t \sum_{j=1}^{k-1} \tilde{a}_{kj} \mathbf{E}(\mathbf{U}^{(j)}),$$

Numerical diffusion limit II

Thanks to the definitions of \mathbf{f} and \mathbf{g} we get

$$\begin{aligned} \mathbf{U}^{(k)} = & \mathbf{U}^n - 2\Delta t \mathbf{D} \sum_{j=1}^{k-1} \tilde{a}_{kj} \left(\xi^2 \frac{\partial^2 \mathbf{U}^{(j)}}{\partial x^2} + 2\xi \eta \mathbf{J} \frac{\partial^2 \mathbf{U}^{(j)}}{\partial x \partial y} + \eta^2 \frac{\partial^2 \mathbf{U}^{(j)}}{\partial y^2} \right) \\ & + \Delta t \sum_{j=1}^{k-1} \tilde{a}_{kj} \mathbf{E} \left(\mathbf{U}^{(j)} \right). \end{aligned}$$

Finally, integrating over the velocity field one has

$$\begin{aligned} S^{(k)} = & S^n - \Delta t D_S \sum_{j=1}^{k-1} \tilde{a}_{kj} \left(\frac{\partial^2 S^{(j)}}{\partial x^2} + \frac{\partial^2 S^{(j)}}{\partial y^2} \right) - \Delta t \sum_{j=1}^{k-1} \tilde{a}_{kj} F_I(S^{(j)}, I_T^{(j)}) + F_A(S^{(j)}, A_T^{(j)}), \\ R^{(k)} = & R^n - \Delta t D_R \sum_{j=1}^{k-1} \tilde{a}_{kj} \left(\frac{\partial^2 R^{(j)}}{\partial x^2} + \frac{\partial^2 R^{(j)}}{\partial y^2} \right) + \Delta t \sum_{j=1}^{k-1} \tilde{a}_{kj} \gamma_I I^{(j)} + \gamma_A A^{(j)} \end{aligned}$$

and the same for $A^{(k)}$, $I^{(k)}$ and $E^{(k)}$.

Thus, the **internal stages correspond to the stages of the explicit scheme applied to the reaction-diffusion system**. Moreover, the last stage is equivalent to the numerical solution.

Stochastic collocation method

- We restrict to the case in which there is only **one stochastic variable** z in the system.
- The probability density function is supposed known. The approximate solution for the commuters $\mathbf{Q}_M(x, v, t, z)$ and the non commuters $\mathbf{Q}_M^u(x, t, z)$ are represented as **truncations of the series of the orthonormal polynomials**, i.e.

$$\mathbf{Q}_M(x, v, t, z) = \sum_{j=1}^M \hat{\mathbf{Q}}_j(x, v, t) \phi_j(z), \quad \mathbf{Q}_M^u(x, t, z) = \sum_{j=1}^M \hat{\mathbf{Q}}_j^u(x, t) \phi_j(z)$$

where M is the number of terms of the truncated series and $\phi_j(z)$ are orthonormal polynomials, with respect to the measure $\rho_z(z) dz$.

- The **expansion coefficients** are obtained by

$$\hat{\mathbf{Q}}_j(x, v, t) = \int_{\Gamma} \mathbf{Q}(x, v, t, z) \phi_j(z) \rho_z(z) dz, \quad j = 1, \dots, M.$$

$$\hat{\mathbf{Q}}_j^u(x, t) = \int_{\Gamma} \mathbf{Q}^u(x, t, z) \phi_j(z) \rho_z(z) dz, \quad j = 1, \dots, M.$$

Stochastic collocation method II

- The exact integrals for the expansion coefficients are replaced by a **suitable quadrature formula** characterized by the set $\{z_n, w_n\}_{n=1}^{N_p}$, z_n the collocation point, w_n the corresponding weight, N_p the number of quadrature points

$$\hat{\mathbf{Q}}_j(x, v, t) \approx \sum_{n=1}^{N_p} \mathbf{Q}(x, v, t, z_n) \phi_j(z_n) w_n, \quad \hat{\mathbf{Q}}_j^u(x, t) \approx \sum_{n=1}^{N_p} \mathbf{Q}^u(x, t, z_n) \phi_j(z_n) w_n,$$

where $\mathbf{Q}(x, v, t, z_n)$ and $\mathbf{Q}^u(x, t, z_n)$ with $n = 1, \dots, N_p$ are the solutions of the problem evaluated at the n -th collocation point.

- All quantities of interest** concerning the random variable can then be computed. For example, the expectations are approximated as

$$\mathbb{E}[\mathbf{Q}] \approx \mathbb{E}[\mathbf{Q}_M] = \int_{\Gamma} \mathbf{Q}_M(x, v, t, z) \rho_z(z) dz \approx \sum_{n=1}^{N_p} \mathbf{Q}(x, v, t, z_n) w_n,$$

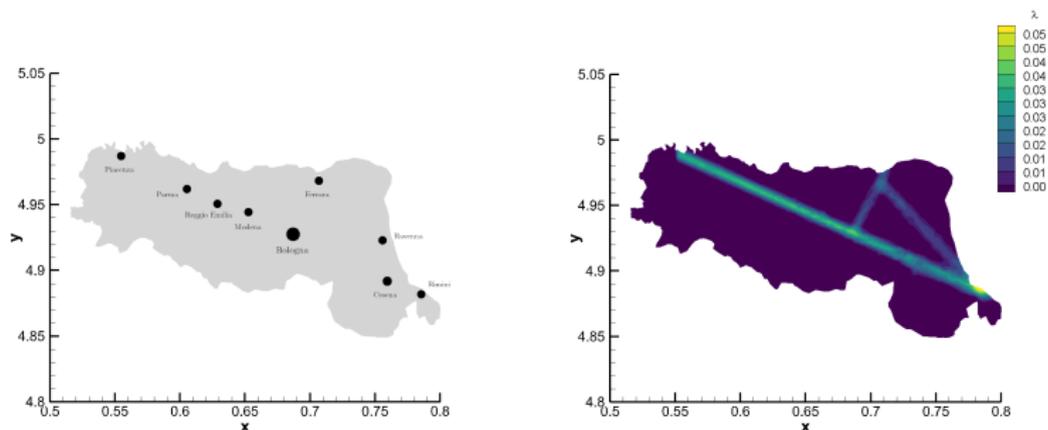
$$\mathbb{E}[\mathbf{Q}^u] \approx \mathbb{E}[\mathbf{Q}_M^u] = \int_{\Gamma} \mathbf{Q}_M^u(x, t, z) \rho_z(z) dz \approx \sum_{n=1}^{N_p} \mathbf{Q}^u(x, t, z_n) w_n.$$

Outline

- 1 Introduction and motivation
- 2 A multiscale kinetic transport model
- 3 A numerical method capturing the diffusive limit
- 4 Application to the emergence of COVID-19 in Italy

Application to the spread of COVID-19 in a realistic geographical scenario

COVID-19 outbreak in Emilia-Romagna (Italy), 1-10 March 2020



From \ To	PC	PR	RE	MO	BO	FE	RA	FC	RN	C [%]
PC	-	4178	-	-	-	-	-	-	-	1.45
PR	1707	-	5142	-	-	-	-	-	-	1.51
RE	-	8969	-	19841	-	-	-	-	-	5.42
MO	-	-	11488	-	13034	1173	-	-	-	3.63
BO	-	-	-	6842	-	5983	3887	-	-	1.64
FE	-	-	-	2682	16865	-	2610	-	-	6.80
RA	-	-	-	-	9808	1016	-	9211	-	5.14
FC	-	-	-	-	-	-	6646	-	6944	3.41
RN	-	-	-	-	-	-	-	6075	-	1.79

Population

Initial distribution of a generic population $f(x, y)$ assigned to each main city c

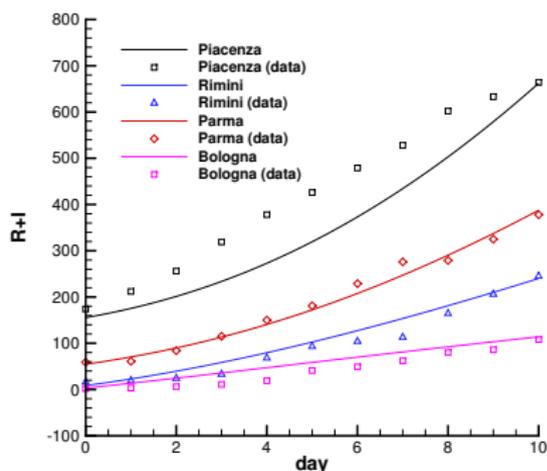
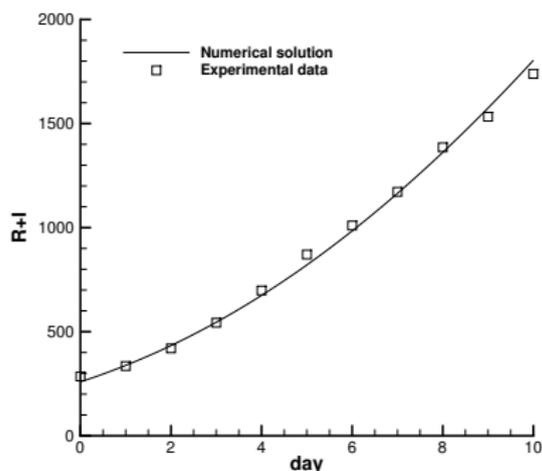
$$f(x, y) = \frac{1}{2\pi r_c} e^{-\frac{(x-x_c)^2 + (y-y_c)^2}{2r_c^2}} f_c$$

where r_c is the radius of the urban area.

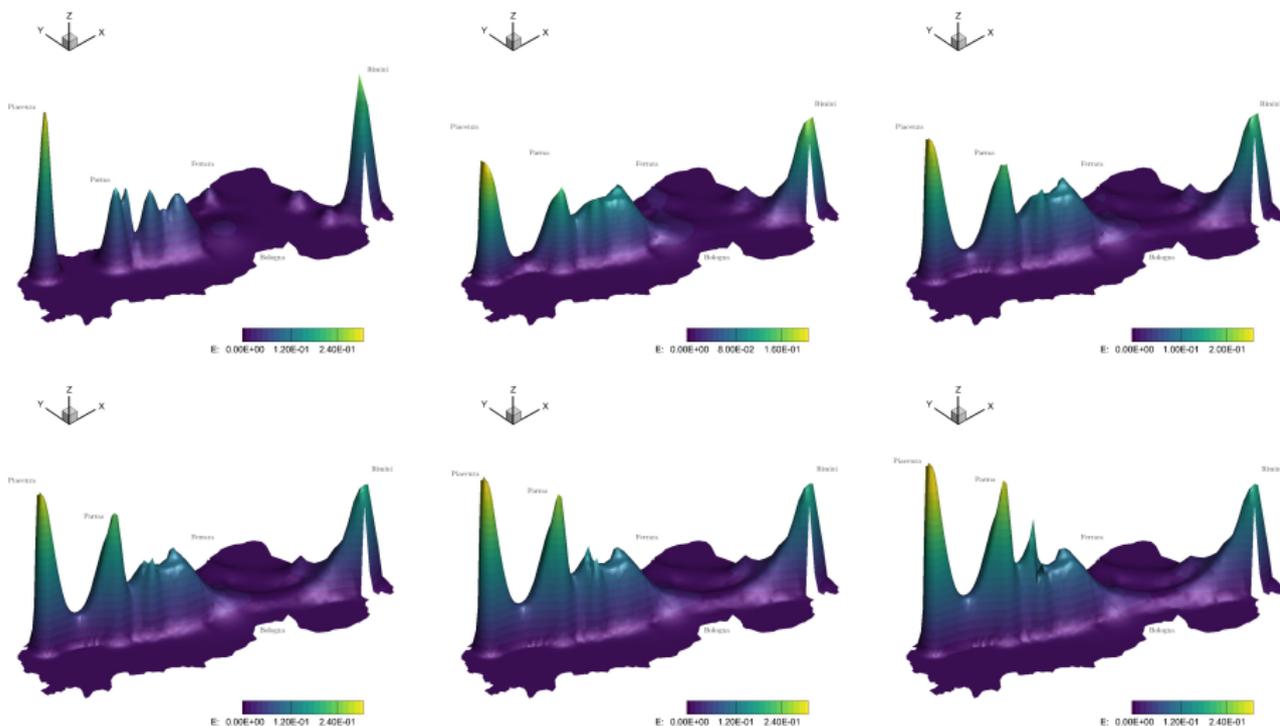
Model calibration

Using the connection matrix and the road distances for the traveling speed λ_i ,
 $i \in \{S, E, I, R\}$

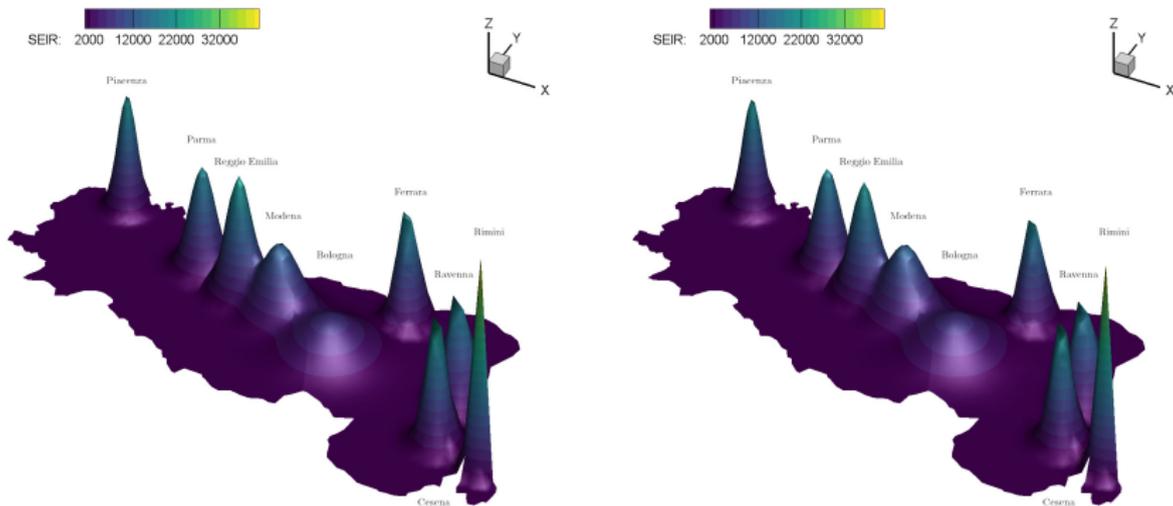
Using the corresponding SEIR ODE model for the epidemic parameters



Time evolution of total infected and recovered population ($R + I$) compared against **experimental data** for the region Emilia-Romagna. Right: time evolution of total infected and recovered population ($R + I$) compared against experimental data.



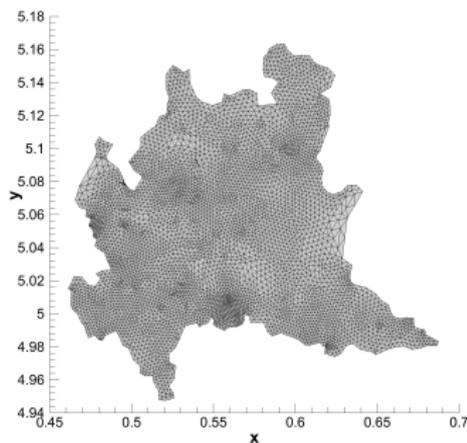
Distribution of exposed population E including asymptomatic at times $t = 0$, $t = 2$, $t = 4$ (top), $t = 6$, $t = 8$ and $t = 10$ (bottom).



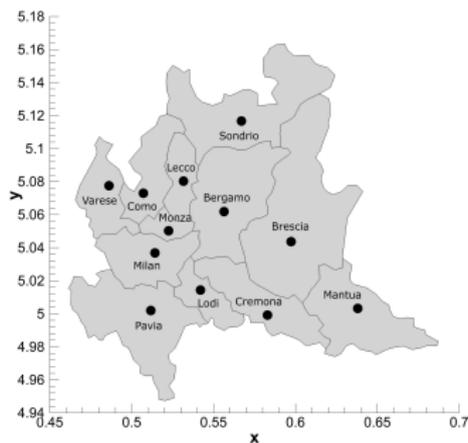
Distribution of total population $(S + S_u) + (E + E_u) + (I + I_u) + (R + R_u)$ at initial time $t = 0$ (left) and at the final time $t = 10$ (right). The model bring back at regular intervals from day to day the situation to the initial condition in which all commuters and non-commuters are located in the main urban areas.

The Lombardy network

From/To	LO	MI	BG	BS	CR	Total commuters
LO	-	56717 (80.53%)	-	-	13712 (19.47%)	70429 (30.97%)
MI	55397 (31.14%)	-	74168 (41.69%)	26709 (15.01%)	21622 (12.15%)	177896 (5.45%)
BG	-	76337 (45.79%)	-	78348 (47.00%)	12016 (7.21%)	166701 (15.04%)
BS	-	26594 (23.24%)	70879 (61.93%)	-	16967 (14.83%)	114440 (9.12%)
CR	13264 (20.06%)	23142 (35.00%)	12025 (18.19%)	17681 (26.75%)	-	66112 (18.58%)

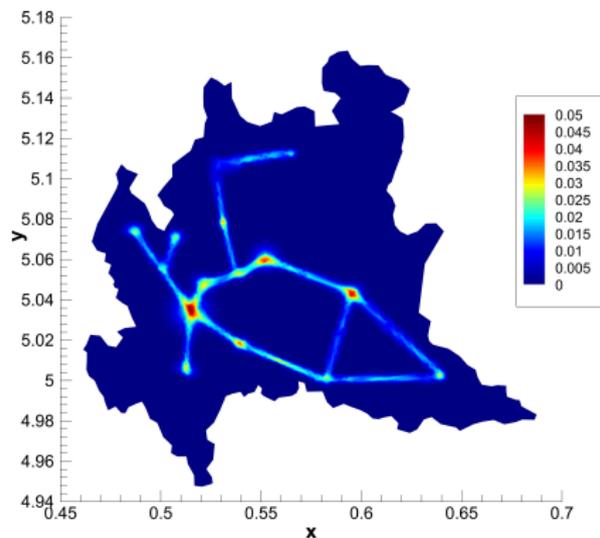


mesh grid

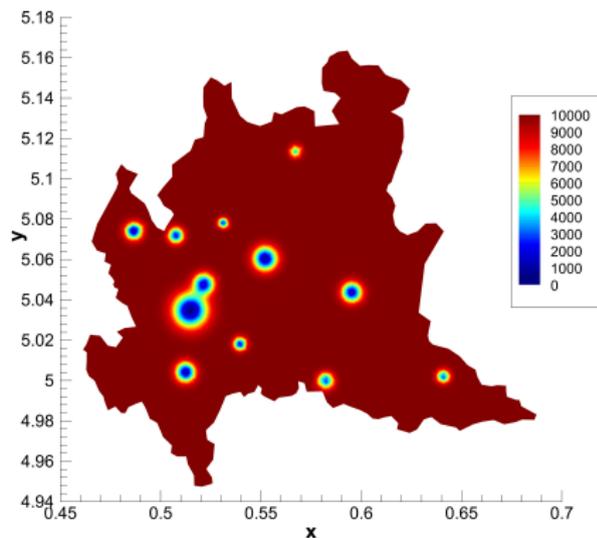


provinces

The Lombardy network II



$$\lambda_i, i \in \{S, E, A, R\}$$



$$\tau_i, i \in \{S, E, I, A, R\}$$

Figure: Left initial condition imposed for characteristic speeds λ . Right relaxation times τ .

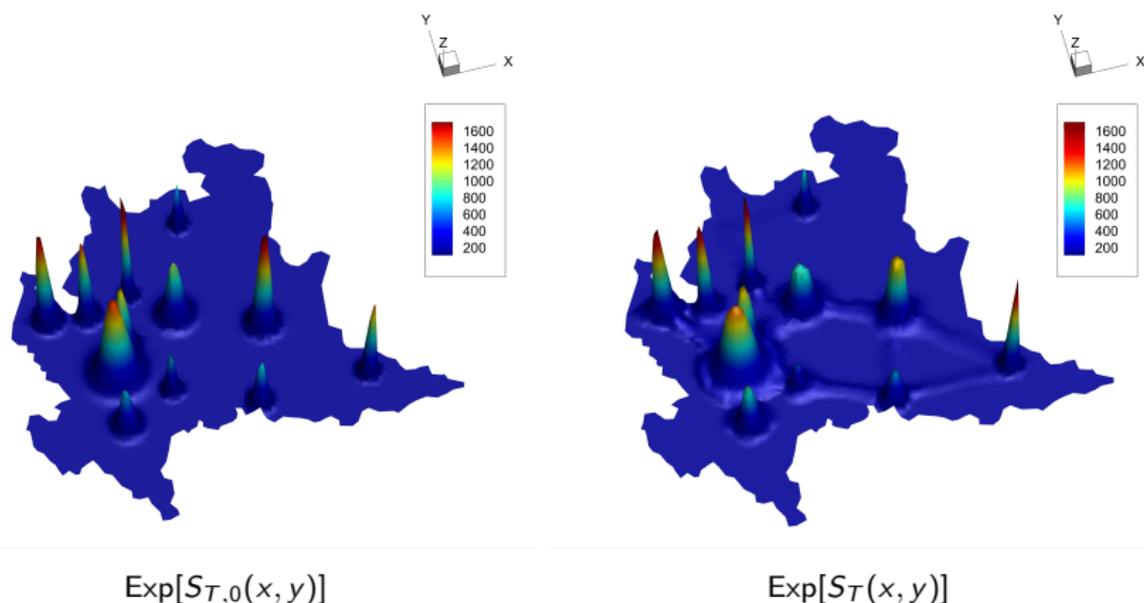
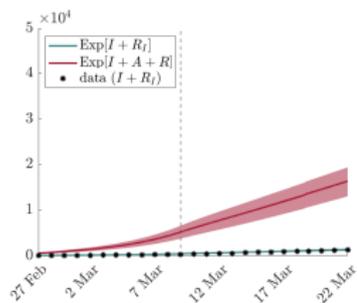
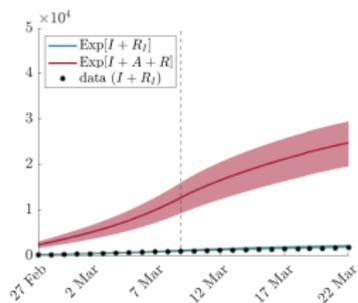


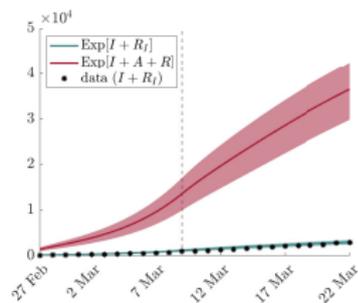
Figure: Numerical results of the simulation of the first outbreak of COVID-19 in Lombardy, Italy. Left expectation of the susceptible population S_T on the initial day simulated (February 27, 2020) and right at the end of the simulation (March 22, 2020).



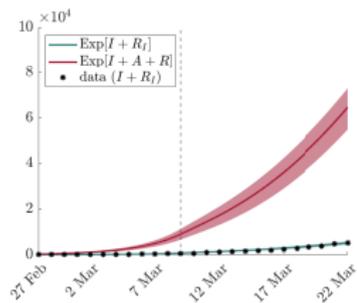
Pavia



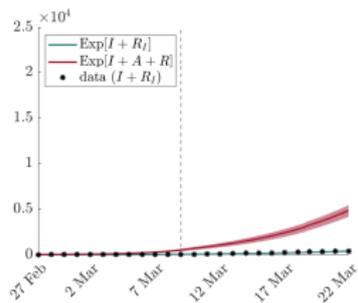
Lodi



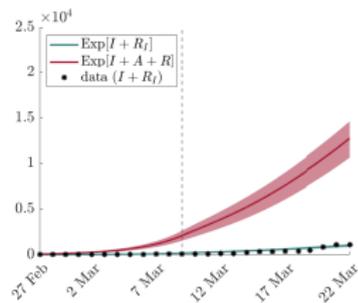
Cremona



Milan



Varese



Monza-Brianza

Figure: Numerical results, with 95% confidence intervals, of the simulation of the first outbreak of COVID-19 in Lombardy, Italy. Data are taken from the COVID-19 repository of the Civil Protection Department of Italy. Vertical dashed lines identify the onset of governmental lockdown restrictions.

Research perspectives

Conclusions

- **Multiscale** kinetic models for the description of the spread of an epidemic disease in a spatially heterogeneous context
- Introduction of **commuters** and non-commuters population.
- Application to a **real epidemic spread** (COVID-19 outbreak in Emilia-Romagna and Lombardy).
- Second order asymptotic preserving time discretization and finite volume spatial discretization made of **unstructured meshes** of arbitrary shape.

Outlook

- Multidimensional **uncertainty quantification**.
- **Data fitting** for extrapolation of the model parameter and construction of data driven models.
- Incorporation of an **age-structured** population.
- **Control**.

The material for this part is mostly based on the recent survey: "**Kinetic modelling of epidemic dynamics: social contacts, control with uncertain data, and multiscale spatial dynamics.**" G. Albi, G. Bertaglia, W. Boscheri, G.D., L. Pareschi, G. Toscani, M. Zanella. Predicting Pandemics in a Globally Connected World, Vol. 1, Birkhauser-Springer Series: Modeling and Simulation in Science, Engineering and Technology, 2022.