

Subsurface flow with uncertainty : applications and numerical analysis issues

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1 Subsurface flow with uncertainty : a simple model

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What is hydrogeology ?

“The word “hydrogeology” can be understood as a combination of “hydraulics” and “geology”...

“Hydrogeology” is thus the science where the two are combined : finding the solution of the flow (and transport) equations in a complex, only partly identified, geological system“.

Dealing with spatial heterogeneity,
Marsily *et al.*,
Hydrogeology Journal, 2005

What is hydrogeology ?

- Hydraulics : well-known PDEs (Darcy, transport,...)
- Geological properties : lack of data (few well tests/pumping tests) and more importantly multiple scales of Heterogeneity (from pore to aquifer)

How to deal with heterogeneity in hydrogeology ?

Early approach : **average** to get equivalent homogeneous properties

- The **experiments themselves** provide an average (Darcy 1856 until \simeq 1960)
- Quantitative hydrogeology and numerical modelling (from \simeq 1960) :
 - ▶ A few pumping test values (wells) provide **local averaged permeability values** which are **interpolated** in the aquifer.
 - ▶ Matheron 1967 : aquifer properties are described as **random variables** (for the purpose of averaging and not to describe heterogeneity).
Lognormal laws are used to describe the permeability.
- This approach works quite well to compute pressure
- It is limited for predictions of oil recovery, groundwater contamination problems
 \rightsquigarrow such problems are very sensitive to high-permeability channels, faults, low permeability barriers
- pressure variations due to heterogeneity are small, whereas those of velocities and travel times are large

Stochastic modeling to describe heterogeneity

- Geostatistics (Delhomme 1976...) and stochastic hydrogeology (Gelhar 1976, Dagan 1985...) use **random variables** (lognormal laws) taking into account **spatial covariance** to describe heterogeneity.
- Heterogeneity can be described by a **“structure” defined by the spatial covariance** (Marsily 1986, Chiles 1999...)
- Calibration of the parameters (**inverse problems**) and provide **Monte-Carlo simulations** of aquifer models to **estimate uncertainty** on the flow and transport (Delhomme 1979, Ramaro et al. 1995, Zimmermann et al. 1998...)
- “Geostatistics make better use of the data without asking for more” : few degrees of freedom (variance range and type of covariance)

Groundwater steady flow

- We consider an isotropic porous medium with constant porosity and permeability \mathbf{a} .
- We denote by u the hydraulic head and by \mathbf{v} the Darcy velocity.
- **Darcy law** $\mathbf{v} = -\mathbf{a}\nabla p$ together with **mass conservation** $\operatorname{div} \mathbf{v} = 0$ yield the steady flow equation:

$$\operatorname{div}(\mathbf{a}(x)\nabla u(x)) = 0 \quad \forall x \in D \subset \mathbb{R}^d$$

+boundary conditions.

\rightsquigarrow in practice : typically we take a box with mixed homogeneous Neumann conditions (up/down) and non-homogeneous Dirichlet conditions (left/right).

Uncertainty in groundwater steady flow

- Because of :
 - ▶ the heterogeneity of the permeability a
 - ▶ the lack of data
- ↪ we use a stochastic model : the permeability a is not known exactly instead we suppose that we know its law : the permeability is then a random field $a(\omega, x)$
- More precisely a is a function $a : \Omega \times D \rightarrow \mathbb{R}$. We can see a as a random variable taking values into $\mathcal{C}^0(\bar{D})$
- the steady flow equation with uncertainty is then:

$$\operatorname{div}_x(a(\omega, x)\nabla_x u(\omega, x)) = 0 \quad \forall x \in D \subset \mathbb{R}^d, \omega \text{ a.e.}$$

+boundary conditions.

Equation and first assumptions

- Let D be a bounded open C^2 domain of \mathbb{R}^d , (Ω, \mathcal{F}, P) a probability space and $f \in L^2(D)$.
- We look for $u : \Omega \times D \rightarrow \mathbb{R}$ such that for almost every ω

$$\begin{cases} -\operatorname{div}_x(a(\omega, x)\nabla u_x(\omega, x)) & = f(x) & x \in D \\ u(\omega, \cdot) & = 0 & \text{on } \partial D. \end{cases}$$

- Remark : to ensure that the equation is well posed we will require that we have for almost all $\omega : a(\omega, \cdot) \in L^\infty(D)$

$$0 < a_{\min}(\omega) \leq a(\omega, x) < a_{\max}(\omega) < +\infty,$$

for almost every x .

A model for the law of the permeability : lognormal law

- A widely used model in hydrogeology : $a : \Omega \times \bar{D} \rightarrow \mathbb{R}$ is a **lognormal homogeneous** random field, which means that : $a(\omega, x) = e^{g(\omega, x)}$, where g is a **gaussian random field**,
 \rightsquigarrow i.e. any linear combination $\lambda_1 g(x_1, \omega) + \dots + \lambda_n g(x_n, \omega)$ is a gaussian random variable.
- The law of g is determined by its **expected value** $x \mapsto \mathbb{E}[g(\omega, x)]$ and its **covariance function** :
$$\text{cov}[g](x, y) = \mathbb{E}[(g(\omega, x) - \mathbb{E}[g(\omega, x)])(g(\omega, y) - \mathbb{E}[g(\omega, y)])]$$
- we suppose g to be **homogeneous**,
i.e. $\text{cov}[g]$ only depends on $\|x - y\|$.
 \rightsquigarrow it means that the law of the permeability field is invariant by any affine isometry of the spatial domain.
- Such a choice enables to modelize such **very heterogeneous fields**

Typical examples of covariance functions (Mattern class)

- the widely use case of an **exponential covariance** (Hoeksema et al. 1985, Gelhar 1986...) : $\text{cov}[g](x, y) = \sigma^2 e^{-\frac{\|x-y\|}{\lambda}}$.
- the case of a **gaussian covariance** $\text{cov}[g](x, y) = \sigma^2 e^{-\left(\frac{\|x-y\|}{\lambda}\right)^2}$.

↪ two examples leading to **very different mathematical and numerical properties**.

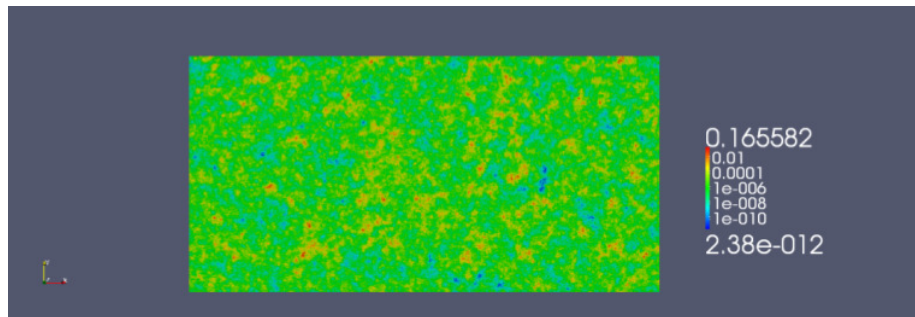


Figure : case of an exponential covariance with $\lambda = 10$, $\sigma = 3$, $\mathbb{E}[g] = -14$, realized by Géraldine Pichot with the software platform H2OLAB

Domains of application

- Study of groundwater pollution (management of groundwater resources)
- Oil and gas recovery
- Storage of nuclear waste
- Geological sequestration of carbon dioxide
- ...

Various quantities of interest

- The law of $u : \Omega \rightarrow H^1(D)$:
the law of u is determined by the knowledge of all the values of the $\mathbb{E}[\varphi(u)]$ for a certain class of functions φ
- In practice it is not possible (and not really interesting) to know the law completely. We consider only particular test functions φ .
 - ▶ mean point values of the pressure $\mathbb{E}[u(x)]$
 - ▶ variance of point values of the pressure $\mathbb{E}[(u(x) - \mathbb{E}[u(x)])^2]$,
 - ▶ mean value of some norm of the pressure (typically $\mathbb{E}[\|u\|_{L^2(D)}], \mathbb{E}[\|u\|_{L^\infty(D)}], \dots$),
 - ▶ mean value of outflow through a part Γ of the boundary $\mathbb{E} \left[\int_{\Gamma} -a(\omega, x) \nabla u(\omega, x) \cdot d\vec{\nu} \right]$
- exit times of transported particles (with or without diffusion).
- Failure probabilities and cdf.
 - ▶ pressure at a fault $\mathbb{P}(p(x) \geq c)$
 - ▶ exit time $\mathbb{P}(T_{exit} \leq t_{critical})$
- Densities of some functionals of the solution.
- Rare events (failure probabilities, rare events).
- Inverse problems.

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Spatial regularity of lognormal fields

Since g is homogeneous we can write $\text{cov}[g](x, y) = k(x - y)$.

We suppose in what follows that $k \in \mathcal{C}^{0,1}(\mathbb{R}^d, \mathbb{R})$. We deduce the following **spatial regularity result** :

Proposition

There exists a version of a , still denoted a , such that for almost all ω , $a(\omega, \cdot) \in \mathcal{C}^{0,\beta}(D)$ for any $\beta < \frac{1}{2}$.

Proof : Since k a L -lipschitz continuous function, we deduce that

$$\begin{aligned}\mathbb{E}[|g(x) - g(y)|^2] &= \mathbb{E}[g(x)^2] - 2\mathbb{E}[g(x)g(y)] + \mathbb{E}[g(y)^2] \\ &= 2(k(0) - k(x - y)) \leq 2L\|x - y\|.\end{aligned}$$

$g(x) - g(y)$ is a mean-free gaussian random variable, so for any positive integer p ,

$$\mathbb{E}[|g(x) - g(y)|^{2p}] \leq c_p(2L)^p\|x - y\|^p.$$

Proof of the spatial regularity result

We use the Kolmogorov's continuity theorem

Theorem (Kolmogorov)

Let $X(\omega, x) : \Omega \times D \subset \mathbb{R}^d \rightarrow \mathbb{R}^n$ be a stochastic process such that there exists constants $C, p > 1$ and $\varepsilon > 0$ such that for any $x, y \in D$ we have

$$\mathbb{E}[\|X(\omega, x) - X(\omega, y)\|^p] \leq C \|x - y\|^{d+\varepsilon},$$

then X admits a version \tilde{X} such that for almost all ω , $\tilde{X}(\omega, \cdot) \in \mathcal{C}^{0,\beta}(\bar{D})$ for any $\beta < \frac{\varepsilon}{p}$.

and deduce that here exists a version of g which is a.s. Hölder-continuous with any exponent $\beta < \frac{p-d}{2p}$, it remains to let $p \rightarrow +\infty$.

More spatial regularity

Using the same arguments we can prove that if the **covariance function is more regular** (more precisely if the function k is more regular), the realizations of a will also have more spatial regularity.

Proposition

If $k \in \mathcal{C}^{n,\alpha}(\mathbb{R})$ (with $\alpha > 0$) then there exists a version of a , still denoted a such that for almost all ω , $a(\omega, \cdot) \in \mathcal{C}^{[\beta],\beta-[\beta]}(D)$ for any $\beta < \frac{n+\alpha}{2}$.

Coming back to our two examples :

- **exponential covariance** : $\text{cov}[g](x, y) = \sigma^2 e^{-\frac{\|x-y\|}{\ell}}$.
 $\text{cov}[g]$ has only lipschitz regularity (as a function of $x - y$)
 \Rightarrow the realizations of a are only Hölder continuous on \bar{D} with any exponent $\beta < 1/2$.
- **gaussian covariance** : $\text{cov}[g](x, y) = \sigma^2 e^{-\left(\frac{\|x-y\|}{\ell}\right)^2}$.
 $\text{cov}[g]$ has \mathcal{C}^∞ regularity (as a function of $x - y$)
 \Rightarrow the realizations of a are \mathcal{C}^∞ on \bar{D} .

Lognormal fields are not uniformly bounded from above or below

- a is neither **uniformly bounded from above nor below** with respect to ω . In other words, $a(\omega, x)$ can be **arbitrary close to 0** and **arbitrary close to $+\infty$** .
- However, since we have seen that for almost every ω , $x \mapsto a(\omega, x)$ is **continuous on \bar{D}** , we can then define for almost all ω :
$$a_{min}(\omega) = \min_{x \in \bar{D}} a(\omega, x) > 0 \text{ and } a_{max}(\omega) = \max_{x \in \bar{D}} a(\omega, x) < +\infty.$$
- Conclusion : $\frac{1}{a_{min}(\omega)} \notin L^\infty(\Omega)$ and $a_{max}(\omega) \notin L^\infty(\Omega)$.

An integrability result

Proposition

$\frac{1}{a_{\min}(\omega)} \in L^p(\Omega)$ and $a_{\max}(\omega) \in L^p(\Omega) \forall p > 0$. Moreover we have $a \in L^p(\Omega, \mathcal{C}_0^{0,\beta}(\bar{D})) \forall p > 0$ and $\beta < 1/2$.

Proof :

$g : \Omega \rightarrow \mathcal{C}^0(\bar{D})/\mathcal{C}_0^{0,\beta}(\bar{D})$ is a gaussian random variable which takes values into a separable Banach space. We can apply Fernique's theorem

Theorem (Fernique)

If E is a separable Banach space and X a mean-free gaussian random variable with values in E , for any finite $p > 0$, we have $\mathbb{E}[e^{p\|X\|_E}] < \infty$.

An important consequence : an integrability property of the solution

Proposition

The equation

$$\begin{cases} -\operatorname{div}(a(\omega, \cdot) \nabla u(\omega, \cdot)) = f(x) & \text{on } D \\ u(\omega, \cdot) = 0 & \text{on } \partial D. \end{cases}$$

admits a unique solution $u \in L^p(\Omega, H_0^1(D))$, $\forall p > 0$.

Proof

For a.e. ω , the equation admits a unique solution $u(\omega, \cdot) \in H_0^1(D)$ with

$$\|u(\omega, \cdot)\|_{H_0^1(D)} \leq \frac{C_D}{a_{\min}(\omega)} \|f\|_{L^2(D)}.$$

Therefore

$$\|u\|_{L^p(\Omega, H_0^1(D))} \leq C_D \|f\|_{L^2(D)} \left\| \frac{1}{a_{\min}} \right\|_{L^p(\Omega)}.$$

Approximation of a

We approximate the random field $a(\omega, x)$ by a function of x and of N random variables, i.e. in a finite dimensional stochastic space:

$$a(\omega, x) \rightsquigarrow \tilde{a}(Y_1(\omega), \dots, Y_N(\omega), x).$$

Why approximate a ?

- It can be used to simulate a : we need only to simulate N random variables to simulate \tilde{a} .
- It is the first and fundamental step of several numerical methods: in particular stochastic collocation methods.

Approximation of a using a Karhunen Loève expansion

- We consider the Hilbert-Schmidt operator:

$$f \in L^2(D) \longmapsto \left(x \mapsto \int_D \text{cov}[g](x, y) f(y) dy \right) \in L^2(D)$$

- It is a **compact self-adjoint operator**, hence there exists a sequence $(\lambda_n, b_n)_{n \in \mathbb{N}}$ of eigenpairs such that $\lambda_1 \geq \lambda_2 \geq \dots \geq 0$ and such that $(b_n)_{n \geq 0}$ is an **hilbertian basis of $L^2(D)$** .
- For $n \in \mathbb{N}$, the **normalized coordinate of $g(\omega, \cdot)$** in this hilbertian basis with respect to b_n is $Y_n(\omega) = \frac{1}{\sqrt{\lambda_n}} \int_D g(\omega, x) b_n(x) dx$.
- Here, since g is gaussian, the $(Y_n)_{n \geq 1}$ are more precisely **independent gaussian** random variables.
- Then the **Karhunen-Loève expansion of g** is:

$$g(\omega, x) \stackrel{L^2(\Omega \times D)}{=} \sum_{n=1}^{+\infty} \sqrt{\lambda_n} b_n(x) Y_n(\omega).$$

Decay of the eigenvalues : exponential covariance

- We recall the definition of exponential covariance :
 $\text{cov}[g](x, y) = \sigma^2 e^{-\frac{\|x-y\|}{\ell}}$. If we take the norm 1, we have analytic expressions and are able to deduce properties.
- We have for some constant c , $\lambda_n \leq \frac{c\sigma^2}{\ell n^2}$ and $\|b_n\|_\infty \leq C$.
- We can also observe numerically a plateau (of size about $1/\ell$) in the decrease of the λ_n

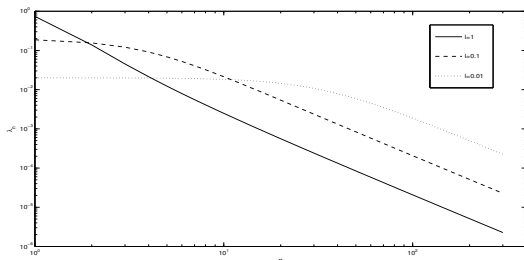


Figure : λ_n versus n , in logarithmic scale, for $\sigma = 1$ and different values of ℓ .

Decay of the eigenvalues : analytic covariance

We suppose that the covariance function $\text{cov}[g]$ is analytic on $D \times D$, then:

Theorem (Frauenfelder, Schwab, Todor, 2005)

There exists two constants $c_1, c_2 > 0$ such that for all $n \geq 1$

$$\lambda_n \leq c_1 e^{-c_2 n^{1/d}}.$$

For any $s > 0$ there exists a constant c_s such that for any $n \geq 1$,

$$\|b_n\|_\infty \leq c_s |\lambda_n|^{-s} \text{ and } \|\nabla b_n\|_\infty \leq c_s |\lambda_n|^{-s}.$$

Exemple of analytic covariance function : the gaussian covariance

$$\text{cov}[g](x, y) = \sigma^2 e^{-\frac{\|x-y\|^2}{\ell^2}}.$$

- **More generally** : the **more regular the covariance function** is, the **more faster the eigenvalues decrease**, and the **more faster the KL expansion converges**.
- The value of the covariance length has also influence on the speed of convergence of the KL expansion.

Approximation of a using a Karhunen Loève expansion

- We define the **approximation a_N of a** :

$$a_N(\omega, x) = e^{g_N(\omega, x)} = e^{\sum_{n=1}^N \sqrt{\lambda_n} b_n(x) Y_n(\omega)}.$$

- We define the **approximation u_N of u** as the solution of:

$$\begin{cases} -\operatorname{div}(a_N(\omega, \cdot) \nabla u_N(\omega, \cdot)) & = f(x) & \text{on } D \\ u_N(\omega, \cdot) & = 0 & \text{on } \partial D. \end{cases}$$

- We have then $\tilde{a}_N(y, x) = e^{\sum_{n=1}^N \sqrt{\lambda_n} b_n(x) y_n}$.
- Hence $u_N(\omega, x) = \tilde{u}_N(Y(\omega, x))$ where \tilde{u}_N is the solution of the deterministic parametrized in \mathbb{R}^N PDE.
- The cost of stochastic collocation methods increases with N (exponentially for the basic collocation).

Strong and weak KL truncature error estimates : assumptions

Assumption

- i) *The eigenfunctions b_n are Hölder continuous with exponent α_0 , where $0 < \alpha_0 < 1/2$*
- ii) *The series $\sum_{n \geq 1} \lambda_n \|b_n\|_{C^{0,\alpha_0}(\bar{D})}^2$ is convergent*
- iii) *$f \in L^p(D)$ for some $p > d$*

\rightsquigarrow This is in particular the case for an **exponential covariance** with norm 1 and for a **gaussian covariance** .

Strong error bound

Under the previous assumption we define, for α such that $0 \leq \alpha \leq \alpha_0$ and $N \in \mathbb{N}$,

$$R_N^\alpha = \sum_{n>N} \lambda_n \|b_n\|_{C^{0,\alpha}(\bar{D})}^2.$$

Theorem (J.C, A.Debussche, 2014)

There exists for any α, β with $0 < \beta < \alpha < \alpha_0$ and $\beta < 1 - \frac{d}{p}$ a constant $C_s(p, q, \alpha, \beta)$ such that for any $N, q \in \mathbb{N}$ we have

$$\|u - u_N\|_{L^q(\Omega, C^{1,\beta}(\bar{D}))} \leq C_s(p, q, \alpha, \beta) \sqrt{R_N^\alpha} \|f\|_{L^p(D)}.$$

Weak error bounds

Our goal is to compute the law of u , therefore we consider the **weak error**, i.e. the error committed by approximating the law of u by the law of u_N .

Theorem (J.C, A. Debussche, 2014)

Let α, β such that $0 < \beta < \alpha < \alpha_0$ and $\beta < 1 - \frac{d}{p}$, then

- for any $\varphi \in C^6(\mathbb{R}, \mathbb{R})$ whose derivatives have at most polynomial growth, there exists a constant $C_{w1}(\beta, f, p, \varphi)$ such that for all $N \in \mathbb{N}$

$$\|\mathbb{E}_w[\varphi(u_N) - \varphi(u)]\|_{C^{1,\beta}(D)} \leq C_{w1}(\beta, f, p, \varphi) R_N^\beta.$$

- for any $\psi \in C^4(C^{0,\beta}(\bar{D}) \times C^{1,\beta}(\bar{D}), \mathbb{R})$ whose differentials have at most polynomial growth, there exists a constant $C_{w2}(\beta, f, p, \varphi)$ such that for all $N \in \mathbb{N}$

$$|\mathbb{E}[\psi(a, u) - \psi(a_N, u_N)]| \leq C_{w2}(\beta, f, p, \varphi) R_N^\beta.$$

Conclusion

Specificities of the random fields frequently used in hydrogeology need to be taken into account :

- The **bounds are not uniform** wrt the random variable, it requires to be careful to get L^p bounds.
↪ you need to track all constants.
- The **spatial regularity** is often law and depend and the choice of the correlation function.
↪ you don't always have usual orders for the spatial discretization.
- The **dimension of the randomness** can be very high and depends strongly on the structure of the correlation function for the asymptotic behaviour but also strongly on the correlation length.
↪ It has to be considered in the choice of a numerical method.

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Transport-diffusion of a solute

- An inert **solute is injected** in the porous media at initial time : typically a pollutant.
- We consider only **molecular diffusion, assumed to be homogeneous and isotropic**.
- We denote by $c(\omega, x, t)$ the **solute concentration**
- The migration of the solute is then described by :

$$\frac{\partial c(\omega, x, t)}{\partial t} + v(\omega, x) \cdot \nabla_x c(\omega, x, t) - D \Delta_x c(\omega, x, t) = 0,$$

and boundary/initial conditions

↪ we recall that v is the random Darcy velocity and we consider the advection-dominated case (Peclet number $\gg 1$).

Quantities of interest

- **First quantity of interest** : the **mean spread** of the solute $S(t) = \mathbb{E}_\omega[S(\omega, t)]$, where

$$S_d(\omega, t) = \int_O c(\omega, x, t)(x_d - G_d(\omega, t))^2 dx$$

and

$$G_d(\omega, t) = \int_O c(\omega, x, t)x_d dx$$

\rightsquigarrow it is the mean value of the **spatial spread in each direction**.

- **Second quantity of interest** : the **mean macro-dispersion** $\mathcal{D}(t)$ defined by

$$\mathcal{D}_d(t) = \mathbb{E}_\omega \left[\frac{dS_d(\omega, t)}{dt} \right].$$

\rightsquigarrow it is the mean value of the **speed of spreading in each direction**.

- **Goal** : compute the mean values of spreading and macro-dispersion
 \rightsquigarrow asymptotic values of macrodispersion ?
- **Question** : **how do molecular diffusion and high heterogeneity impact asymptotic macrodispersion ?**

- A Monte-Carlo method to deal with uncertainty
 N independent realizations of the permeability field $a : a^1, \dots, a^N$
- For each i , p_h^i is a finite element approximation (h is the mesh size) of the solution p^i of the flow equation

$$\operatorname{div}(a^i(\omega, x)\nabla p^i(x)) = 0$$

- $v_h^i = -a\nabla p_h^i$ is then the approximation of the Darcy velocity

Approximation of the mean spread

- We note that the solution $c(x, t)$ of

$$\frac{\partial c(x, t)}{\partial t} + v(x) \cdot \nabla_x c(x, t) - D \Delta_x c(x, t) = 0,$$

is the **density of the solution of the SDE**

$$dX(t) = v(X(t))dt + \sqrt{2D}dW(t)$$

- Therefore **the spread $\mathcal{S}(t)$** can be expressed as

$$\mathcal{S}(t) = \mathbb{E}[(X(t) - \mathbb{E}[X(t)])^2]$$

- The solution of the SDE is approximated through an **Euler scheme**:

$$Y(t_{k+1}) = Y(t_k) + v_h(Y(t_k))\Delta t + \sqrt{2D}\Delta B_k,$$

- Approximation of the mean spread $\mathcal{S}(t)$:

$$\frac{1}{NM} \sum_{i=1}^N \sum_{j=1}^M \left(Y^{ij}(t) - \frac{1}{M} \sum_{j=1}^M Y^{ij}(t) \right)^2$$

Approximation of the macro-dispersion

- Thanks to **Itô formula** we have

$$\begin{aligned}\mathcal{D}_d(t) &= \mathcal{S}'_d(t) \\ &= \mathbb{E} [2(X_d(t) - \mathbb{E}[X_d(t)])(v^i(X_d(t)) - \mathbb{E}[v^i(X_d(t))]) + 2D] \\ &= \mathbb{E}[f(X_d(t))],\end{aligned}$$

where f depends on v .

- Therefore using an **Euler scheme** and a **Monte Carlo method** we compute the approximation of the mean macro-dispersion $\mathcal{S}(t)$:

$$\frac{1}{NM} \sum_{i=1}^N \sum_{j=1}^M f_h^i(Y^{ij}(t)),$$

where

$$dY^{ij}(t_{k+1}) = v_h^i(Y^{ij}(t_k))\Delta t + \sqrt{2D\Delta}B_k^{ij}.$$

We consider two assumptions :

- **Assumption 1** : there exists $0 < \alpha < 1$, $r > d$ with $\alpha < 1 - \frac{d}{r}$ such that for any finite $q \geq 1$ $a \in L^q(\Omega, \mathcal{C}_b^{0,\alpha}(\mathbb{R}^d))$ and $\frac{1}{a_{min}} \in L^q(\Omega)$.
 \rightsquigarrow it is fulfilled in the case of a lognormal field with **exponential covariance**
- **Assumption 2** : there exists $0 < \alpha < 1$, $r > d$ with $\alpha < 1 - \frac{d}{r}$ such that for any finite $q \geq 1$ $a \in L^q(\Omega, \mathcal{C}_b^{1,\alpha}(\mathbb{R}^d))$ and $\frac{1}{a_{min}} \in L^q(\Omega)$.
 \rightsquigarrow it requires **more regularity**, it is among others fulfilled in the case of a lognormal field with gaussian covariance.

Proposition

- Under **Assumption 1**, we get that $p \in L^q(\Omega, \mathcal{C}_b^{1,\alpha}(\mathbb{R}^d))$ and $v \in L^q(\Omega, \mathcal{C}_b^{0,\alpha}(\mathbb{R}^d))$ for any finite $q \geq 1$.
- Under **Assumption 2**, we get that $p \in L^q(\Omega, \mathcal{C}_b^{2,\alpha}(\mathbb{R}^d))$ and $v \in L^q(\Omega, \mathcal{C}_b^{1,\alpha}(\mathbb{R}^d))$ for any finite $q \geq 1$.

Total error on the generalized spread

We define the **generalized spread** :

$$\mathbb{E}_\omega[\psi(\mathbb{E}_\xi[\varphi(X(\omega, \xi, T))])].$$

Theorem (J.C. 2015)

Let $\varphi \in \mathcal{C}_b^{1,\alpha}(\mathbb{R}^d, \mathbb{R}^{d'})$ and $\psi \in \mathcal{C}_b^1(\mathbb{R}^{d'}, \mathbb{R}^{d''})$.

- Under **Assumption 1** we have the bound :

$$\|er\|_{L^2(\Omega \times \Omega')} \leq C \left((\Delta t)^{\frac{\alpha}{2}} + h^\alpha |\ln h| + \frac{1}{\sqrt{M}} + \frac{1}{\sqrt{N}} \right).$$

- Under **Assumption 2** we have the bound :

$$\|er\|_{L^2(\Omega \times \Omega')} \leq C \left((\Delta t)^{\frac{1+\alpha}{2}} + h |\ln h| + \frac{1}{\sqrt{M}} + \frac{1}{\sqrt{N}} \right).$$

Total error on the generalized macro-dispersion

- For some vector-valued functions φ and ψ we define the **generalized spread** :

$$\frac{d}{dt} \mathbb{E}_\omega [\psi(\mathbb{E}_\xi[\varphi(X(\omega, \xi, T))])],$$

which thanks to **Itô formula** is equal to

$$\mathbb{E}_\omega [D\psi(\mathbb{E}_\xi[\varphi(X)]) \cdot \mathbb{E}_\xi [D\varphi(X) \cdot v(X) + D\Delta\varphi(X)]].$$

- The **corresponding error** is then

$$\begin{aligned} \bar{E}r &= \mathbb{E}_\omega [D\psi(\mathbb{E}_\xi[\varphi(X)]) \cdot \mathbb{E}_\xi [D\varphi(X) \cdot v(X) + D\Delta\varphi(X)]] \\ &- \frac{1}{N} \sum_{i=1}^N \left[D\psi \left(\frac{1}{M} \sum_{j=1}^M \varphi(X_{n,h}^{i,j}) \right) \cdot \frac{1}{M} \sum_{j=1}^M (D\varphi(X_{n,h}^{i,j})) \cdot v_h^i(X_{n,h}^{i,j}) + D\Delta\varphi(X_{n,h}^{i,j}) \right] \end{aligned}$$

\rightsquigarrow we use similar techniques as for the spread, but we **need more regularity**.

Total error on the generalized macro-dispersion

Theorem (J.C 2015)

Let $\bar{\varphi} \in \mathcal{C}_b^{3,\alpha}(\mathbb{R}^d, \mathbb{R}^{d'})$ and $\bar{\psi} \in \mathcal{C}_b^2(\mathbb{R}^{d'}, \mathbb{R}^{d''})$, then under *Assumption 2*, we have

$$\|\bar{E}r(\omega, \xi)\|_{L_\omega^2} \leq C \left((\Delta t)^{\frac{1+\alpha}{2}} + h|\ln(h)| + \frac{1}{\sqrt{M}} + \frac{1}{\sqrt{N}} \right)$$

Difficulties :

- v has not enough **regularity** to apply the classical results
- we need **explicit and sharp bounds** (because the constant will depend on ω and we need integrability with respect to ω)
- we need to deal with the **spatial approximation** at the same time as the time discretization in the weak error
- in the case of the macro-spreading, we have to deal with the presence of the random function v_h inside “the test function”.

Numerical results and answer to the hydrogeology question

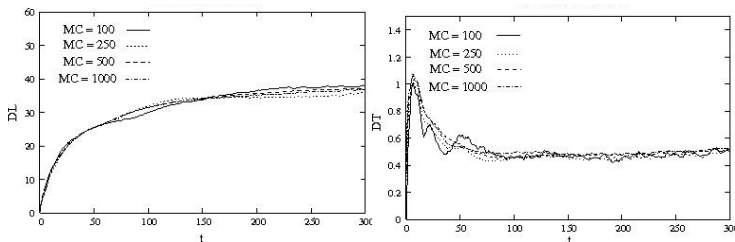


Figure : Case of an exponential covariance with $\lambda = 10$, $\sigma = 1, 5$, $D = 0, 1$, $d = 2$, 10000 particles, realized with the software platform H2OLAB

“For large heterogeneities ($\sigma^2 > 1$), diffusion induces a significant longitudinal macro-dispersion decrease and a transverse macro-dispersion increase larger than expected.”

Asymptotic dispersion in 2D heterogeneous porous media determined by parallel numerical simulations, de Dreuzy *et al.*, Water Resources Research,

Perspectives

- Use **spatial ergodicity** type properties to get better theoretical convergence rates in order to :
 - ▶ justify theoretically the use of stochastic particular methods
 - ▶ estimate the speed of convergence to the asymptotic value
 - ▶ get sharp a priori error estimates
 - ▶ fit the discretization parameters efficiently
- More complicated problems :
 - ▶ fractured media
 - ▶ chemical reactions
 - ▶ add cinematic diffusion
 - ▶ computation of exit times

Thank you for your attention !