

Anim3D Project

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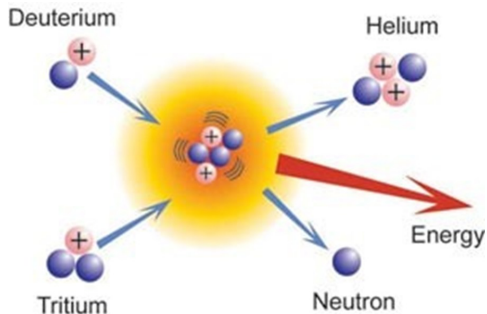
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- Motivations
- IsoGeometric Analysis
- Numerical results
- Conclusion and perspectives

Plasma Physics

- **Plasma:** For very high temperatures, the gas are ionized and gives a plasma which can be controlled by magnetic and electric fields.
- **Thermonuclear fusion:** The MHD allows to describe some configuration where the collision are not so small or for long time behavior.
- **Astrophysics:** The MHD describe a lot of astrophysics configuration: supernovae explosion, solar wind and instabilities etc.
- **Context:** in the case we consider the application of the MHD to the simulation of **Tokamak instabilities**.



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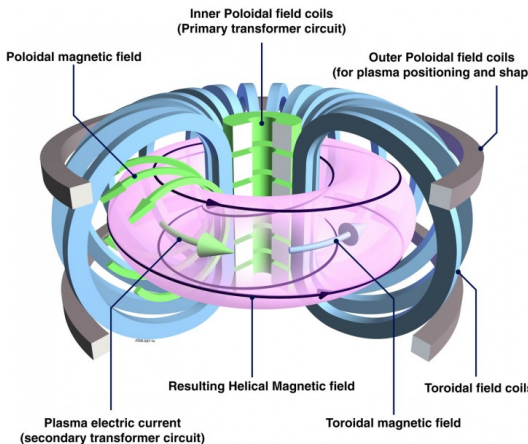
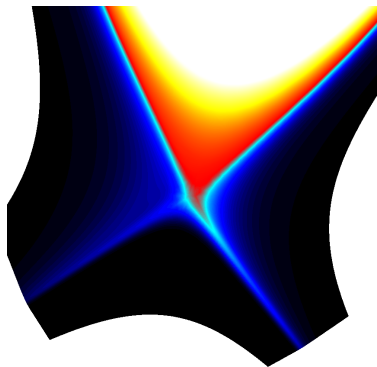


Figure: Tokamak

Example of Application : MHD and ELM

- In the tokamak **some instabilities** can appear in the plasma.
- The simulation of these instabilities is an **important subject for ITER**.
- Example of Instabilities in the tokamak :
 - **Disruptions**: Violent instabilities which can critically damage the Tokamak.
 - **Edge Localized Modes (ELM)**: Periodic edge instabilities which can damage the Tokamak.
- These instabilities are linked to the **very large gradient of pressure and very large current** at the edge.
- Many aspects of these instabilities are described by **fluid models** (MHD resistive and diamagnetic or extended)

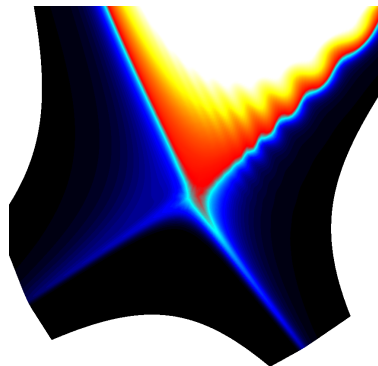
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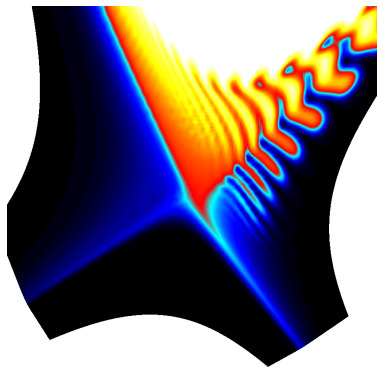
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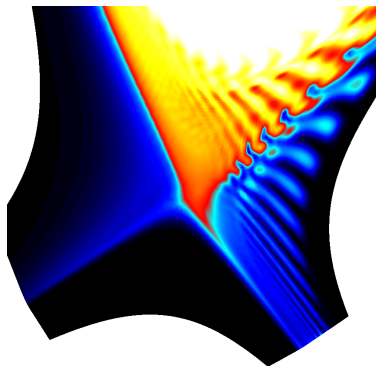
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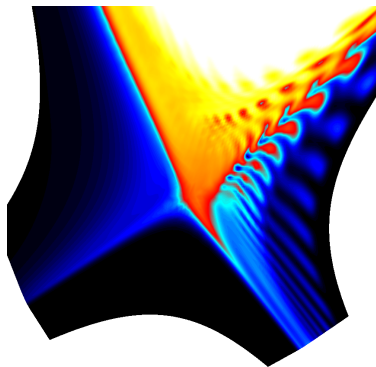
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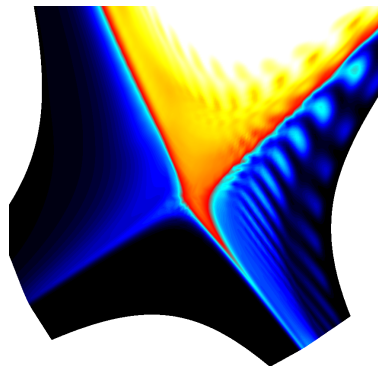
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Discretization

B-Splines

To create a family of *B-splines*, we need a non-decreasing sequence of knots $T = (t_i)_{1 \leq i \leq N+k}$, also called **knot vector**, with $k = p + 1$.

Each set of knots $T_j = \{t_j, \dots, t_{j+p}\}$ will generate a *B-spline* N_j .

Definition (B-Spline serie)

The j -th B-Spline of order k is defined by the recurrence relation:

$$N_j^k = w_j^k N_j^{k-1} + (1 - w_{j+1}^k) N_{j+1}^{k-1}$$

where,

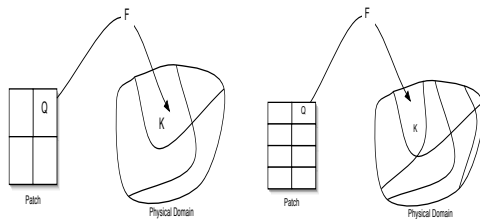
$$w_j^k(x) = \frac{x - t_j}{t_{j+k-1} - t_j}$$

$$N_j^1(x) = \chi_{[t_j, t_{j+1}[}(x)$$

for $k \geq 1$ and $1 \leq j \leq N$.

Discretization

The IsoGeometric Approach



Grid generation: the use of $h/p/k$ -refinement keeps the mapping F unchanged.

- Compact support
- Partition of Unity
- Affine covariance
- IsoParametric concept
- Error estimates in Sobolev norms
- Exacte DeRham discrete sequence

Discretization

Refinement strategies in IGA

Refinement strategies

Refining the grid can be done in 3 different ways. This is the most interesting aspects of B-splines basis.

h-refinement by inserting new knots. It is the equivalent of mesh refinement of the classical finite element method.

p-refinement by elevating the B-spline degree. It is the equivalent of using higher finite element order in the classical FEM.

k-refinement by increasing / decreasing the regularity of the basis functions (increasing / decreasing multiplicity of inserted knots).

r-refinement moving the control points to reduce a given error estimate

the use of **k-refinement** strategy is more efficient than the classical p-refinement, as it reduces the dimension of the basis.

Parallelism

Domain Decomposition

Available algorithms

- Tensor decomposition, when using Tensor Spaces
- Metis (ParMetis will be added later)

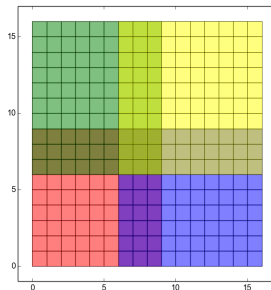
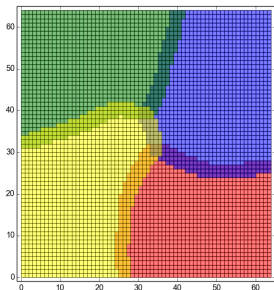
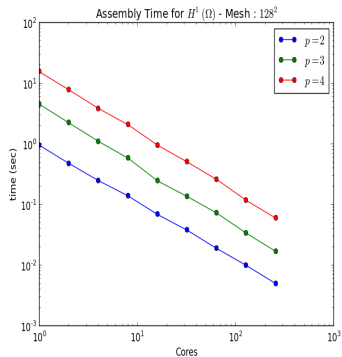
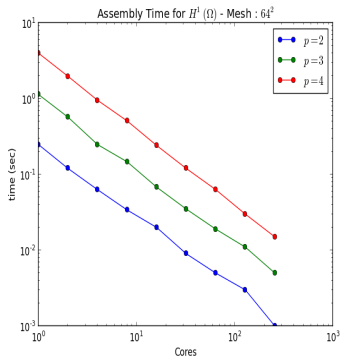


Figure: Metis (left) and tensor (right) partitioning.

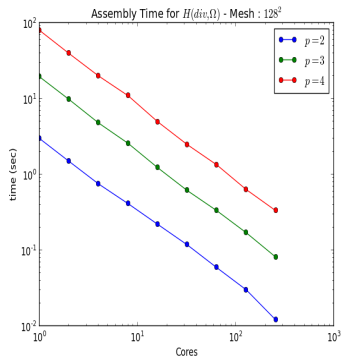
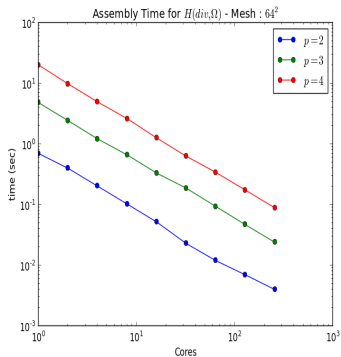
Numerical results: Parallel runs

The 2D case



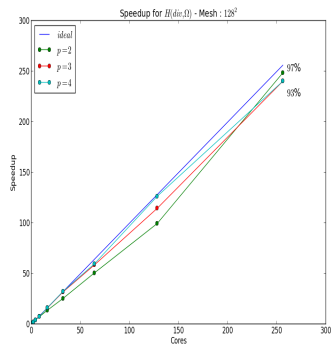
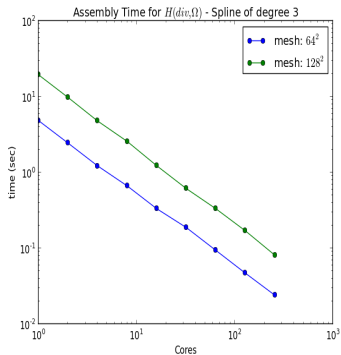
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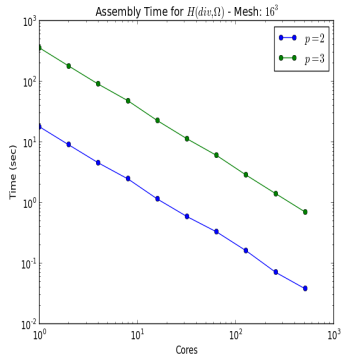
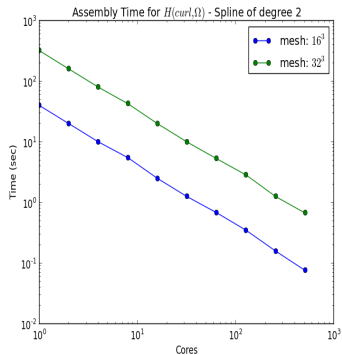
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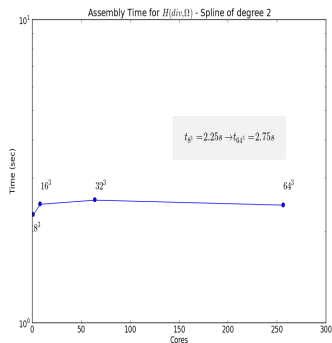
Numerical results: Parallel runs

The 3D case



Numerical results: Parallel runs

The 3D case



Statistics: Quadratic Splines on a grid 32^3 :

- 23'101'440 non zeros for $H(\text{curl})$
- 98'304 dofs for $H(\text{curl})$
- 13'860'864 non zeros for $H(\text{div})$
- 98'304 dofs for $H(\text{div})$

Equilibrium

- We consider the resistive MHD with $\mathbf{v} = 0$. We obtain the following equilibrium

$$\begin{cases} \mathbf{v} = 0, & \mathbf{J} \times \mathbf{B} = \nabla p \\ \partial_t \mathbf{B} = \frac{\eta}{\nu_0} \Delta \mathbf{B} \end{cases}$$

- $\tau \ll \tau_{diff}$ with τ and τ_{diff} the characteristic time of transport and the diffusion.

MHD equilibrium

- The equilibrium is mainly defined by the force balanced

$$\mathbf{J} \times \mathbf{B} = \nabla p$$

- The equilibrium induces that $\mathbf{B} \cdot \nabla p = 0$, $\nabla \cdot \mathbf{J} = 0$ and we assume that $\nabla p \cdot \mathbf{e}_\phi = 0$.
- In a Tokamak we assume that

$$\mathbf{B} = \mu_0 \frac{F(\psi, Z)}{R} \mathbf{e}_\phi + \frac{1}{R} (\nabla \psi \times \mathbf{e}_\phi)$$

- After some computation we obtain the following equation

Grad-Shafranov equation

$$\Delta^* \psi = -\mu_0 R^2 \frac{dp(\psi)}{d\psi} - \mu_0^2 F(\psi) \frac{dF(\psi)}{d\psi}$$

Grad-Shafranov Shift and β plasma

Shift

- **Property of GS operator:** induce a shift of the magnetic surface
- **Shift estimation:** $\frac{\Delta}{r} \approx \beta_p \frac{r}{R_0}$
- with r and R_0 the minor and major radius.
- $\beta_p = \frac{2\mu_0 |p|}{|B_p|}$ the ratio of the pressure and poloidal magnetic pressure.

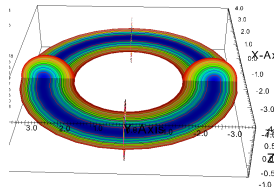


Figure: 3D equilibrium

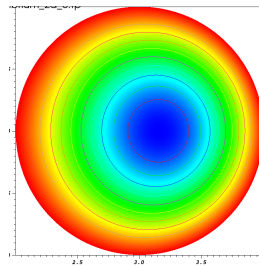


Figure: poloidal cut of equilibrium

Test case

- **Discretization:** 32*32*8 cells with third order B-Splines
- **Solvers** Picard nonlinear solver + GMRES
- **Physical problems:** $\beta_p \approx 10^{-2}$, $R_0 = 3$ and $\frac{\beta}{\beta_p} \approx 10$

Anisotropic diffusion

- Model :

$$\partial_t T - \nabla \cdot \left((k_{\parallel} - k_{\perp})(\mathbf{b} \otimes \mathbf{b}) \nabla T + k_{\perp} \nabla T \right) = 0$$

with $k_{\parallel} \gg k_{\perp}$.

- The magnetic field is construct solving the equilibrium.
- In this case $k_{\parallel} = 200$ and $k_{\perp} = 0$. The diffusion is along the magnetic lines.

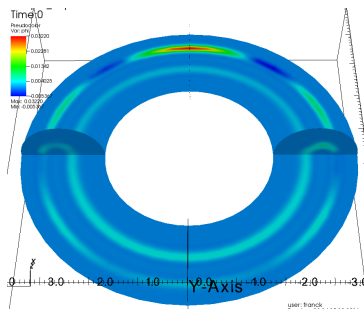
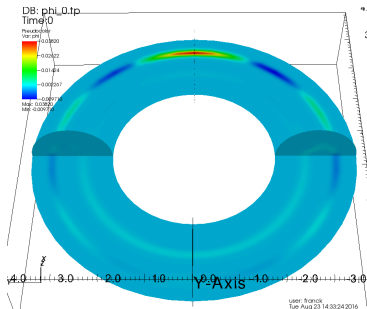


Figure: Left: solution after time $T = 0.5$. Right: solution after time $T = 7$

Reduce MHD model

Single fluid resistive MHD

$$\begin{cases} \partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = 0, \\ \rho \partial_t \mathbf{v} + \rho \mathbf{v} \cdot \nabla \mathbf{v} + \nabla p = \mathbf{J} \times \mathbf{B} - \nabla \cdot \bar{\bar{\mathbf{n}}}, \\ \partial_t \rho + \mathbf{v} \cdot \nabla \rho + \rho \nabla \cdot \mathbf{v} + \nabla \cdot \mathbf{q} = 0 \\ \partial_t \mathbf{B} = -\nabla \times (-\mathbf{v} \times \mathbf{B} + \eta \mathbf{J}), \\ \nabla \cdot \mathbf{B} = 0, \quad \nabla \times \mathbf{B} = \mathbf{J}. \end{cases}$$

- **Reduced MHD model:** Reduce the number of variables and eliminate the fast waves in the reduced MHD model.
- We consider the cylindrical coordinate $(R, Z, \phi) \in \Omega \times [0, 2\pi]$.

Reduced MHD: Assumption

$$\mathbf{B} = \frac{F_0}{R} \mathbf{e}_\phi + \frac{1}{R} \nabla \psi \times \mathbf{e}_\phi, \quad \mathbf{v} = -R \nabla u \times \mathbf{e}_\phi + v_{\parallel} \mathbf{B}$$

with u the electrical potential, ψ the magnetic poloidal flux, v_{\parallel} the parallel velocity.

- **Initialization:** we use ψ and pressure equilibrium, a zero velocity ($u = v_{\parallel} = 0$).
- **Wave structure:** low Mach and low β regime \rightarrow a large ratio between wave speeds.
- This problem coupled with hyperbolic structure generate **ill-conditioned** problem.

Preconditioning

- The implicit system after linearization is given by

$$\begin{pmatrix} \mathbf{B}^{n+1} \\ p^{n+1} \\ \mathbf{u}^{n+1} \end{pmatrix} = \begin{pmatrix} A_{\mathbf{B},p} & C_{\mathbf{B},p,\mathbf{u}} \\ C_{\mathbf{u},\mathbf{B},p} & A_{\mathbf{u}} \end{pmatrix}^{-1} \begin{pmatrix} R_{\mathbf{B}} \\ R_p \\ R_{\mathbf{u}} \end{pmatrix}$$

- with $A_{\mathbf{B},p}$ and $A_{\mathbf{u}}$ the advection terms linked to \mathbf{B} and p (resp \mathbf{u}), $C_{\mathbf{B},p,\mathbf{u}}$ and $C_{\mathbf{u},\mathbf{B},p}$ the coupling terms which gives the **Alfven and acoustic waves**.
- The solution of the system is given by

$$\begin{pmatrix} \mathbf{B}^{n+1} \\ p^{n+1} \\ \mathbf{u}^{n+1} \end{pmatrix} = \begin{pmatrix} I_d & A_{\mathbf{B},p}^{-1} C_{\mathbf{B},p,\mathbf{u}} \\ 0 & I_d \end{pmatrix} \begin{pmatrix} A_{\mathbf{B},p}^{-1} & 0 \\ 0 & P_{schur}^{-1} \end{pmatrix} \begin{pmatrix} I_d & 0 \\ -C_{\mathbf{u},\mathbf{B},p} A_{\mathbf{B},p}^{-1} & I_d \end{pmatrix} \begin{pmatrix} R_{\mathbf{B}} \\ R_p \\ R_{\mathbf{u}} \end{pmatrix}$$

- Using the previous Schur decomposition, we obtain the following algorithm:

$$\left\{ \begin{array}{l} \text{Predictor : } A_{\mathbf{B},p} \begin{pmatrix} \mathbf{B}^* \\ p^* \end{pmatrix} = \begin{pmatrix} R_{\mathbf{B}} \\ R_p \end{pmatrix} \\ \text{Velocity evolution : } P_{schur} \mathbf{u}^{n+1} = \begin{pmatrix} -C_{\mathbf{u},\mathbf{B},p} \begin{pmatrix} \mathbf{B}^{n+1} \\ p^{n+1} \end{pmatrix} + R_{\mathbf{u}} \end{pmatrix} \\ \text{Corrector : } A_{\mathbf{B},p} \begin{pmatrix} \mathbf{B}^{n+1} \\ p^{n+1} \end{pmatrix} = A_{\mathbf{B},p} \begin{pmatrix} \mathbf{B}^* \\ p^* \end{pmatrix} - C_{\mathbf{B},p,\mathbf{u}} \mathbf{u}^{n+1} \end{array} \right.$$

- Preconditioning**: we approximate the Schur complement by a **multi-scale elliptic operator**.
- Using classical **Multi-grids and auxiliary-space theory** we can perform the invert of the Schur approximation.

These developments were done in the **JorekDjango** framework. It is written in *Fortran2008*, using MPI.

- **CLAPPIO** Input/Output Library
- **PLAF** Parallel Linear Algebra Library
- **SPL** Library for NURBS/B-Splines
- **DISCO** Abstract Discretization Context Library
- **FEMA** Library of Finite Elements Assemblers

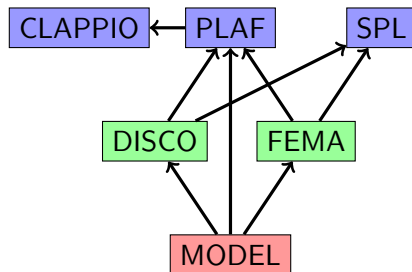


Figure: Structure of the JorekDjango Framework

Conclusions

- Weak scalability. In $2D$, we get a speedup of 104% on H^1 and 94 – 97% for $H(\text{div})$
- Different models have been implemented during this summer-school (Equilibrium, Anisotropic Diffusion, Reduced-MHD)

Ongoing work and Perspectives

- Validation of OpenMP
- Physics-Based Preconditioner for the Reduced-MHD (model199)
- Add more physics (model303)

Thanks!