

Simulations of Kinetic Electrostatic Electron Nonlinear (KEEN) Waves with Variable Velocity Resolution Grids and High-Order Time-Splitting

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Joint work : B. Afeyan, F. Casas, N. Crouseilles, A. Dodhy, E. Faou, M. Mehrenberger, E. Sonnendrücker, submitted to EPJD, special issue for the Vlasovia Conference 2013.

A short introduction to the Physics of KEEN waves

- KEEN waves are asymptotic states in Vlasov plasmas
 - non-stationary
 - non linear
 - self-organized
- Driven by a ponderomotive force generated by 2 crossing laser beams. Perturb a range in velocity depending on :
 - the **amplitude** of the drive
 - the **duration** of the drive

⇒ We want here to study weakly driven cases

The equations to solve

Periodic $1D \times 1D$ Vlasov-Poisson equation

$$\partial_t f + v \partial_x f + (E - E_{\text{Pond}}) \partial_v f = 0, \quad \partial_x E = \rho(t, x) - 1, \quad \rho(t, x) = \int_{\mathbb{R}} f dv.$$

Here $E_{\text{Pond}}(t, x)$ is of the form

$$E_{\text{Pond}}(x, t) = a_{\text{Dr}} k_{\text{Dr}} a(t) \sin(k_{\text{Dr}} x - \omega_{\text{Dr}} t),$$

with

$$a(t) = \frac{g(t) - g(t_0)}{1 - g(t_0)},$$

$$g(t) = 0.5(\tanh(\frac{t - t_L}{t_{wL}}) - \tanh(\frac{t - t_R}{t_{wR}})),$$

$t_0 = 0$, $t_L = 69$, $t_{wL} = t_{wR} = 20$, $t_R = 207 + T_{\text{Dr}}$, $k_{\text{Dr}} = 0.26$,
 $\omega_{\text{Dr}} = 0.37$. The initial condition is

$$f_0(x, v) = \frac{1}{\sqrt{2\pi}} \exp\left(-\frac{v^2}{2}\right), \quad (x, v) \in [0, 2\pi/k] \times [-6, 6].$$

Outline

- CEMRACS'2012 : VOG project
- Non uniform cubic splines
- High order splitting
- Numerical results
 - L^2 norm
 - ρ -harmonics
- Towards physical results
 - RMS quantities
 - δ -f distribution function
- Conclusion/perspectives

KEEN wave simulation on uniform grid

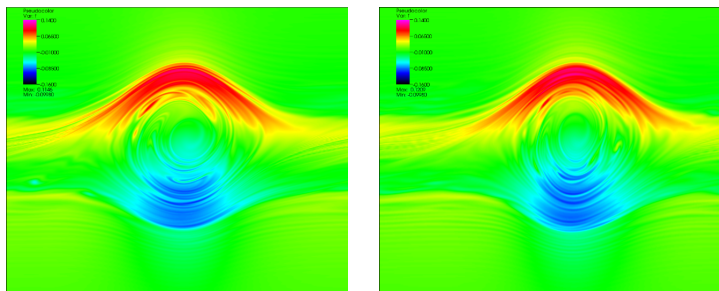


FIGURE: $f(1000, x, v) - f_0(x, v)$, $\Delta t = 0.05$. Cubic splines, with $N_x = N_v = 4096$ on CPU (left) and Lagrange interpolation of degree 17 $N_x = N_v = 2048$ on GPU double precision (right).

Canonical drive $a_{Dr} = 0.2$, $T_{Dr} = 100$.

How to change the uniform grid ?

- Two-grid
- Non-uniform grid
- Fully adaptative

First simple strategy : non-uniform cubic splines on a two-grid mesh

Other alternatives :

- history of adaptive schemes in CALVI INRIA project
- some more recent contributions
 - real two-grid method/sparse grid (Crouseilles, M.M., Sonnendrücker, Kormann...)
 - velocity discretization (Helluy, Navoret, Pham...)
 - discontinuous Galerkin (Madaule-Restelli, Qiu et al ; Steiner...)
 - forward PIC (Campos-Pinto ; Larson & Young)

Difficulties of **efficient** adaptive schemes

Need to measure the efficiency of the implementation

- common for PIC codes
- still not so familiar for grid based codes

A measure of the efficiency

Number of advected points in double precision in 1 d per μ -second, per processor. We need

- The number of points in x N_x
- The number of points in v N_v
- The time step Δt and final time T
- The number of sub-steps s per time step (2 or 3 for Strang splitting)
- The time t of the simulation in seconds
- The number of processors n_{proc}

Formula is

$$E = \frac{N_x \cdot N_v \cdot T / \Delta t \cdot s}{10^6 \cdot t \cdot n_{\text{proc}}}$$

Some examples of efficiency

| Size | Sequential | GPU |
|--------------------|------------|-----|
| 256×256 | 29 | 73 |
| 2048×2048 | 15 | 414 |

In our code (simulation in Selalib), efficiency is around 4, grids until 2^{29} points ; 256 processors are generally used but up to 2048 ; parallelization is mixed MPI/OPENMP

- Still room for improvements !
- GPU no more competitive, using 128 or more processors.
- PIC ($2D \times 2D$) in Selalib : 13 points advected per μ -second (30 using linear splines) [Communicated by S. Hirstoaga]
- PIC up to $10^{12} \simeq 2^{40}$ points [see. S. Collombi talk]

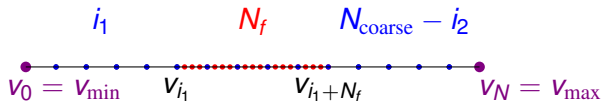
A simple mesh with a refined zone

$$\Delta v_{\text{coarse}} = \frac{v_{\text{max}} - v_{\text{min}}}{N_{\text{coarse}}}, \quad \Delta v_{\text{fine}} = \frac{v_{\text{max}} - v_{\text{min}}}{N_{\text{fine}}}$$

N_{fine} is an integer multiple of N_{coarse} .

The refined zone is chosen with $0 \leq i_1 < i_2 \leq N_{\text{coarse}}$ and the total number of cells is

$$N = i_1 + N_f + N_{\text{coarse}} - i_2, \quad N_f = \frac{N_{\text{fine}}}{N_{\text{coarse}}}(i_2 - i_1)$$



Mesh under consideration

- Canonical drive ($a_{Dr} = 0.2$) : 1/6 of velocity range

$$v_{i_1} \simeq 0.375, v_{i_2} \simeq 2.25$$

- Weak drive ($a_{Dr} = 0.00625$) : 1/30 of velocity range

$$v_{i_1} \simeq 1.2, v_{i_2} \simeq 1.6$$

We have fixed the ratio

$$\frac{N_{\text{fine}}}{N_{\text{coarse}}} = \frac{\Delta v_{\text{coarse}}}{\Delta v_{\text{fine}}} = 32.$$

*Method still parametrized by the **total** number of points $N = N_v$*

Interpolation on the mesh

- Thanks to dimensional splitting, we have to solve the 1d constant advection

$$\partial_t u + c \partial_v u = 0$$

- Use of cubic splines on non uniform mesh
- Conservative form

$$\int_{v_j}^{v_{j+1}} u(v, \Delta t) dv = \int_{v_j - c\Delta t}^{v_{j+1} - c\Delta t} u(v, 0) dv.$$

- Unknowns are in the **middle** of the velocity cells
- Cubic splines interpolation for the primitive
- Remove/add total mass to deal with primitive that is periodic

Note : Lagrange of degree 17 is used for advection in x . (see Charles-Després-M, SINUM 2013)

Strang splitting

Transport in x over $\Delta t/2$

$$\partial_t f(t, x, v) + v \partial_x f(t, x, v) = 0, \quad t \rightarrow t + \frac{\Delta t}{2}$$

Update of E through the Poisson equation

Transport in v over Δt

$$\partial_t f(t, x, v) + E(t, x) \partial_v f(t, x, v) = 0$$

Transport in x over $\Delta t/2$

$$\partial_t f(t, x, v) + v \partial_x f(t, x, v) = 0, \quad t \rightarrow t + \frac{\Delta t}{2}$$

Splitting is parametrized by

$$s = 3, \sigma_{\text{init}} = 1, a_1 = 1/2, a_2 = 1, a_3 = 1/2.$$

6-th order time splitting

$$s = 11,$$

$$\begin{aligned}
 a_1 &= 0.0490864609761162454914412 \\
 &\quad -2\Delta t^2(0.0000697287150553050840999), \\
 a_2 &= 0.1687359505634374224481957, \\
 a_3 &= 0.2641776098889767002001462 \\
 &\quad -2\Delta t^2(0.000625704827430047189169) \\
 &\quad +4\Delta t^4(-2.91660045768984781644 \cdot 10^{-6}), \\
 a_4 &= 0.377851589220928303880766, \\
 a_5 &= 0.1867359291349070543084126 \\
 &\quad -2\Delta t^2(0.00221308512404532556163) \\
 &\quad +4\Delta t^4(0.0000304848026170003878868) \\
 &\quad -8\Delta t^6(4.98554938787506812159 \cdot 10^{-7}), \\
 a_6 &= -0.0931750795687314526579244,
 \end{aligned}$$

together with $a_{6+i} = a_{6-i}$, $i = 1, \dots, 5$ and $\sigma_{\text{init}} = 0$.

History of time splitting

- Strang : second order
- Yoshida : classical triple jump
- Blanes & Moan : optimized coefficients
- For an updated litterature, see Y. Güclü et al. JCP 270 (2014), p716.
 - Strang by far the most popular
 - High order splitting is nice when combined with high order phase-space discretization with fine mesh
 - Reducing the number of stages/ size of negative coefficients is mandatory
- Here : new optimized coefficients for Vlasov-Poisson

How to find the coefficients ? Use of abstract form

- Hamiltonian is defined as $H(f) = T(f) + U(f)$, with

$$H(f) = \int_{(0,2\pi) \times \mathbb{R}} \frac{|v|^2}{2} f(x, v) dx dv + \int_{(0,2\pi)} \frac{1}{2} |E(f)(x)|^2 dx.$$

- Vlasov-Poisson can be written as

$$\partial_t f - \left\{ \frac{\delta H}{\delta f}(f), f \right\} = 0,$$

with Poisson bracket $\{f, g\} = \partial_x f \cdot \partial_v g - \partial_v f \cdot \partial_x g$

- Equivalent to

$$\forall G, \quad \frac{d}{dt} G(f) = [H, G](f)$$

with Poisson bracket for functionals

$$[H, G] = \int_{(0,2\pi) \times \mathbb{R}} \frac{\delta H}{\delta f}(f) \left\{ \frac{\delta G}{\delta f}(f), f \right\} dx dv = -[G, H].$$

Relation between T and U

- We have the relation (in $1D \times 1D$)

$$[[T, U], U] = 2U$$

- This implies RKN condition

$$[[[T, U], U], U] = 0.$$

Derivation of high order methods

- We write

$$\psi_p^\tau = e^{b_1\tau U} e^{a_1\tau T} e^{b_2\tau U} \dots e^{b_2\tau U} e^{a_1\tau T} e^{b_1\tau U}$$

- Using for example the Baker-Campbell-Hausdorff formula, we get

$$\psi_p^\tau = e^{\tau(T+U)+R(\tau)}$$

where

$$\begin{aligned} R(\tau) = & \tau^2 p_{21}[T, U] + \tau^3 (p_{31}[[T, U], T] + p_{32}[[T, U], U]) + \\ & \tau^4 (p_{41}[[[T, U], T], T] + p_{42}[[[T, U], U], T] \\ & + p_{43}[[[T, U], U], U]) + \mathcal{O}(\tau^5) \end{aligned}$$

p_{ij} are polynomials in the parameters a_i, b_j .

Use of modified potentials

- We can incorporate the flow associated to $[[T, U], U]$.

$$\psi_p^\tau = e^{b_1\tau U + c_1\tau^3[[T,U],U]} e^{a_1\tau T} \dots e^{a_1\tau T} e^{b_1\tau U + c_1\tau^3[[T,U],U]}$$

\Rightarrow Achievement of given order with less stages

- further reduction of $e^{\tau C_i}$, when $[[T, U], U] = 2U$, with

$$\begin{aligned} C_i &= b_i U + c_i \tau^2 [[T, U], U] + d_i \tau^4 W_{5,1} + e_i \tau^6 W_{7,1} \\ &= (b_i + 2c_i \tau^2 + 4d_i \tau^4 - 8e_i \tau^6) U \end{aligned}$$

with

$$\begin{aligned} W_{5,1} &= [U, [U, [T, [T, U]]]] \\ W_{7,1} &= [U, [T, [U, [U, [T, [T, U]]]]]]. \end{aligned}$$

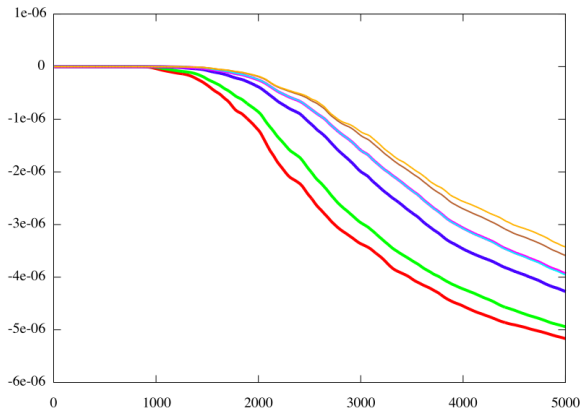
L_2 norm

FIGURE: Time evolution of relative L_2 -norm for the **small drive** amplitude case for different runs, with uniform and non-uniform velocity mesh. The 6-th order time splitting scheme is used for all the runs.

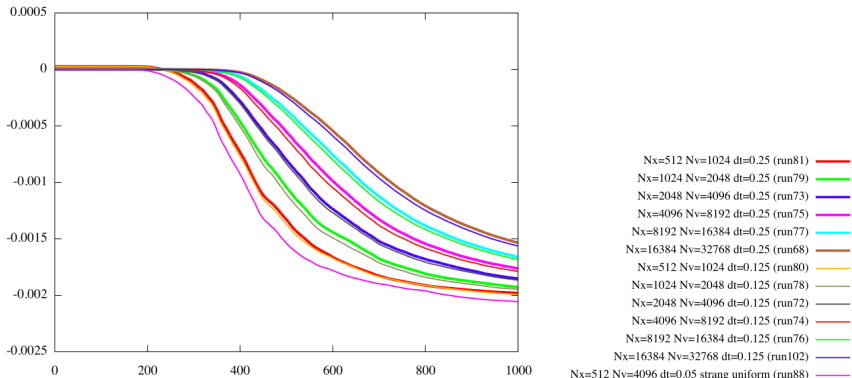
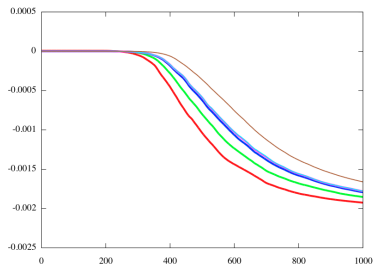
L_2 norm

FIGURE: Time evolution of relative L_2 -norm for the canonical drive amplitude case, with non-uniform velocity grids, 6-th order time splitting and different phase-space resolutions and time steps ; only the last plot uses Strang splitting and uniform velocity grid.

L_2 norm

- $N_x=2048$ $N_v=2048$ $dt=0.25$ non-uniform (run45) —
- $N_x=2048$ $N_v=4096$ $dt=0.25$ non-uniform (run46) —
- $N_x=2048$ $N_v=8192$ $dt=0.25$ non-uniform (run47) —
- $N_x=2048$ $N_v=16384$ $dt=0.25$ non-uniform (run48) —
- $N_x=2048$ $N_v=131072$ $dt=0.25$ non-uniform (run55) —
- $N_x=8192$ $N_v=16384$ $dt=0.25$ non-uniform (run77) —

FIGURE: Time evolution of relative L_2 -norm for the canonical drive amplitude case : comparison of non-uniform runs with different velocity resolution and **space resolution** $N_x = 2048$, and with canonical non-uniform run, using $N_x = 8192$. Here $\Delta t = 0.25$.

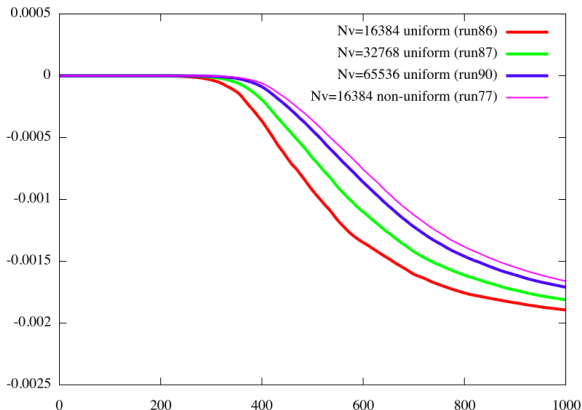
L_2 norm

FIGURE: Time evolution of relative L_2 -norm for the canonical drive case : comparison of **uniform** grid runs with different velocity resolution and a canonical non-uniform velocity grid run. Here $N_x = 8192$ throughout.

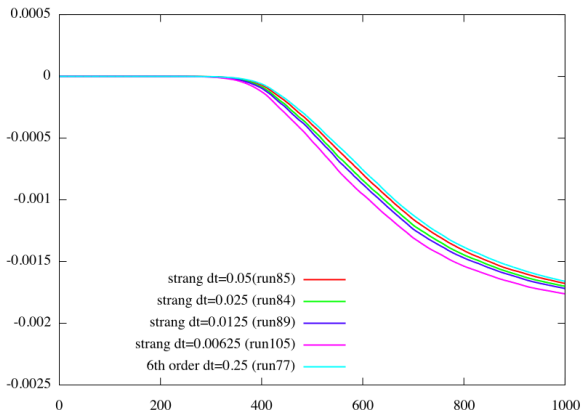
L_2 norm

FIGURE: Time evolution of relative L_2 -norm for the canonical drive amplitude case : comparison of non-uniform runs with **Strang** splitting and different time steps ; and also with canonical non-uniform run with 6th order scheme. Here $N_x = 8192$, $N_v = 16384$.

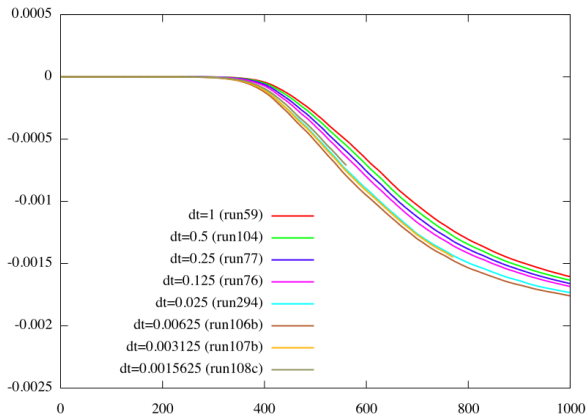
L_2 norm

FIGURE: Time evolution of relative L_2 -norm for the canonical drive amplitude case : comparison of non-uniform runs with **6th order** scheme and different time steps. Here $N_x = 8192$, $N_v = 16384$.

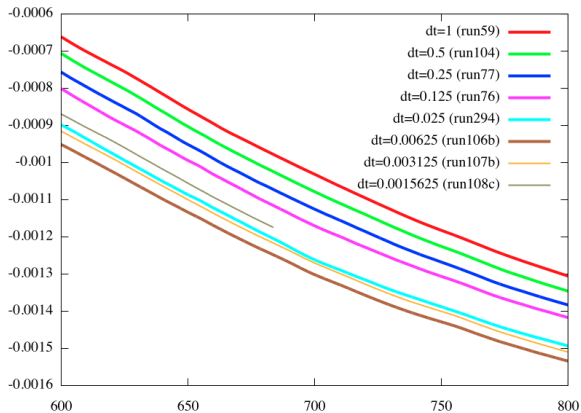
L_2 norm

FIGURE: Zoom of previous picture for $t \in [600, 800]$.

Comparisons for the small drive

- default scheme : non-uniform cubic splines with 6-th order time splitting
- $N_v = 262144$ (uniform) comparable to $N_v = 16384$ (top-left)
- $\Delta t = 0.0078125$ (Strang) comparable to $\Delta t = 0.5$ (bottom-left)
- $N_v = 16384$ comparable to $N_v = 65536$ (top-right)
- $\Delta t = 0.5$ comparable to $\Delta t = 0.25$ (bottom-right)

Comparisons on the ρ -harmonics (small drive)

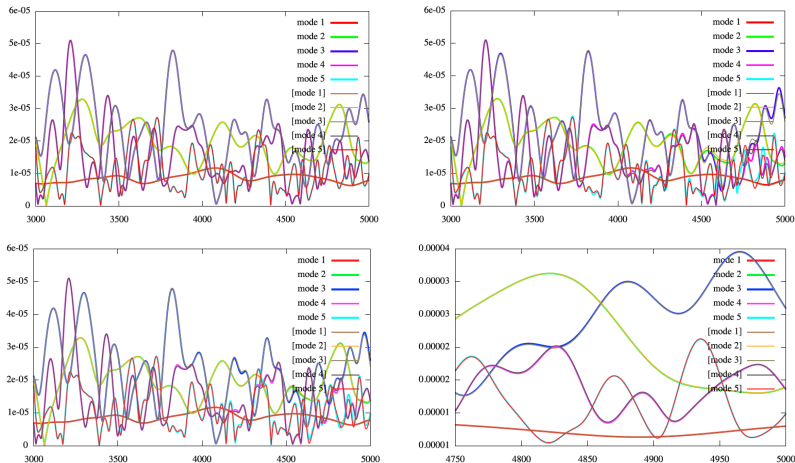


FIGURE: Comparison of the time evolution of the first 5 ρ harmonics, in the **small drive** amplitude case.

Comparisons for the canonical drive

- default scheme : non-uniform cubic splines with 6-th order time splitting
- $N_v = 65536$ (uniform) comparable to $N_v = 16384$ (top-left)
- $\Delta t = 0.0125$ (Strang) comparable to $\Delta t = 0.25$ (bottom-left)
- $(N_x, N_v) = (8192, 16384)$ comparable to $(N_x, N_v) = (16384, 32768)$ (top-right)
- $\Delta t = 0.25$ comparable to $\Delta t = 0.125$ (bottom-right)

Comparisons on the ρ -harmonics (canonical drive)

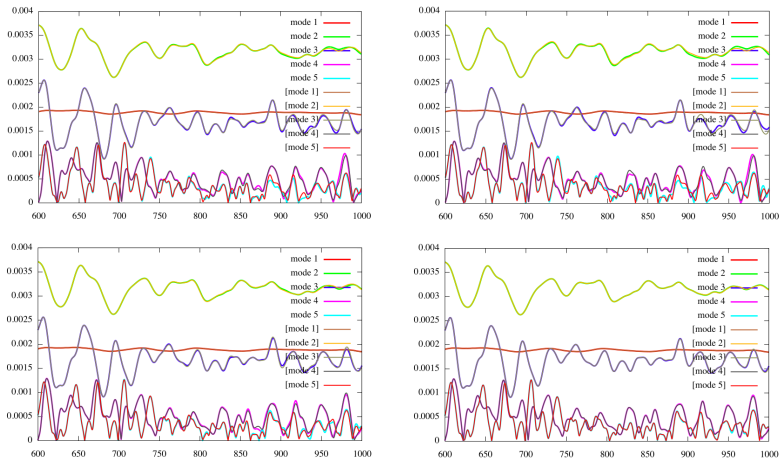


FIGURE: Comparison of the time evolution of the first 5 ρ harmonics, in the **canonical drive** case

Study of KEEN waves formation

- Other diagnostics

- RMS quantity $\sqrt{\int_{\mathbb{R}} |\hat{f}^{(k)}(t, v)|^2 dv}$
- δ -f distribution function

- Link with KEEN waves formation

- consider weak drive $a_{Dr} = 0.00625$
- change the duration of the drive : $T_{Dr} \in \{100, 150, 200\}$

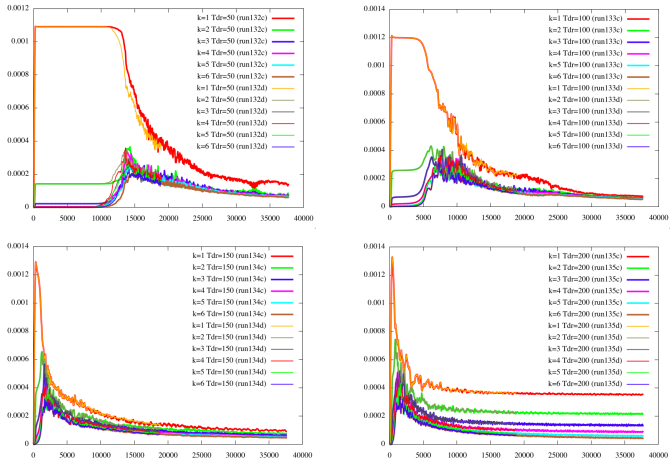


FIGURE: Time evolution of $\sqrt{\int_{\mathbb{R}} |\hat{f}^{(k)}(t, v)|^2 dv}$ for $k = 1, 2, \dots, 6$ and for drive time 50, 100, 150, 200 (run132c-135c and run132d-135d) Parameters are $N_x = 4096$, $N_v = 32768$, non-uniform, 6-th order time scheme and $\Delta t = \mathbf{0.5}$ for run132c-135c, $\Delta t = \mathbf{0.25}$ for run132d-135d.

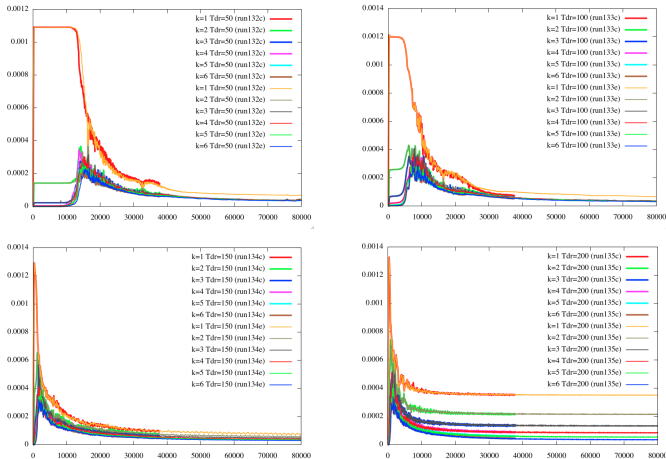


FIGURE: Time evolution of $\sqrt{\int_{\mathbb{R}} |\hat{f}^{(k)}(t, v)|^2 dv}$ for $k = 1, 2, \dots, 6$ and for drive time 50, 100, 150, 200 (run132c-135c and run132e-135e) Parameters are $\Delta t = 0.5$, non-uniform, 6-th order time scheme and $(N_x, N_v) = (4096, 32768)$ for run132c-135c, $(N_x, N_v) = (2048, 16384)$ for run132e-135e.

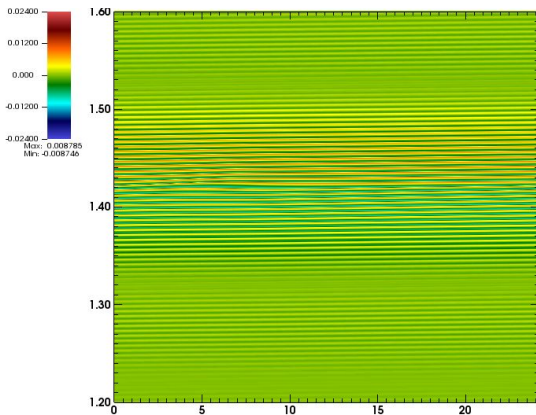


FIGURE: Case $a_{Dr} = 0.00625$ and $T_{Dr} = \mathbf{100}$. δf distribution function $f - f_0$ at time $T = 5000$ as a function of $(x, v) \in [0, 4\pi] \times [1.2, 1.6]$ (from top to bottom). Parameters are $N_x = 2048$, $N_v = 16384$, $\Delta t = 0.5$ non-uniform, 6-th order time scheme (run13).

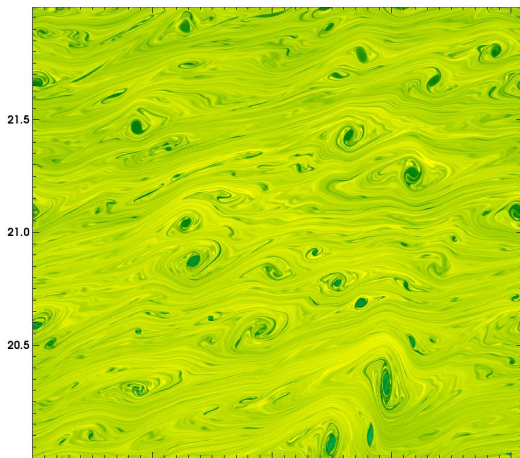


FIGURE: Details (1/3) of δf distribution function $(f - f_0)(x_i, v_j)$ at time $T = 36000$ as a function of $(i, j) \in [0, N_x] \times [16000, 22000]$ Parameters are $N_x = 4096$, $N_v = 32768$, $\Delta t = 0.5$, non-uniform, 6-th order time scheme, $T_{Dr} = \mathbf{100}$, $a_{Dr} = 0.00625$ (run133c).

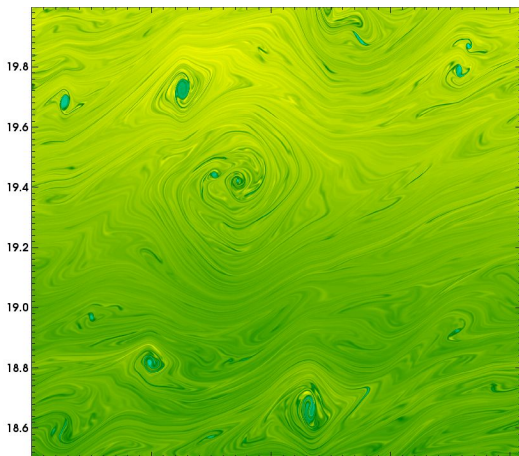


FIGURE: Details (2/3) of δf distribution function $(f - f_0)(x_i, v_j)$ at time $T = 36000$ as a function of $(i, j) \in [0, N_x] \times [16000, 22000]$ Parameters are $N_x = 4096$, $N_v = 32768$, $\Delta t = 0.5$, non-uniform, 6-th order time scheme, $T_{Dr} = \mathbf{100}$, $a_{Dr} = 0.00625$ (run133c).

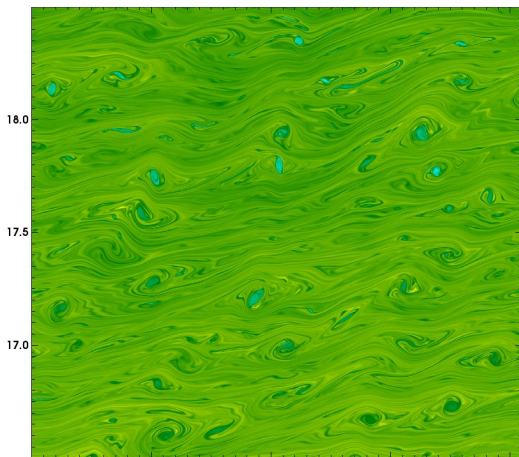


FIGURE: Details (3/3) of δf distribution function $(f - f_0)(x_i, v_j)$ at time $T = 36000$ as a function of $(i, j) \in [0, N_x] \times [16000, 22000]$ Parameters are $N_x = 4096$, $N_v = 32768$, $\Delta t = 0.5$, non-uniform, 6-th order time scheme, $T_{Dr} = \mathbf{100}$, $a_{Dr} = 0.00625$ (run133c).

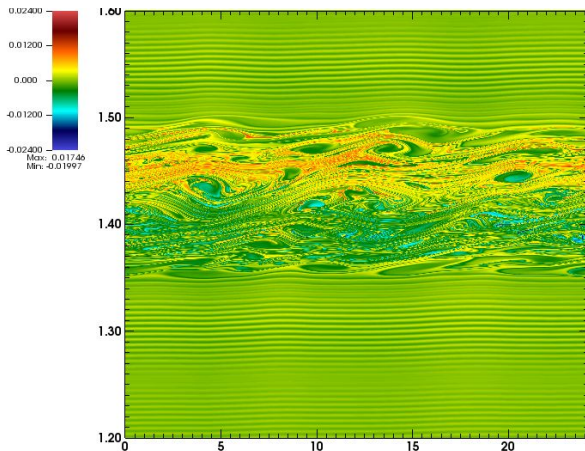


FIGURE: $a_{\text{dr}} = 0.00625$ and $T_{\text{dr}} = \mathbf{150}$ Parameters are $N_x = 2048$, $N_y = 16384$, $\Delta t = 0.5$ non-uniform, 6-th order time scheme.

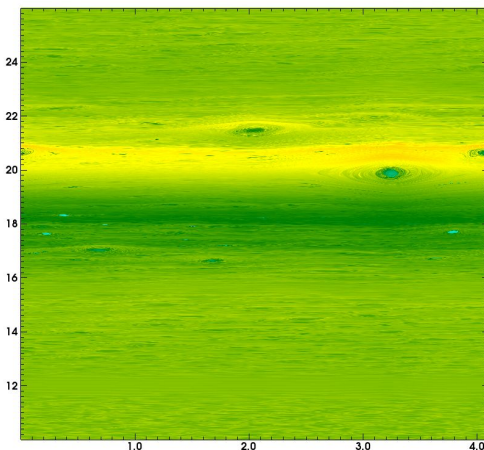


FIGURE: δf distribution function $(f - f_0)(x_i, v_j)$ at time $T = 36000$ as a function of $(i, j) \in [0, N_x] \times [13000, 26000]$. Parameters are $N_x = 4096$, $N_v = 32768$, $\Delta t = 0.5$, non-uniform, 6-th order time scheme, $T_{\text{Dr}} = \mathbf{150}$, $a_{\text{Dr}} = 0.00625$ (run134c).

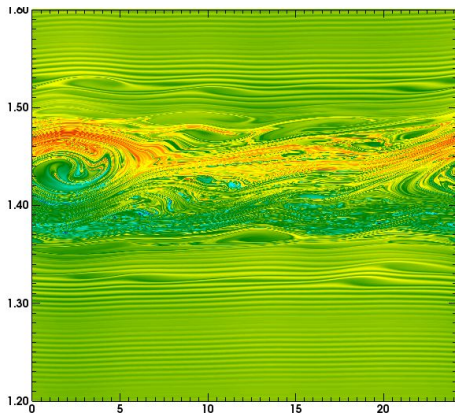


FIGURE: δf distribution function $f - f_0$ at time $T = 5000$ as a function of $(x, v) \in [0, 4\pi] \times [1.2, 1.6]$ $T_{\text{Dr}} = \mathbf{200}$, $a_{\text{Dr}} = 0.00625$: run with 6th order time scheme and non-uniform velocity mesh. Parameters are $N_x = 2048$, $N_v = 16384$, $\Delta t = 0.5$ non-uniform, 6-th order time scheme (run13).

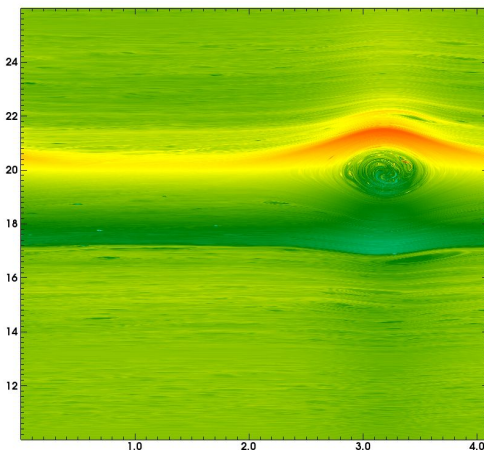


FIGURE: δf distribution function $(f - f_0)(x_i, v_j)$ at time $T = 36000$ as a function of $(i, j) \in [0, N_x] \times [13000, 26000]$. Parameters are $N_x = 4096$, $N_v = 32768$, $\Delta t = 0.5$, non-uniform, 6-th order time scheme, $T_{\text{Dr}} = \mathbf{200}$, $a_{\text{Dr}} = 0.00625$ (run135c).

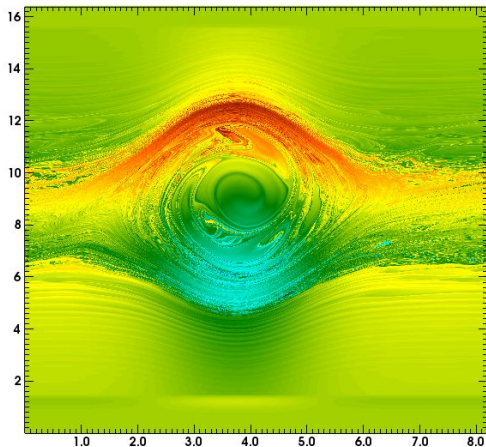


FIGURE: δf distribution function $(f - f_0)(x_i, v_j)$ at time $T = 1000$ as a function of $(i, j) \in [0, N_x] \times [0, N_v]$, in the **canonical drive** case : canonical run with 6th order time scheme and non-uniform velocity mesh. Parameters are $N_x = 8192$, $N_v = 16384$, $\Delta t = 0.25$ non-uniform, 6-th order time scheme (run77).

Conclusion

- Numerical simulation of KEEN waves with weak drive and different duration
- Study of convergence and long time behavior
- Proliferation of tiny vortices which have not coalesced vs well formed KEEN wave
- rms diagnostic helps to distinguish
- development of new numerical methods relevant for such test problem

Possible perspectives

- 4D simulations
- Variations of physics parameters and diagnostics
- Comparisons of numerical methods
 - discontinuous Galerkin (classical/semi-Lagrangian)
 - PIC
- Error computations, restart strategies, enhanced movie generation
- Improve performance ; multi-GPU/MIC
- How to choose the numerical parameters ?
- Convergence vs right long term averaged quantities

Some references

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